

BUCKLING ANALYSIS OF RECTANGULAR PLATES WITH CENTERED CUT-OUT SUBJECTED TO IN-PLANE TWO DIRECTIONS UNDER DIFFERENT BOUNDARY CONDITIONS

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Received 18 July 2018; accepted 18 January 2019

Abstract: In this research, buckling behavior of rectangular plates of symmetric and anti-symmetric laminated composite having centered circular hole under in-plane static loadings is analyzed with the aid of first shear deformation theory and the finite element method is used to find critical loads. The presence of hole may cause redistribution of stresses in plates with reduction of stability. The aim of the current paper is to find critical buckling load. The loads depend on many parameters like geometric aspect ratios (a/b) and (d/b), plate thickness (t), diameter of the circular hole (d), orientation of ply and boundary conditions. Numerical simulations for various boundary conditions obtained are shown in tables and graphical forms and compared with each other.

Keywords: Laminated composite, Buckling, Finite element, Boundary conditions

1. Introduction

Laminated composite plates are widely used in many engineering structures including automotive bridges wind-turbine blades and aviation, due to their lower weights compared to metal structures. Frequently, such plates contain holes of various shapes, for example for access or for weight savings and this perforation or openings often require some type of reinforcing structure to control local structural deformations and stresses near the cutout. In addition, these structures may experience compression loads during operation, and thus their buckling response characteristics must be

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understood and accurately predicted in order to determine effective designs and safe operating conditions for these structures. Thus in the design context buckling analysis plays a crucial role.

Buckling of rectangular plates it loss of stability [1] and it was widely discussed in literature, Reza Eslami [2], Timoshenko and Gere [3], Waszczyszyn et al. [4]. The existence of analytical solutions for certain special cases of loads and boundary conditions creates a good basis examination of numerical results Wang et al. [5]. Hence, the analytical buckling solutions are inaccessible, and various numerical methods need to be developed for analyzing these types of plates. In the literature, few published studies investigated the buckling of laminated composite plates with a cutout [6], [7].

Lin and Kuo [6] presented a nine-node Lagrangian finite element analysis for the buckling of cross-ply and angle-ply-laminated plates with circular holes under in-plane static loadings. Altan and Kartal [7] also used the finite element approach for investigating the buckling of reinforced concrete plates with central rectangular hole.

Schlack investigated the elastic stability of pierced square plates by analytical method and experimental [8], [9]. Falkowicz and Dębski [10] studied numerical and experimental stability of rectangular plate with rectangular cut-out. Mouhat and et al. [11], [12] studied dynamical buckling rectangular stiffened plates with geometric imperfection and effect of pulse duration.

This paper deals with buckling of symmetrically and anti-symmetrically laminated composite plates with circular hole under biaxial static loadings, *Fig. 1*. The First Shear Deformation Theory (FSDT) plate deformation theory and the variational energy method are employed in mathematical formulation, and the finite element technique is used for finding critical loads. The effects on critical load by whole size, plate thickness ratio, material modulus ratio, ply-lamination geometry, loading types, and boundary conditions are investigated. The results are shown in the graphical form for various boundary conditions.

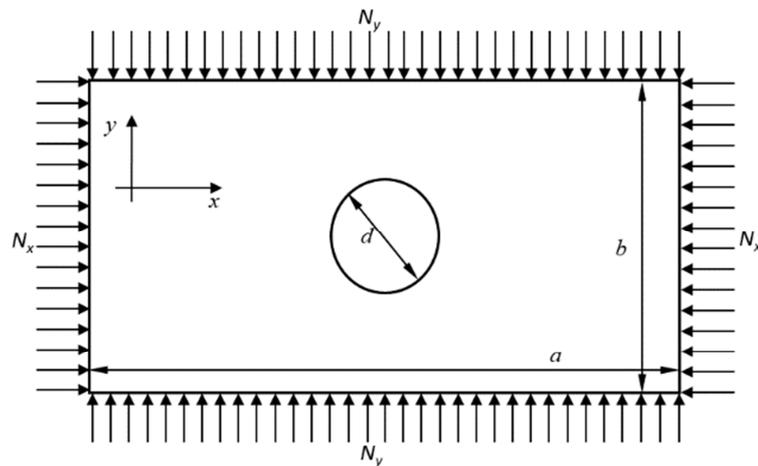


Fig. 1. Rectangular plate under biaxial compressive load ($N_x = -N_0$ and $N_y = -\xi N_0$)

2. Theoretical background

2.1. Kinematics

Consider a symmetrically or anti-symmetrically laminated rectangular plate of length a , width b , central circular hole of diameter d and thickness t which consists of six orthotropic laminas and fibers orientation $[(90^\circ/+ \theta / - \theta)]_{\text{sym}}$ for symmetric and $[(90^\circ/+ \theta / - \theta)]_{\text{asym}}$ for anti-symmetric.

The displacement field according to the FSDT can be expressed in the form [13], Fig. 2:

$$\begin{aligned} u(x, y, z) &= u_0(x, y) + z\phi_x(z, y), \\ v(x, y, z) &= v_0(x, y) + z\phi_y(z, y), \\ w(x, y, z) &= w_0(x, y), \end{aligned} \tag{1}$$

where $u_0(x, y)$, $v_0(x, y)$, $w_0(x, y)$ denote the displacements of a point on the plane $z = 0$, and ϕ_x and ϕ_y are the rotations of a transverse normal about the y and x axes, respectively, Fig. 2,

$$\frac{\partial u}{\partial z} = \phi_x, \quad \frac{\partial v}{\partial z} = \phi_y \tag{2}$$

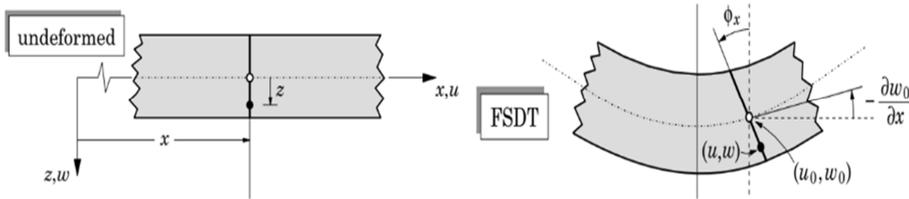


Fig. 2. Deformation of a transverse normal according to first shear deformation plate theory

The strains associated with the displacement field (1) are given by

$$\begin{aligned} \epsilon_{xx} &= \frac{\partial u_0}{\partial x} + z \frac{\partial \phi_x}{\partial x}, \\ \epsilon_{yy} &= \frac{\partial v_0}{\partial y} + z \frac{\partial \phi_y}{\partial y}, \end{aligned} \tag{3a}$$

$$\gamma_{xy} = \left(\frac{\partial u_0}{\partial y} + \frac{\partial v_0}{\partial x} + \frac{\partial w_0}{\partial x} \frac{\partial w_0}{\partial y} \right) + z \left(\frac{\partial \phi_x}{\partial y} + \frac{\partial \phi_y}{\partial x} \right),$$

$$\gamma_{xz} = \frac{\partial w_0}{\partial x} + \phi_x, \quad \gamma_{yz} = \frac{\partial w_0}{\partial y} + \phi_y, \quad \varepsilon_{zz} = 0. \quad (3b)$$

The constitutive equation for an orthotropic layer can be expressed by [14]

$$\begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \tau_{xy} \end{bmatrix} = \begin{bmatrix} Q_{11} & Q_{12} & Q_{16} \\ Q_{12} & Q_{22} & Q_{26} \\ Q_{16} & Q_{26} & Q_{66} \end{bmatrix} \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{yy} \\ \gamma_{xy} \end{bmatrix}, \quad \begin{bmatrix} \tau_{yz} \\ \tau_{xz} \end{bmatrix} = \begin{bmatrix} Q_{44} & Q_{45} \\ Q_{45} & Q_{55} \end{bmatrix} \begin{bmatrix} \gamma_{yz} \\ \gamma_{xz} \end{bmatrix}. \quad (4)$$

The resulting forces and moments occurred in the plate are obtained by integrating the stresses through the thickness of laminas [6], [7].

$$\begin{bmatrix} N_{xx} & M_{xx} \\ N_{yy} & M_{yy} \\ N_{xy} & M_{xy} \end{bmatrix} = \int_{-t/2}^{t/2} \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \tau_{xy} \end{bmatrix} [1 \quad z] dz = \sum_{i=1}^n \int_{t_{i-1}}^{t_i} \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \tau_{xy} \end{bmatrix} [1 \quad z] dz, \quad (5)$$

and the shear forces to the shear strains can be computed as:

$$\begin{bmatrix} R_{yz} \\ R_{xz} \end{bmatrix} = k_s \int_{-t/2}^{t/2} \begin{bmatrix} \tau_{yz} \\ \tau_{xz} \end{bmatrix} dz = k_s \sum_{i=1}^n \int_{t_{i-1}}^{t_i} \begin{bmatrix} \tau_{yz} \\ \tau_{xz} \end{bmatrix} dz, \quad (6)$$

where $k_s (= 5/6)$ is the shear correction factor [15], [16].

For the buckling analysis, the in-plane compression edge forces are assumed the only applied loads (*Fig. 1*), defined by:

$$N_x = -N_0, \quad N_y = -\xi N_0. \quad (7)$$

The equation set that solves the buckling problem takes the form:

$$\frac{\partial M_{xx}}{\partial x} + \frac{\partial M_{xy}}{\partial y} - R_{xy} = 0, \quad (8a)$$

$$\frac{\partial M_{xx}}{\partial x} + \frac{\partial M_{xy}}{\partial y} - R_{yz} = 0, \tag{8b}$$

$$\frac{\partial R_{xz}}{\partial x} + \frac{\partial R_{yz}}{\partial y} - N_0 \left(\frac{\partial^2 w}{\partial x^2} - \xi \frac{\partial^2 w}{\partial y^2} \right) R_{xy} = 0. \tag{8c}$$

2.2. Finite element model

The total potential energy Π of a laminated plate under biaxial loadings is equal to strain energy U can be written as [4], [17], [18]:

$$\Pi = U_b + U_s + V, \tag{9}$$

where U_b is bending term, U_s is the strain energy of shear and V denotes the potential energy of in-plane forces.

The equilibrium requires that the potential energy Π must be stationary. The equilibrium equations of the cross-ply laminated plate can be derived from the variational principle by using

$$\delta \Pi = 0. \tag{10}$$

The finite element method is applied in this study; the following linear algebraic equations are obtained by using equations (8), (9) and (10):

$$([K] + \lambda[K_G])[U] = 0, \tag{11}$$

where $[K]$ is the stiffness matrix, $[K_G]$ is the geometric stiffness matrix, λ is the eigenvalue and $[U]$ is the nodal displacements. The critical buckling load $N_{x,cr}$ is that which corresponds to the least eigenvalue λ determined by:

$$\det([K] + \lambda[K_G]) = 0, \tag{12}$$

and the stiffness components is based on the following 13 weighted-residual statements of the FSDT [19] and [20]:

$$\int_{\Omega_e} \left(\frac{\partial \bar{W}_1}{\partial x} N_{xx} + \frac{\partial \bar{W}_1}{\partial y} N_{xy} \right) dx dy - \int_{\Gamma_e} \bar{W}_1 (n_x N_{xx} + n_y N_{xy}) ds = 0,$$

$$\int_{\Omega_e} \left(\frac{\partial \bar{W}_2}{\partial x} N_{xy} + \frac{\partial \bar{W}_2}{\partial y} N_{yy} \right) dx dy - \int_{\Gamma_e} \bar{W}_2 (n_x N_{xy} + n_y N_{yy}) ds = 0,$$

$$\begin{aligned}
& \int_{\Omega_e} \left\{ \frac{\partial \bar{W}_3}{\partial x} \left(N_{xx} \frac{\partial w_0}{\partial x} + N_{xy} \frac{\partial w_0}{\partial y} \right) + \frac{\partial \bar{W}_3}{\partial y} \left(N_{xy} \frac{\partial w_0}{\partial x} + N_{yy} \frac{\partial w_0}{\partial y} \right) \right\} dx dy \\
& + \int_{\Gamma_e} \bar{W}_3 \left\{ \left(R_{xz} + N_{xx} \frac{\partial w_0}{\partial x} + N_{xy} \frac{\partial w_0}{\partial y} \right) n_x + \left(R_{yz} + N_{xy} \frac{\partial w_0}{\partial x} + N_{yy} \frac{\partial w_0}{\partial y} \right) n_y \right\} ds = 0, \\
& \int_{\Omega_e} \left(\frac{\partial \bar{W}_4}{\partial x} M_{xx} + \frac{\partial \bar{W}_4}{\partial y} M_{xy} + \bar{W}_4 R_{xz} \right) dx dy + \int_{\Gamma_e} \bar{W}_4 (M_{xx} n_x + M_{xy} n_y) ds = 0, \\
& \int_{\Omega_e} \left(\frac{\partial \bar{W}_5}{\partial x} M_{xy} + \frac{\partial \bar{W}_5}{\partial y} M_{yy} + \bar{W}_5 R_{yz} \right) dx dy + \int_{\Gamma_e} \bar{W}_5 (M_{xy} n_x + M_{yy} n_y) ds = 0, \\
& \int_{\Omega_e} \bar{W}_6 \left(-A_{11}^* N_{xx} - A_{12}^* N_{yy} + \frac{\partial u_0}{\partial x} \right) dx dy = 0, \tag{13} \\
& \int_{\Omega_e} \bar{W}_7 \left(-A_{12}^* N_{xx} - A_{22}^* N_{yy} + \frac{\partial v_0}{\partial x} \right) dx dy = 0, \\
& \int_{\Omega_e} \bar{W}_8 \left(-A_{66}^* N_{xy} + \frac{\partial u_0}{\partial y} + \frac{\partial v_0}{\partial x} + \frac{\partial w_0}{\partial x} \frac{\partial w_0}{\partial y} \right) dx dy = 0, \\
& \int_{\Omega_e} \bar{W}_9 \left(-\frac{R_{xz}}{k_s A_{55}} + \frac{\partial w_0}{\partial x} + \phi_x \right) dx dy = 0, \\
& \int_{\Omega_e} \bar{W}_{10} \left(-\frac{R_{yz}}{k_s A_{44}} + \frac{\partial w_0}{\partial y} + \phi_y \right) dx dy = 0, \\
& \int_{\Omega_e} \bar{W}_{11} \left(-D_{11}^* M_{xx} - D_{12}^* M_{yy} + \frac{\partial \phi_x}{\partial x} \right) dx dy = 0, \\
& \int_{\Omega_e} \bar{W}_{12} \left(-D_{12}^* M_{xx} - D_{22}^* M_{yy} + \frac{\partial \phi_y}{\partial y} \right) dx dy = 0, \\
& \int_{\Omega_e} \bar{W}_{13} \left(-D_{66}^* M_{xy} + \frac{\partial \phi_x}{\partial y} + \frac{\partial \phi_y}{\partial x} \right) dx dy = 0,
\end{aligned}$$

where Ω_e and Γ_e denote the element region and its boundary, respectively.

3. Results and discussion

The numerical simulations and appropriate calculations have been conducted using the finite element software ANSYS [21], [22]. The finite element SHELL181 has been used for discretization of created multi-layered composite plate model (Fig. 3). This is four nodes element with six degrees of freedom at each. It is suitable for analyzing geometrically nonlinear problems and modeling of different material properties. Its option *Shell Section Type* gives a possibility of defining a multi-layered cross-section, their thickness, number of integration points across each layer thickness and of introducing different material properties for separate layer. This approach of modeling orthotropic plate as multi-layered one is common in Finite Element Method (FEM) buckling analysis [7] and [18].

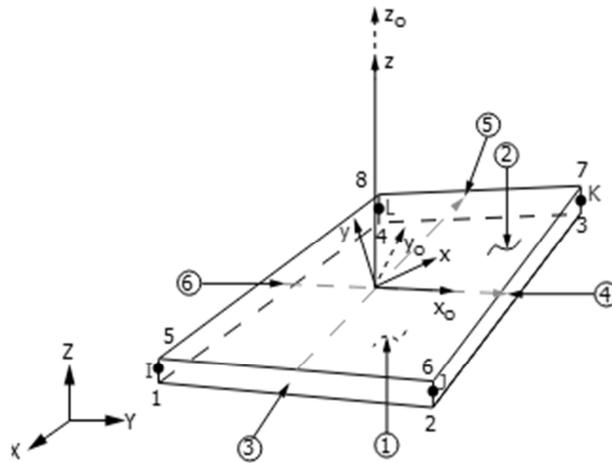


Fig. 3. SHELL 181 element geometry [22]

The composite laminated plates were analysed under three boundary conditions:

- Simple-Simple-Simple-Simple (SSSS)

$$\begin{aligned}
 x = 0, \quad w(0, y) = 0, \quad M_{xx}(0, y) = 0, \\
 x = a, \quad w(a, y) = 0, \quad M_{xx}(a, y) = 0, \\
 y = 0, \quad w(x, 0) = 0, \quad M_{xx}(x, 0) = 0, \\
 y = b, \quad w(x, b) = 0, \quad M_{xx}(x, b) = 0,
 \end{aligned} \tag{14}$$

- Simple-Simple-Simple-Free (SSSF)

$$\begin{aligned}
 y = 0, \quad w(x,0) = 0, \quad M_{yy}(x,0) = 0, \\
 y = b, \quad M_{yy}(x,b) = 0, \quad D_{22} \frac{\partial^3 w}{\partial y^3} + (D_{12} + 4D_{66}) \frac{\partial^3 w}{\partial x^2 \partial y} = 0,
 \end{aligned}
 \tag{15}$$

- Simple-Clamped-Simple-Free (SCSF)

$$\begin{aligned}
 y = 0, \quad w(x,0) = 0, \quad \frac{\partial w}{\partial y} = 0, \\
 y = b, \quad M_{yy}(x,b) = 0, \quad D_{22} \frac{\partial^3 w}{\partial y^3} + (D_{12} + 4D_{66}) \frac{\partial^3 w}{\partial x^2 \partial y} = 0.
 \end{aligned}
 \tag{16}$$

The mechanical properties of material used in this study are given in *Table I*.

Table I

Mechanical Properties of the laminated plate in E-Glass/Epoxy

Mechanical properties		Values
E ₁		5.49E+006 Psi
E ₂		9.53E+005 Psi
E ₃		9.53E+005 Psi
G ₁₂ =G ₁₃		3.47E+005 Psi
v ₁₂ =v ₁₃		0.3001
v ₂₃ =v ₃₂		0.3

Results obtained are shown in *Table II*, *Table III* and *Table IV*.

Table II

Critical buckling load of the laminated plates of [(90°/+θ/-θ)]_{sym} and [(90°/+θ/-θ)]_{asym} under the boundary conditions Simple-Simple-Simple-Simple (SSSS)

Plate dimensions (in)				N _{x,cr} (kpsi)							
				Symmetric				Anti-Symmetric			
a	b	t	d	15°	30°	45°	60°	15°	30°	45°	60°
6	6	0.23	0.312	1.132	1.19	1.25	1.25	1.13	1.195	1.256	1.24
	5			1.403	1.49	1.569	1.567	1.383	1.467	1.549	1.538
	4			1.845	1.927	1.987	1.949	1.826	1.911	1.98	1.933
	3			2.8	2.89	2.902	2.766	2.746	2.824	2.848	2.719
	6	0.3		2.471	2.588	2.717	2.71	2.46	2.597	2.728	2.695
	5			3.054	3.23	3.399	3.387	2.997	3.174	3.350	3.324
	4			3.983	4.153	4.277	4.187	3.931	4.108	4.256	4.152
	3			5.974	6.164	6.189	5.8964	5.845	6.003	6.069	5.803
	6	0.355		4.023	4.215	4.42	4.405	4.006	4.224	4.437	4.377
	5			4.952	5.246	5.52	5.488	4.86	5.141	5.428	5.386
	4			6.441	6.707	6.906	6.754	6.34	6.614	6.858	6.69
	3			9.566	9.86	9.908	9.448	9.326	9.567	9.701	9.302

Table III

Critical buckling load of the laminated plates of $[(90^\circ/+0^\circ/-0^\circ)]_{\text{sym}}$ and $[(90^\circ/+0^\circ/-0^\circ)]_{\text{asym}}$ under the boundary conditions Simple-Simple-Simple-Free (SSSF)

Plate dimensions (in)				$N_{x,cr}$ (kpsi)							
a	b	t	d	Symmetric				Anti-Symmetric			
				15°	30°	45°	60°	15°	30°	45°	60°
6	6	0.23	0.312	0.749	0.763	0.756	0.732	0.75	0.769	0.764	0.736
	5			0.724	0.745	0.739	0.709	0.723	0.751	0.749	0.714
	4			0.678	0.721	0.736	0.708	0.682	0.733	0.747	0.709
	3			0.626	0.692	0.724	0.688	0.634	0.709	0.738	0.692
	6	0.3		1.633	1.665	1.647	1.593	1.634	1.678	1.664	1.602
	5			1.575	1.623	1.611	1.542	1.572	1.63	1.631	1.552
	4			1.477	1.571	1.602	1.539	1.484	1.597	1.625	1.541
	3			1.361	1.5045	1.573	1.491	1.378	1.54	1.602	1.5018
	6	0.355		2.662	2.717	2.686	2.593	2.662	2.737	2.712	2.608
	5			2.566	2.647	2.626	2.509	2.559	2.667	2.657	2.525
	4			2.406	2.561	2.609	2.504	2.417	2.602	2.645	2.506
	3			2.216	2.449	2.557	2.422	2.243	2.505	2.603	2.439

Table IV

Critical buckling load of the laminated plates of $[(90^\circ/+0^\circ/-0^\circ)]_{\text{sym}}$ and $[(90^\circ/+0^\circ/-0^\circ)]_{\text{asym}}$ under the boundary conditions Simple-Clamped-Simple-Free (SCSF)

Plate dimensions (in)				$N_{x,cr}$ (kpsi)							
a	b	t	d	Symmetric				Anti-Symmetric			
				15°	30°	45°	60°	15°	30°	45°	60°
6	6	0.23	0.312	0.792	0.816	0.804	0.769	0.795	0.827	0.816	0.745
	5			0.804	0.829	0.812	0.762	0.801	0.835	0.819	0.763
	4			0.84	0.879	0.864	0.848	0.839	0.886	0.869	0.867
	3			0.992	1.038	1.012	0.911	0.989	1.039	1.0131	0.911
	6	0.3		1.723	1.778	1.752	1.671	1.729	1.799	1.775	1.683
	5			1.745	1.805	1.765	1.653	1.739	1.814	1.779	1.657
	4			1.825	1.909	1.876	1.734	1.822	1.922	1.886	1.729
	3			2.147	2.247	2.186	1.968	2.141	2.248	2.188	1.968
	6	0.355		2.806	2.897	2.887	2.717	2.814	2.929	2.887	2.735
	5			2.839	2.938	2.873	2.685	2.826	2.949	2.892	2.692
	4			2.968	3.106	3.048	2.816	2.961	3.124	3.063	2.808
	3			3.485	3.642	3.541	3.188	3.473	3.642	3.542	3.187

It is easily seemed from Tables that while thickness increase the critical buckling load $N_{x,cr}$ increase. Fig. 4 refers that critical buckling of the boundary condition (SSSS) plates are approximately 2 times greater than the buckling load of (SSSF) and (SCSF)

plates. The effects of a/b ratio on buckling load are illustrated in *Fig. 5*, according to it the buckling loads increase with increasing of a/b ratios. Furthermore, the highest values of buckling loads are calculated $[(90^\circ/60^\circ/-60^\circ)]_{\text{sym}}$ (SSSS) plates, while the lower values are computed angle-ply $[(90^\circ/+ \theta^\circ/-\theta^\circ)]_{\text{sym}}$ (SSSF) plates. The effects of ply orientation symmetric and anti-symmetric on the buckling load can be seen in *Table II*, *Table III* and *Table IV*. The results show that the change in buckling load is neglected. The effect of circular cut-out is taken into account. The analysis indicates that the variation of the buckling loads is very sensitive to the presence of hole. It can be seen that the buckling load generally decreases with presence of cut-out (*Fig. 5*), the predicted buckling loads of plates without hole is 70% higher than with hole in (SSSS) plates and 57% in (SSSF) plate and 67% in (SCSF) plate.

In *Fig. 6*, effects of hole size on critical buckling loads of rectangular orthotropic plates under biaxial compressions are compared with those obtained by El Youbi et al. [18]. It can be seen that in dashed circle (*Fig. 6*), that the influence of hole is more significant for thinner plate $t=0.23$ in.

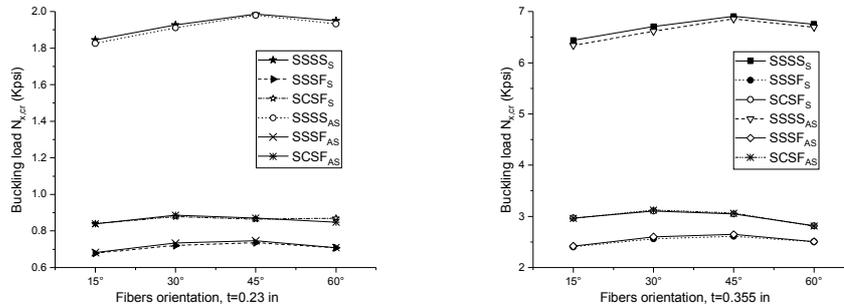


Fig. 4. Variation of the critical buckling load $N_{x,cr}$ compared to the fibers orientation symmetric and anti-symmetric of a plate 6×4 in² for thicknesses $t=0.23$ in and $t=0.355$ in

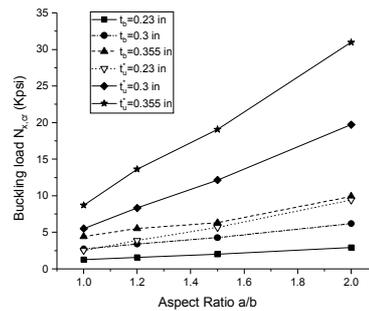


Fig. 5. Variation of the critical buckling load $N_{x,cr}$ in function of aspect ratio a/b of biaxial plate compression (t_b) compared to uniaxial load plate (t_u)

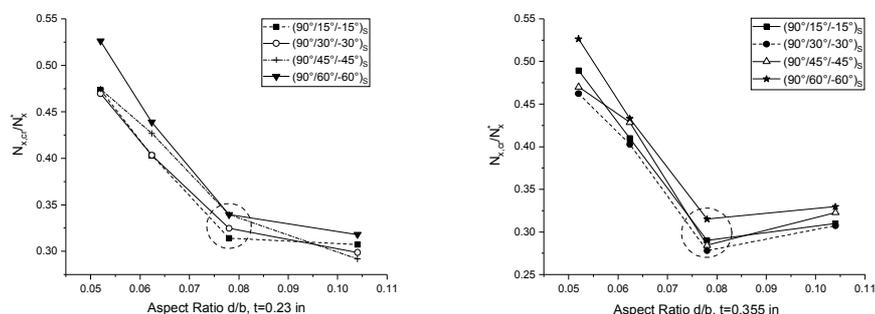


Fig. 6. Hole effect on critical load of $[(90^\circ/\theta/-\theta)]_{\text{sym}}$ cross-ply under boundary conditions (SSSS) at different thickness ratio a/b under biaxial compression, N_x^* is the critical load of plate without hole [18]

4. Conclusion

In this study, the buckling load of E-Glass/Epoxy composite laminated rectangular plate with central circular cut-out bi-axially loaded is investigated. Additionally, the effect of a/b and t/b ratios on buckling loads are calculated. Solutions were done with FEM using ANSYS finite element program. From the present study, the following conclusions can be made:

1. The reduction of the buckling load due to hole is found to be significant, this means big holes cause the weakest plates under the pressure;
2. The increasing of buckling loads by increasing a/b ratios;
3. Boundary conditions of composite have the strongest effect on the buckling load, among the other parameters;
4. Buckling is loss of stability due to geometric effects and boundary conditions than the material failure;
5. The cross-ply $[(90^\circ/45^\circ/-45^\circ)]_{\text{sym}}$ composite plates is stronger than other analyzed angle-ply laminated plates;
6. The $[(90^\circ/15^\circ/-15^\circ)]_{\text{sym}}$ laminated plate is the weakest angle-ply plate;
7. From this study, no big difference between the cross-ply and angle-ply laminates with symmetric arrangement and anti-symmetric arrangement of layers.

Comparison of results of the finite element with experimental test has to be done.

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