# Almost sure, $L_{1}$ - and $L_{2}$-growth behavior of supercritical multi-type continuous state and continuous time branching processes with immigration 

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#### Abstract

Under a first order moment condition on the immigration mechanism, we show that an appropriately scaled supercritical and irreducible multi-type continuous state and continuous time branching process with immigration (CBI process) converges almost surely. If an $x \log (x)$ moment condition on the branching mechanism does not hold, then the limit is zero. If this $x \log (x)$ moment condition holds, then we prove $L_{1}$ convergence as well. The projection of the limit on any left non-Perron eigenvector of the branching mean matrix is vanishing. If, in addition, a suitable extra power moment condition on the branching mechanism holds, then we provide the correct scaling for the projection of a CBI process on certain left non-Perron eigenvectors of the branching mean matrix in order to have almost sure and $L_{1}$ limit. Moreover, under a second order moment condition on the branching and immigration mechanisms, we prove $L_{2}$ convergence of an appropriately scaled process and the above mentioned projections as well. A representation of the limits is also provided under the same moment conditions.


## 1 Introduction

The description of the asymptotic behavior of branching processes without or with immigration has a long history. For multi-type Galton-Watson processes without immigration see, e.g.,

[^0]Athreya and Ney [3, Sections 4-8 in Chapter V]. For supercritical multi-type Galton-Watson processes with immigration see, e.g., Kaplan [12].

Let us consider a multi-type continuous state and continuous time branching process with immigration (CBI process) which can be represented as a pathwise unique strong solution of the stochastic differential equation (SDE)

$$
\begin{align*}
\boldsymbol{X}_{t}= & \boldsymbol{X}_{0}+\int_{0}^{t}\left(\boldsymbol{\beta}+\widetilde{\boldsymbol{B}} \boldsymbol{X}_{u}\right) \mathrm{d} u+\sum_{\ell=1}^{d} \int_{0}^{t} \sqrt{2 c_{\ell} \max \left\{0, X_{u, \ell}\right\}} \mathrm{d} W_{u, \ell} \boldsymbol{e}_{\ell}  \tag{1.1}\\
& +\sum_{\ell=1}^{d} \int_{0}^{t} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \boldsymbol{z} \mathbb{1}_{\left\{w \leqslant X_{u-, \ell}\right.} \widetilde{N}_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w)+\int_{0}^{t} \int_{\mathcal{U}_{d}} \boldsymbol{r} M(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r})
\end{align*}
$$

for $t \in[0, \infty)$, see, Theorem 4.6 and Section 5 in Barczy et al. [5], where (1.1) was proved only for $d \in\{1,2\}$, but their method clearly works for all $d \in\{1,2, \ldots\}$. Here $d \in$ $\{1,2, \ldots\}$ is the number of types, $X_{t, \ell}, \quad \ell \in\{1, \ldots, d\}$, denotes the $\ell^{t h}$ coordinate of $\boldsymbol{X}_{t}$, $\mathbb{P}\left(\boldsymbol{X}_{0} \in[0, \infty)^{d}\right)=1, \quad \boldsymbol{\beta} \in[0, \infty)^{d}, c_{1}, \ldots, c_{d} \in[0, \infty), \boldsymbol{e}_{1}, \ldots, \boldsymbol{e}_{d}$ denotes the natural basis in $\mathbb{R}^{d}, \quad \mathcal{U}_{d}:=[0, \infty)^{d} \backslash\{(0, \ldots, 0)\}, \quad\left(W_{t, 1}\right)_{t \geqslant 0}, \ldots, \quad\left(W_{t, d}\right)_{t \geqslant 0} \quad$ are independent standard Wiener processes, $N_{\ell}, \quad \ell \in\{1, \ldots, d\}$, and $M$ are independent Poisson random measures on $(0, \infty) \times \mathcal{U}_{d} \times(0, \infty)$ and on $(0, \infty) \times \mathcal{U}_{d}$ with intensity measures $\mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w, \quad \ell \in$ $\{1, \ldots, d\}$, and $\mathrm{d} u \nu(\mathrm{~d} \boldsymbol{r})$, respectively, and $\widetilde{N}_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w):=N_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w)-\mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w$, $\ell \in\{1, \ldots, d\}$. We suppose that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)<\infty$, the Borel measures $\mu_{\ell}, \ell \in\{1, \ldots, d\}$, and $\nu$ on $\mathcal{U}_{d}$ satisfy the moment conditions given in parts (v), (vi) of Definition 2.2 and (2.3), and $\boldsymbol{X}_{0},\left(W_{t, 1}\right)_{t \geqslant 0}, \ldots,\left(W_{t, d}\right)_{t \geqslant 0}, N_{1}, \ldots, N_{d}$ and $M$ are independent. Moreover, $\widetilde{\boldsymbol{B}}=\left(\widetilde{b}_{i, j}\right)_{i, j \in\{1, \ldots, d\}} \in \mathbb{R}^{d \times d}$ is a matrix satisfying $\widetilde{b}_{i, j} \geqslant \int_{\mathcal{U}_{d}} z_{i} \mu_{j}(\mathrm{~d} \boldsymbol{z})$ for all $i, j \in\{1, \ldots, d\}$ with $i \neq j$.

A multi-type CBI process $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is called irreducible if $\widetilde{\boldsymbol{B}}$ is irreducible, see Definition 2.8. An irreducible multi-type CBI process is called subcritical, critical or supercritical if the logarithm $s(\widetilde{\boldsymbol{B}})$ of the Perron eigenvalue of the branching mean matrix $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$ is negative, zero or positive, respectively, see Definition 2.9, A multi-type CBI process $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is called a multi-type CB process if there is no immigration, i.e., $\boldsymbol{\beta}=\mathbf{0}$ and $\nu=0$.

In case of a subcritical or critical single-type CBI process (when it is necessarily irreducible) with a non-vanishing branching mechanism, $\quad \boldsymbol{X}_{t} \xrightarrow{\mathcal{D}} \pi$ as $t \rightarrow \infty$ with a probability measure $\pi$ on $[0, \infty)$ if and only if certain integrability condition holds for the branching and immigration mechanisms, see, e.g., Li [15, Theorem 3.20].

In case of a supercritical single-type CB process, under the $x \log (x)$ moment condition (3.2) with $\lambda=s(\widetilde{\boldsymbol{B}})$ on the branching mechanism, Li [15, Corollary 3.16 and Theorem 3.8] proved that $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t}$ converges almost surely as $t \rightarrow \infty$ towards a non-negative random variable, and the probability that this limit is zero equals to the probability of the event that the extinction time is finite.

In case of a critical and irreducible multi-type CBI process, under fourth order moment conditions on the branching and immigration mechanisms, Barczy and Pap [7, Theorem 4.1]
proved that the sequence $\left(n^{-1} \boldsymbol{X}_{\lfloor n t\rfloor}\right)_{t \in[0, \infty)}, \quad n \in\{1,2, \ldots\}$, of scaled random step functions converges weakly towards a squared Bessel process (in other words, a Feller diffusion) supported by a ray determined by the right Perron vector $\widetilde{\boldsymbol{u}}$ of the branching mean matrix $e^{\widetilde{B}}$.

Recently, there is a renewed interest for studying asymptotic behavior of supercritical branching processes. In case of a supercritical and irreducible multi-type CB process, Kyprianou et al. [14, Theorem 1.3] described the asymptotic behavior of the projection $\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$, where $\boldsymbol{u}$ denotes the left Perron eigenvector of the branching mean matrix $\mathrm{e}^{\widetilde{B}}$. Namely, they proved that if an $x \log (x)$ moment condition on the branching mechanism holds, then $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle \rightarrow w_{\boldsymbol{u}, \boldsymbol{X}_{0}}$ almost surely and in $L_{1}$ as $t \rightarrow \infty$, where $w_{\boldsymbol{u}, \boldsymbol{X}_{0}}$ is a nonnegative random variable, otherwise $\mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle \rightarrow 0$ almost surely as $t \rightarrow \infty$. Note that their $x \log (x)$ moment condition is equivalent to our moment condition (3.2) with $\lambda=s(\widetilde{\boldsymbol{B}})$, since for $\mathbb{R}^{d}$, all norms are equivalent. Moreover, in case of a supercritical and irreducible multi-type CB process, Kyprianou et al. [14, Theorem 1.4] proved that $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \rightarrow w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}}$ almost surely as $t \rightarrow \infty$.

Ren et al. [20] investigated central limit theorems for supercritical branching Markov processes, and Ren et al. [21, 22] studied some properties of strong limits for supercritical superprocesses. Moreover, Chen et al. [8] and Ren et al. [19] studied spine decomposition and an $x \log x$ criterion for supercritical superprocesses with non-local branching mechanisms.

Recently, Marks and Miłoś [17, Theorem 3.2] considered a branching particle system with particles moving according to a multi-dimensional Ornstein-Uhlenbeck process with a positive drift and branching according to a law in the domain of attraction of a stable law having stability index in (1,2), and in the so-called "large branching case" (see [17, page 3]) they proved almost sure and $L_{1}$ convergence of appropriately normalized projections of the particle system in question onto certain twice differentiable real-valued functions defined on the real line of polynomial growth together with a description of the limit in which the whole genealogical structure is somewhat preserved. These projections include projections onto certain eigenfunctions of the semigroup associated to the infinitesimal generator of the underlying Ornstein-Uhlenbeck process.

Very recently, Ren et al. [18] derived stable central limit theorems for some kind of projections of (measure-valued) super Ornstein-Uhlenbeck processes having a branching mechanism which is close to a function of the form $-a_{1} z+a_{2} z^{2}+a_{3} z^{1+\alpha}, z \geqslant 0$, with $a_{1}>0, a_{2} \geqslant 0$, $a_{3}>0$ and $\alpha \in(0,1)$ in some sense (see the Assumption 2 in Ren et al. [18).

As a new result, in case of a supercritical and irreducible multi-type CBI process, under the first order moment condition (2.3) on the immigration mechanism, we show $\mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \rightarrow$ $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}}$ almost surely as $t \rightarrow \infty$, where $w_{\boldsymbol{u}, \boldsymbol{X}_{0}}$ is a non-negative random variable, see Theorem 3.3. If the $x \log (x)$ moment condition (3.2) with $\lambda=s(\widetilde{\boldsymbol{B}})$ does not hold, then $\mathbb{P}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}=0\right)=1$, see Theorem 3.1. If this $x \log (x)$ moment condition holds, then we prove $L_{1}$ convergence, see Theorem [3.3, and we give a representation of $w_{\boldsymbol{u}, \boldsymbol{X}_{0}}$ as well, see (3.4). Note that $\mathbb{P}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}=0\right)=1$ if and only if $\mathbb{P}\left(\boldsymbol{X}_{t}=\mathbf{0}\right)=1$ for all $t \in \mathbb{R}_{+}$, see Theorem 3.1. Hence here the scaling factor $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}$ is correct. If $\boldsymbol{v}$ is a left non-Perron eigenvector of the
branching mean matrix $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$, then this result implies that $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle \rightarrow w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\langle\boldsymbol{v}, \widetilde{\boldsymbol{u}}\rangle=0$ almost surely as $t \rightarrow \infty$, since $\langle\boldsymbol{v}, \widetilde{\boldsymbol{u}}\rangle=0$ due to the so-called principle of biorthogonality (see, e.g., Horn and Johnson [10, Theorem 1.4.7(a)]), consequently, the scaling factor $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}$ is not appropriate for describing the asymptotic behavior of the projection $\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$. It turns out that, under the extra power moment condition (3.2) with $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)$ on the branching mechanism and the first order moment condition (2.3) on the immigration mechanism, we can show $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle \rightarrow w_{\boldsymbol{v}, \boldsymbol{X}_{0}}$ almost surely and in $L_{1}$ as $t \rightarrow \infty$, where $\lambda$ is a non-Perron eigenvalue of the branching mean matrix $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$ with $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)$, $\boldsymbol{v}$ is a left eigenvector corresponding to $\lambda$, and $w_{\boldsymbol{v}, \boldsymbol{X}_{0}}$ is a complex random variable, see Theorem 3.1, where we give a representation of $w_{\boldsymbol{v}, \boldsymbol{X}_{0}}$ as well, see (3.4). Here the scaling factor $\mathrm{e}^{-\lambda t}$ is correct if $\left\langle\boldsymbol{v}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)+\lambda^{-1} \widetilde{\boldsymbol{\beta}}\right\rangle \neq 0$, since then $\mathbb{P}\left(w_{\boldsymbol{v}, \boldsymbol{X}_{0}}=0\right)<1$, see Theorem 3.1. In Remark 3.2 we explain why we do not have any result in the case when the moment condition (3.2) does not hold for $\lambda \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)$ formulating some open problems as well. Note that the asymptotic behavior of the second moment $\mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right|^{2}\right)$ as $t \rightarrow \infty$ explains the role of the assumption $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$, see Proposition B.1.

Further, in case of a supercritical and irreducible multi-type CBI process, under the second order moment condition (3.55) on the branching and immigration mechanisms, we show $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \rightarrow w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}}$ and $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle \rightarrow w_{\boldsymbol{v}, \boldsymbol{X}_{0}}$ in $L_{2}$ as $t \rightarrow \infty$ as well, where $\lambda$ is a eigenvalue of the branching mean matrix $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$ with $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$ and $\boldsymbol{v}$ is a left eigenvector corresponding to $\lambda$, see Theorem 3.4.

The paper is structured as follows. In Section 2, for completeness and better readability, from Barczy et al. [5], we recall some notions and statements for multi-type CBI processes such as a formula for their first moment, an appropriate transformation which results in a $d$ dimensional martingale in Lemma 2.6, a useful representation of $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$in Lemma 2.7, the definition of subcritical, critical and supercritical irreducible CBI processes, see Definitions 2.8 and 2.9. Section 3 contains our main results detailed above, see Theorems 3.1, 3.3 and 3.4. For the proofs, we use heavily the representation of $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$in Lemma 2.7 based on the SDE (1.1). In the course of the proof of Theorem [3.3, we follow the steps and methods of the proof of Theorem 1.4 in Kyprianou et al. [14]. We close the paper with two Appendices. We present a useful decomposition of a CBI process as an independent sum of a CBI process starting from $\mathbf{0}$ and a CB process, see Appendix A. In Appendix B , we describe the asymptotic behavior of the second moment of $\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right|$ as $t \rightarrow \infty$ for each left eigenvector $\boldsymbol{v} \in \mathbb{C}^{d}$ of $\widetilde{\boldsymbol{B}}$ corresponding to an arbitrary eigenvalue $\lambda \in \sigma(\widetilde{\boldsymbol{B}})$ in case of a supercritical and irreducible CBI process.

Now, we summarize the novelties of the paper. We point out that we investigate the asymptotic behavior of the projections of a multi-type CBI process on certain left non-Perron eigenvectors of its branching mean matrix. According to our knowledge, this type of question has not been studied so far for multi-type CBI processes. A new phenomenon appears compared to the left Perron eigenvector case, namely, a moment type condition on the branching mechanism of the CBI process in question. Furthermore, if the $x \log (x)$ moment condition (3.2) with $\lambda=s(\widetilde{\boldsymbol{B}})$ on the branching mechanism does not hold, then one usually uses a so-called spine decomposition technique in order to show that $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \stackrel{\text { a.s. }}{=} 0$ (see, e.g., the proof
of Theorem 1.3 in Kyprianou et al. [14] or that of Theorem 6.2 in Ren et al. [19]). In this paper, we use that the law of a multi-type CBI process $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$at time $t+T, t, T \in \mathbb{R}_{+}$, coincides with the law of an independent sum of a multi-type CB process at time $t$ starting from an initial value having distribution as that of $\boldsymbol{X}_{T}$ and a multi-type CBI process at time $t$ starting from $\mathbf{0}$, presented in Lemma A.1, and that the corresponding result $w_{\boldsymbol{u}, \boldsymbol{\boldsymbol { X } _ { 0 }}} \stackrel{\text { a.s. }}{=} 0$ is already known for CB processes due to Kyprianou et al. [14, Theorem 1.3].

Finally, we mention a possible extension of the present results which can be a topic of future work. Since the $d$-dimensional matrix $\widetilde{\boldsymbol{B}}$ is not symmetric in general, its left eigenvectors may not generate $\mathbb{C}^{d}$, so it is natural to study the asymptotic behaviour of $\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$, where $\boldsymbol{v}$ is an arbitrary vector in $\mathbb{C}^{d}$. This type of question was investigated by Kesten and Stigun [13] and Badalbaev and Mukhitdinov [4] for supercritical and irreducible multi-type discrete time Galton-Watson processes without immigration under second order moment assumptions, and, by Athreya [1, 2], for supercritical and positively regular multitype continuous time Markov branching processes without immigration under second order moment assumptions. The above mentioned four references are for some branching processes without immigration, we do not know any corresponding result for branching processes with immigration. Motivated by these references, we think that the Jordan normal form of $\widetilde{\boldsymbol{B}}$ may be well-used in our case as well, where we consider multi-type CBI processes with immigration.

## 2 Multi-type CBI processes

Let $\mathbb{Z}_{+}, \mathbb{N}, \mathbb{R}, \mathbb{R}_{+}, \mathbb{R}_{++}$and $\mathbb{C}$ denote the set of non-negative integers, positive integers, real numbers, non-negative real numbers, positive real numbers and complex numbers, respectively. For $x, y \in \mathbb{R}$, we will use the notations $x \wedge y:=\min \{x, y\}$ and $x^{+}:=\max \{0, x\}$. By $\langle\boldsymbol{x}, \boldsymbol{y}\rangle:=\sum_{j=1}^{d} x_{j} \overline{y_{j}}$, we denote the Euclidean inner product of $\boldsymbol{x}=\left(x_{1}, \ldots, x_{d}\right)^{\top} \in \mathbb{C}^{d}$ and $\boldsymbol{y}=\left(y_{1}, \ldots, y_{d}\right)^{\top} \in \mathbb{C}^{d}$, and by $\|\boldsymbol{x}\|$ and $\|\boldsymbol{A}\|$, we denote the induced norm of $\boldsymbol{x} \in \mathbb{C}^{d}$ and $\boldsymbol{A} \in \mathbb{C}^{d \times d}$, respectively. The null vector and the null matrix will be denoted by $\mathbf{0}$. Moreover, $\boldsymbol{I}_{d} \in \mathbb{R}^{d \times d}$ denotes the identity matrix. By $C_{\mathrm{c}}^{2}\left(\mathbb{R}_{+}^{d}, \mathbb{R}\right)$, we denote the set of twice continuously differentiable real-valued functions on $\mathbb{R}_{+}^{d}$ with compact support. Convergence almost surely, in $L_{1}$ and in $L_{2}$ will be denoted by $\xrightarrow{\text { a.s. }}, \xrightarrow{L_{1}}$ and $\xrightarrow{L_{2}}$, respectively. Almost sure equality will be denoted by $\stackrel{\text { a.s. }}{=}$. Throughout this paper, we make the conventions $\int_{a}^{b}:=\int_{(a, b]}$ and $\int_{a}^{\infty}:=\int_{(a, \infty)}$ for any $a, b \in \mathbb{R}$ with $a<b$.
2.1 Definition. A matrix $\boldsymbol{A}=\left(a_{i, j}\right)_{i, j \in\{1, \ldots, d\}} \in \mathbb{R}^{d \times d}$ is called essentially non-negative if $a_{i, j} \in \mathbb{R}_{+}$whenever $i, j \in\{1, \ldots, d\}$ with $i \neq j$, that is, if $\boldsymbol{A}$ has non-negative off-diagonal entries. The set of essentially non-negative $d \times d$ matrices will be denoted by $\mathbb{R}_{(+)}^{d \times d}$.
2.2 Definition. A tuple $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ is called a set of admissible parameters if
(i) $d \in \mathbb{N}$,
(ii) $\boldsymbol{c}=\left(c_{i}\right)_{i \in\{1, \ldots, d\}} \in \mathbb{R}_{+}^{d}$,
(iii) $\boldsymbol{\beta}=\left(\beta_{i}\right)_{i \in\{1, \ldots, d\}} \in \mathbb{R}_{+}^{d}$,
(iv) $\boldsymbol{B}=\left(b_{i, j}\right)_{i, j \in\{1, \ldots, d\}} \in \mathbb{R}_{(+)}^{d \times d}$,
(v) $\nu$ is a Borel measure on $\mathcal{U}_{d}:=\mathbb{R}_{+}^{d} \backslash\{\mathbf{0}\}$ satisfying $\int_{\mathcal{U}_{d}}(1 \wedge\|\boldsymbol{r}\|) \nu(\mathrm{d} \boldsymbol{r})<\infty$,
(vi) $\boldsymbol{\mu}=\left(\mu_{1}, \ldots, \mu_{d}\right)$, where, for each $i \in\{1, \ldots, d\}, \mu_{i}$ is a Borel measure on $\mathcal{U}_{d}$ satisfying

$$
\begin{equation*}
\int_{\mathcal{U}_{d}}\left[\|\boldsymbol{z}\| \wedge\|\boldsymbol{z}\|^{2}+\sum_{j \in\{1, \ldots, d\} \backslash\{i\}}\left(1 \wedge z_{j}\right)\right] \mu_{i}(\mathrm{~d} \boldsymbol{z})<\infty \tag{2.1}
\end{equation*}
$$

2.3 Remark. Our Definition 2.2 of the set of admissible parameters is a special case of Definition 2.6 in Duffie et al. 9], which is suitable for all affine processes, see Barczy et al. [5, Remark 2.3]. Further, due to Remark 2.3 and (2.12) in Barczy et al. [5], the condition (2.1) is equivalent to

$$
\int_{\mathcal{U}_{d}}\left[\|\boldsymbol{z}\| \wedge\|\boldsymbol{z}\|^{2}+\sum_{j \in\{1, \ldots, d\} \backslash\{i\}} z_{j}\right] \mu_{i}(\mathrm{~d} \boldsymbol{z})<\infty
$$

and also to

$$
\int_{\mathcal{U}_{d}}\left[\left(1 \wedge z_{i}\right)^{2}+\sum_{j \in\{1, \ldots, d\} \backslash\{i\}}\left(1 \wedge z_{j}\right)\right] \mu_{i}(\mathrm{~d} \boldsymbol{z})<\infty \quad \text { and } \quad \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\| \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{i}(\mathrm{~d} \boldsymbol{z})<\infty
$$

2.4 Theorem. Let $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ be a set of admissible parameters. Then, there exists a unique conservative transition semigroup $\left(P_{t}\right)_{t \in \mathbb{R}_{+}}$acting on the Banach space (endowed with the supremum norm) of real-valued bounded Borel-measurable functions on the state space $\mathbb{R}_{+}^{d}$ such that its infinitesimal generator is

$$
\begin{aligned}
(\mathcal{A} f)(\boldsymbol{x})= & \sum_{i=1}^{d} c_{i} x_{i} f_{i, i}^{\prime \prime}(\boldsymbol{x})+\left\langle\boldsymbol{\beta}+\boldsymbol{B} \boldsymbol{x}, \boldsymbol{f}^{\prime}(\boldsymbol{x})\right\rangle+\int_{\mathcal{U}_{d}}(f(\boldsymbol{x}+\boldsymbol{r})-f(\boldsymbol{x})) \nu(\mathrm{d} \boldsymbol{r}) \\
& +\sum_{i=1}^{d} x_{i} \int_{\mathcal{U}_{d}}\left(f(\boldsymbol{x}+\boldsymbol{z})-f(\boldsymbol{x})-f_{i}^{\prime}(\boldsymbol{x})\left(1 \wedge z_{i}\right)\right) \mu_{i}(\mathrm{~d} \boldsymbol{z})
\end{aligned}
$$

for $f \in C_{\mathrm{c}}^{2}\left(\mathbb{R}_{+}^{d}, \mathbb{R}\right)$ and $\boldsymbol{x} \in \mathbb{R}_{+}^{d}$, where $f_{i}^{\prime}$ and $f_{i, i}^{\prime \prime}, i \in\{1, \ldots, d\}$, denote the first and second order partial derivatives of $f$ with respect to its $i$-th variable, respectively, and $\boldsymbol{f}^{\prime}(\boldsymbol{x}):=\left(f_{1}^{\prime}(\boldsymbol{x}), \ldots, f_{d}^{\prime}(\boldsymbol{x})\right)^{\top}$. Moreover, the Laplace transform of the transition semigroup $\left(P_{t}\right)_{t \in \mathbb{R}_{+}}$has a representation

$$
\begin{equation*}
\int_{\mathbb{R}_{+}^{d}} \mathrm{e}^{-\langle\boldsymbol{\lambda}, \boldsymbol{y}\rangle} P_{t}(\boldsymbol{x}, \mathrm{~d} \boldsymbol{y})=\mathrm{e}^{-\langle\boldsymbol{x}, \boldsymbol{v}(t, \boldsymbol{\lambda})\rangle-\int_{0}^{t} \psi(\boldsymbol{v}(s, \boldsymbol{\lambda})) \mathrm{d} s}, \quad \boldsymbol{x} \in \mathbb{R}_{+}^{d}, \quad \boldsymbol{\lambda} \in \mathbb{R}_{+}^{d}, \quad t \in \mathbb{R}_{+} \tag{2.2}
\end{equation*}
$$

where, for any $\boldsymbol{\lambda} \in \mathbb{R}_{+}^{d}$, the continuously differentiable function $\mathbb{R}_{+} \ni t \mapsto \boldsymbol{v}(t, \boldsymbol{\lambda})=$ $\left(v_{1}(t, \boldsymbol{\lambda}), \ldots, v_{d}(t, \boldsymbol{\lambda})\right)^{\top} \in \mathbb{R}_{+}^{d}$ is the unique locally bounded solution to the system of differential equations

$$
\partial_{t} v_{i}(t, \boldsymbol{\lambda})=-\varphi_{i}(\boldsymbol{v}(t, \boldsymbol{\lambda})), \quad v_{i}(0, \boldsymbol{\lambda})=\lambda_{i}, \quad i \in\{1, \ldots, d\}
$$

with

$$
\varphi_{i}(\boldsymbol{\lambda}):=c_{i} \lambda_{i}^{2}-\left\langle\boldsymbol{B} \boldsymbol{e}_{i}, \boldsymbol{\lambda}\right\rangle+\int_{\mathcal{U}_{d}}\left(\mathrm{e}^{-\langle\boldsymbol{\lambda}, \boldsymbol{z}\rangle}-1+\lambda_{i}\left(1 \wedge z_{i}\right)\right) \mu_{i}(\mathrm{~d} \boldsymbol{z})
$$

for $\boldsymbol{\lambda} \in \mathbb{R}_{+}^{d}, \quad i \in\{1, \ldots, d\}$, and

$$
\psi(\boldsymbol{\lambda}):=\langle\boldsymbol{\beta}, \boldsymbol{\lambda}\rangle+\int_{\mathcal{U}_{d}}\left(1-\mathrm{e}^{-\langle\boldsymbol{\lambda}, \boldsymbol{r}\rangle}\right) \nu(\mathrm{d} \boldsymbol{r}), \quad \boldsymbol{\lambda} \in \mathbb{R}_{+}^{d}
$$

Theorem 2.4 is a special case of Theorem 2.7 of Duffie et al. [9] with $m=d, n=0$ and zero killing rate. For more details, see Remark 2.5 in Barczy et al. [5].
2.5 Definition. A conservative Markov process with state space $\mathbb{R}_{+}^{d}$ and with transition semigroup $\left(P_{t}\right)_{t \in \mathbb{R}_{+}}$given in Theorem 2.4 is called a multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$. The function $\mathbb{R}_{+}^{d} \ni \boldsymbol{\lambda} \mapsto\left(\varphi_{1}(\boldsymbol{\lambda}), \ldots, \varphi_{d}(\boldsymbol{\lambda})\right)^{\top} \in \mathbb{R}^{d}$ is called its branching mechanism, and the function $\mathbb{R}_{+}^{d} \ni \boldsymbol{\lambda} \mapsto \psi(\boldsymbol{\lambda}) \in \mathbb{R}_{+}$is called its immigration mechanism. A multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ is called a CB process (a continuous state and continuous time branching process without immigration) if $\boldsymbol{\beta}=\mathbf{0}$ and $\nu=0$.

Let $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$be a multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ such that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)<\infty$ and the moment condition

$$
\begin{equation*}
\int_{\mathcal{U}_{d}}\|\boldsymbol{r}\| \mathbb{1}_{\{\|\boldsymbol{r}\| \geqslant 1\}} \nu(\mathrm{d} \boldsymbol{r})<\infty \tag{2.3}
\end{equation*}
$$

holds. Then, by formula (3.4) in Barczy et al. [5],

$$
\begin{equation*}
\mathbb{E}\left(\boldsymbol{X}_{t} \mid \boldsymbol{X}_{0}=\boldsymbol{x}\right)=\mathrm{e}^{t \widetilde{\boldsymbol{B}}} \boldsymbol{x}+\int_{0}^{t} \mathrm{e}^{u \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} u, \quad \boldsymbol{x} \in \mathbb{R}_{+}^{d}, \quad t \in \mathbb{R}_{+} \tag{2.4}
\end{equation*}
$$

where

$$
\widetilde{\boldsymbol{B}}:=\left(\widetilde{b}_{i, j}\right)_{i, j \in\{1, \ldots, d\}}, \quad \widetilde{b}_{i, j}:=b_{i, j}+\int_{\mathcal{U}_{\boldsymbol{d}}}\left(z_{i}-\delta_{i, j}\right)^{+} \mu_{j}(\mathrm{~d} \boldsymbol{z}), \quad \widetilde{\boldsymbol{\beta}}:=\boldsymbol{\beta}+\int_{\mathcal{U}_{d}} \boldsymbol{r} \nu(\mathrm{~d} \boldsymbol{r}),
$$

with $\delta_{i, j}:=1$ if $i=j$, and $\delta_{i, j}:=0$ if $i \neq j$. Note that $\widetilde{\boldsymbol{B}} \in \mathbb{R}_{(+)}^{d \times d}$ and $\widetilde{\boldsymbol{\beta}} \in \mathbb{R}_{+}^{d}$, since

$$
\int_{\mathcal{U}_{d}}\|\boldsymbol{r}\| \nu(\mathrm{d} \boldsymbol{r})<\infty, \quad \int_{\mathcal{U}_{d}}\left(z_{i}-\delta_{i, j}\right)^{+} \mu_{j}(\mathrm{~d} \boldsymbol{z})<\infty, \quad i, j \in\{1, \ldots, d\}
$$

see Barczy et al. [5, Section 2]. Further, $\mathbb{E}\left(\boldsymbol{X}_{t} \mid \boldsymbol{X}_{0}=\boldsymbol{x}\right), \boldsymbol{x} \in \mathbb{R}_{+}^{d}$, does not depend on the parameter $\boldsymbol{c}$. One can give probabilistic interpretations of the modified parameters $\widetilde{\boldsymbol{B}}$ and $\widetilde{\boldsymbol{\beta}}$, namely, $\mathrm{e}^{\widetilde{\boldsymbol{B}}} \boldsymbol{e}_{j}=\mathbb{E}\left(\boldsymbol{Y}_{1} \mid \boldsymbol{Y}_{0}=\boldsymbol{e}_{j}\right), j \in\{1, \ldots, d\}$, and $\widetilde{\boldsymbol{\beta}}=\mathbb{E}\left(\boldsymbol{Z}_{1} \mid \boldsymbol{Z}_{0}=\mathbf{0}\right)$, where $\left(\boldsymbol{Y}_{t}\right)_{t \in \mathbb{R}_{+}}$and $\left(\boldsymbol{Z}_{t}\right)_{t \in \mathbb{R}_{+}}$are multi-type CBI processes with parameters $(d, \boldsymbol{c}, \mathbf{0}, \boldsymbol{B}, 0, \boldsymbol{\mu})$ and $(d, \mathbf{0}, \boldsymbol{\beta}, \mathbf{0}, \nu, \mathbf{0})$, respectively, see formula (2.4). The processes $\left(\boldsymbol{Y}_{t}\right)_{t \in \mathbb{R}_{+}}$and $\left(\boldsymbol{Z}_{t}\right)_{t \in \mathbb{R}_{+}}$can be considered as pure branching (without immigration) and pure immigration (without branching) processes, respectively. Consequently, $e^{\widetilde{B}}$ and $\widetilde{\boldsymbol{\beta}}$ may be called the branching mean matrix and the immigration mean vector, respectively. Note that the branching mechanism depends only on the parameters $\boldsymbol{c}, \boldsymbol{B}$ and $\boldsymbol{\mu}$, while the immigration mechanism depends only on the parameters $\boldsymbol{\beta}$ and $\nu$.
2.6 Lemma. Let $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$be a multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ such that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)<\infty$ and the moment condition (2.3) holds. Then the process ( $\mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-$ $\left.\int_{0}^{t} \mathrm{e}^{-u \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} u\right)_{t \in \mathbb{R}_{+}}$is a d-dimensional martingale with respect to the filtration $\mathcal{F}_{t}^{\boldsymbol{X}}:=\sigma\left(\boldsymbol{X}_{u}\right.$ : $u \in[0, t]), \quad t \in \mathbb{R}_{+}$.

Proof. First, note that for all $t \in \mathbb{R}_{+}, \boldsymbol{X}_{t}$ is measurable with respect to $\mathcal{F}_{t}^{\boldsymbol{X}}$, and due to $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)<\infty$ and (2.3), by Lemma 3.4 in Barczy et al. [5], we have $\mathbb{E}\left(\left\|\boldsymbol{X}_{t}\right\|\right)<\infty$. For each $v, t \in \mathbb{R}_{+}$with $v \leqslant t$, we have

$$
\mathbb{E}\left(\boldsymbol{X}_{t} \mid \mathcal{F}_{v}^{\boldsymbol{X}}\right)=\mathbb{E}\left(\boldsymbol{X}_{t} \mid \boldsymbol{X}_{v}\right)=\mathrm{e}^{(t-v) \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{v}+\int_{0}^{t-v} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} w,
$$

since $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is a time-homogeneous Markov process, and we can apply (2.4). Thus for each $v, t \in \mathbb{R}_{+}$with $v \leqslant t$, we obtain

$$
\begin{aligned}
& \mathbb{E}\left(\mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-\int_{0}^{t} \mathrm{e}^{-u \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} u \mid \mathcal{F}_{v}^{\boldsymbol{X}}\right)=\mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \mathrm{e}^{(t-v) \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{v}+\mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \int_{0}^{t-v} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} w-\int_{0}^{t} \mathrm{e}^{-u \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} u \\
& =\mathrm{e}^{-v \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{v}+\int_{0}^{t-v} \mathrm{e}^{(w-t) \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{d} w-\int_{0}^{t} \mathrm{e}^{-u \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} u=\mathrm{e}^{-v \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{v}-\int_{0}^{v} \mathrm{e}^{-u \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} u,
\end{aligned}
$$

and consequently, the process $\left(\mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-\int_{0}^{t} \mathrm{e}^{-u \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} u\right)_{t \in \mathbb{R}_{+}}$is a martingale with respect to the filtration $\left(\mathcal{F}_{t}^{\boldsymbol{X}}\right)_{t \in \mathbb{R}_{+}}$.

By an application of the multidimensional Itô's formula one can derive the following useful representation of $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$, where the drift part is deterministic. The proof can be found in Barczy et al. [6, Lemma 4.1].
2.7 Lemma. Let $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$be a multi-type CBI process with parameters (d, $\left.\boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu}\right)$ such that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)<\infty$ and the moment condition (2.3) holds. Then, for each $s, t \in \mathbb{R}_{+}$ with $s \leqslant t$, we have

$$
\begin{aligned}
\boldsymbol{X}_{t}= & \mathrm{e}^{(t-s) \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{s}+\int_{s}^{t} \mathrm{e}^{(t-u) \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{d} u+\sum_{\ell=1}^{d} \int_{s}^{t} \mathrm{e}^{(t-u) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell} \sqrt{2 c_{\ell} X_{u, \ell}} \mathrm{~d} W_{u, \ell} \\
& +\sum_{\ell=1}^{d} \int_{s}^{t} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{(t-u) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{w \leqslant X_{u-, \ell\}}\right.} \widetilde{N}_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w)+\int_{s}^{t} \int_{\mathcal{U}_{d}} \mathrm{e}^{(t-u) \widetilde{\boldsymbol{B}}} \boldsymbol{r} \widetilde{M}(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r}),
\end{aligned}
$$

where $\widetilde{M}(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r}):=M(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r})-\mathrm{d} u \nu(\mathrm{~d} \boldsymbol{r})$.
Note that the formula for $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$in Lemma 2.7 is a generalization of the formula (3.1) in Xu [23], and the formula (1.5) in Li and Ma [16].

Next we recall a classification of multi-type CBI processes. For a matrix $\boldsymbol{A} \in \mathbb{R}^{d \times d}, \sigma(\boldsymbol{A})$ will denote the spectrum of $\boldsymbol{A}$, that is, the set of all $\lambda \in \mathbb{C}$ that are eigenvalues of $\boldsymbol{A}$. Then $r(\boldsymbol{A}):=\max _{\lambda \in \sigma(\boldsymbol{A})}|\lambda|$ is the spectral radius of $\boldsymbol{A}$. Moreover, we will use the notation

$$
s(\boldsymbol{A}):=\max _{\lambda \in \sigma(\boldsymbol{A})} \operatorname{Re}(\lambda)
$$

A matrix $\boldsymbol{A} \in \mathbb{R}^{d \times d}$ is called reducible if there exist a permutation matrix $\boldsymbol{P} \in \mathbb{R}^{d \times d}$ and an integer $r$ with $1 \leqslant r \leqslant d-1$ such that

$$
\boldsymbol{P}^{\top} \boldsymbol{A} \boldsymbol{P}=\left[\begin{array}{cc}
\boldsymbol{A}_{1} & \boldsymbol{A}_{2} \\
0 & \boldsymbol{A}_{3}
\end{array}\right]
$$

where $\boldsymbol{A}_{1} \in \mathbb{R}^{r \times r}, \quad \boldsymbol{A}_{3} \in \mathbb{R}^{(d-r) \times(d-r)}, \quad \boldsymbol{A}_{2} \in \mathbb{R}^{r \times(d-r)}$, and $\mathbf{0} \in \mathbb{R}^{(d-r) \times r}$ is a null matrix. A matrix $\boldsymbol{A} \in \mathbb{R}^{d \times d}$ is called irreducible if it is not reducible, see, e.g., Horn and Johnson [10, Definitions 6.2.21 and 6.2.22]. We do emphasize that no 1-by-1 matrix is reducible.
2.8 Definition. Let $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$be a multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ such that the moment condition (2.3) holds. Then $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is called irreducible if $\widetilde{\boldsymbol{B}}$ is irreducible.

Recall that if $\widetilde{\boldsymbol{B}} \in \mathbb{R}_{(+)}^{d \times d}$ is irreducible, then $\mathrm{e}^{t \widetilde{\boldsymbol{B}}} \in \mathbb{R}_{++}^{d \times d}$ for all $t \in \mathbb{R}_{++}$, and $s(\widetilde{\boldsymbol{B}})$ is an eigenvalue of $\widetilde{\boldsymbol{B}}$, the algebraic and geometric multiplicities of $s(\widetilde{\boldsymbol{B}})$ is 1 , and the real parts of the other eigenvalues of $\widetilde{\boldsymbol{B}}$ are less than $s(\widetilde{\boldsymbol{B}})$. Moreover, corresponding to the eigenvalue $s(\widetilde{\boldsymbol{B}})$ there exists a unique (right) eigenvector $\widetilde{\boldsymbol{u}} \in \mathbb{R}_{++}^{d}$ of $\widetilde{\boldsymbol{B}}$ such that the sum of its coordinates is 1 which is also the unique (right) eigenvector of $\mathrm{e}_{\widetilde{B}}^{\widetilde{B}}$, called the right Perron vector of $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$, corresponding to the eigenvalue $r\left(\mathrm{e}^{\widetilde{\boldsymbol{B}}}\right)=\mathrm{e}^{s(\widetilde{\boldsymbol{B}})}$ of $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$ such that the sum of its coordinates is 1. Further, there exists a unique left eigenvector $\boldsymbol{u} \in \mathbb{R}_{++}^{d}$ of $\widetilde{\boldsymbol{B}}$ corresponding to the eigenvalue $s(\widetilde{\boldsymbol{B}})$ with $\widetilde{\boldsymbol{u}}^{\top} \boldsymbol{u}=1$, which is also the unique (left) eigenvector of $\mathrm{e}^{\widetilde{B}}$, called the left Perron vector of $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$, corresponding to the eigenvalue $r\left(\mathrm{e}^{\widetilde{\boldsymbol{B}}}\right)=\mathrm{e}^{s(\widetilde{\boldsymbol{B}})}$ of $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$ such that $\widetilde{\boldsymbol{u}}^{\top} \boldsymbol{u}=1$. Moreover, we have

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \mathrm{e}^{t \widetilde{\boldsymbol{B}}} \rightarrow \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \in \mathbb{R}_{++}^{d \times d} \quad \text { as } t \rightarrow \infty
$$

and there exist $C_{1}, C_{2}, C_{3} \in \mathbb{R}_{++}$such that

$$
\begin{equation*}
\left\|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \mathrm{e}^{t \widetilde{\boldsymbol{B}}}-\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right\| \leqslant C_{1} \mathrm{e}^{-C_{2} t}, \quad\left\|\mathrm{e}^{t \widetilde{\boldsymbol{B}}}\right\| \leqslant C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) t}, \quad t \in \mathbb{R}_{+} \tag{2.5}
\end{equation*}
$$

These Frobenius and Perron type results can be found, e.g., in Barczy and Pap [7, Appendix A].
2.9 Definition. Let $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$be an irreducible multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ such that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)<\infty$ and the moment condition (2.3) holds. Then $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is called

$$
\begin{cases}\text { subcritical } & \text { if } s(\widetilde{\boldsymbol{B}})<0 \\ \text { critical } & \text { if } s(\widetilde{\boldsymbol{B}})=0 \\ \text { supercritical } & \text { if } s(\widetilde{\boldsymbol{B}})>0\end{cases}
$$

For motivations of Definitions 2.8 and 2.9, see Barczy and Pap [7, Section 3]. Here we only point out that our classification of multi-type CBI processes is based on the asymptotic behaviour of $\mathbb{E}\left(\boldsymbol{X}_{t}\right)$ as $t \rightarrow \infty$, and this asymptotics is available at the moment only under the assumption of irreducibility of $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$.

## 3 Main results

First we present almost sure and $L_{1}$-convergence results for supercritical and irreducible multitype CBI processes.
3.1 Theorem. Let $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$be a supercritical and irreducible multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ such that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)<\infty$ and the moment condition (2.3) holds. Then, there exists a non-negative random variable $w_{\boldsymbol{u}, \boldsymbol{X}_{0}}$ with $\mathbb{E}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right)<\infty$ such that

$$
\begin{equation*}
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle \xrightarrow{\text { a.s. }} w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \quad \text { as } t \rightarrow \infty . \tag{3.1}
\end{equation*}
$$

Moreover, for each $\lambda \in \sigma(\widetilde{\boldsymbol{B}})$ such that $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$ and the moment condition

$$
\begin{equation*}
\sum_{\ell=1}^{d} \int_{\mathcal{U}_{d}} g(\|\boldsymbol{z}\|) \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty \tag{3.2}
\end{equation*}
$$

with

$$
g(x):=\left\{\begin{array}{ll}
x^{\frac{s(\widetilde{\boldsymbol{B}})}{\operatorname{Re}(\lambda)}} & \text { if } \operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right), \\
x \log (x) & \text { if } \operatorname{Re}(\lambda)=s(\widetilde{\boldsymbol{B}}) \quad(\Longleftrightarrow \lambda=s(\widetilde{\boldsymbol{B}})),
\end{array} \quad x \in \mathbb{R}_{++}\right.
$$

holds, and for each left eigenvector $\boldsymbol{v} \in \mathbb{C}^{d}$ of $\widetilde{\boldsymbol{B}}$ corresponding to the eigenvalue $\lambda$, there exists a complex random variable $w_{\boldsymbol{v}, \boldsymbol{X}_{0}}$ with $\mathbb{E}\left(\left|w_{\boldsymbol{v}, \boldsymbol{X}_{0}}\right|\right)<\infty$ such that

$$
\begin{equation*}
\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle \rightarrow w_{\boldsymbol{v}, \boldsymbol{X}_{0}} \quad \text { as } t \rightarrow \infty \quad \text { in } L_{1} \quad \text { and almost surely, } \tag{3.3}
\end{equation*}
$$

and

$$
\begin{align*}
w_{\boldsymbol{v}, \boldsymbol{X}_{0}} \stackrel{\text { a.s. }}{=} & \left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+\frac{\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle}{\lambda}+\sum_{\ell=1}^{d}\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle \int_{0}^{\infty} \mathrm{e}^{-\lambda u} \sqrt{2 c_{\ell} X_{u, \ell}} \mathrm{~d} W_{u, \ell} \\
& +\sum_{\ell=1}^{d} \int_{0}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{z}\rangle \mathbb{1}_{\left\{w \leqslant X_{u-, \ell}\right\}} \widetilde{N}_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w)  \tag{3.4}\\
& +\int_{0}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{r}\rangle \widetilde{M}(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r}),
\end{align*}
$$

where the improper integrals are convergent in $L_{1}$ and almost surely. Especially, $\mathbb{E}\left(w_{\boldsymbol{v}, \boldsymbol{X}_{0}}\right)=$ $\left\langle\boldsymbol{v}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)+\lambda^{-1} \widetilde{\boldsymbol{\beta}}\right\rangle$. Particularly, if $\left\langle\boldsymbol{v}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)+\lambda^{-1} \widetilde{\boldsymbol{\beta}}\right\rangle \neq 0$, then $\mathbb{P}\left(w_{\boldsymbol{v}, \boldsymbol{X}_{0}}=0\right)<1$. Further, $w_{\boldsymbol{u}, \boldsymbol{X}} \stackrel{\text { a.s. }}{=} 0$ if and only if $\boldsymbol{X}_{0} \stackrel{\text { a.s. }}{=} \mathbf{0}$ and $\widetilde{\boldsymbol{\beta}}=\mathbf{0} \quad$ (equivalently, $\boldsymbol{X}_{0} \stackrel{\text { a.s. }}{=} \mathbf{0}, \boldsymbol{\beta}=\mathbf{0}$ and $\nu=0$ ).

If the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle \xrightarrow{\text { a.s. }} 0$ as $t \rightarrow \infty$, i.e., $\mathbb{P}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}=0\right)=1$.

If the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle$ does not converge in $L_{1}$ as $t \rightarrow \infty$, provided that $\mathbb{P}\left(\boldsymbol{X}_{0}=\mathbf{0}\right)<1$ or $\widetilde{\boldsymbol{\beta}} \neq \mathbf{0}$. If $\mathbb{P}\left(\boldsymbol{X}_{0}=\mathbf{0}\right)=1$ and $\widetilde{\boldsymbol{\beta}}=\mathbf{0}$, then $\mathbb{P}\left(\boldsymbol{X}_{t}=\mathbf{0}\right)=1$ for all $t \in \mathbb{R}_{+}$.

Note that the asymptotic behavior of the second moment $\mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right|^{2}\right)$ as $t \rightarrow \infty$ explains the role of the assumption $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$ in Theorem 3.1, see Proposition B.1.
Proof of Theorem [3.1. By Lemma [2.6, the process $\left(\mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-\int_{0}^{t} \mathrm{e}^{-u \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} u\right)_{t \in \mathbb{R}_{+}}$is a martingale with respect to the filtration $\left(\mathcal{F}_{t}^{\boldsymbol{X}}\right)_{t \in \mathbb{R}_{+}}$. Moreover, for each $t \in \mathbb{R}_{+}$, we have

$$
\begin{align*}
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle & =\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{u}^{\top} \boldsymbol{X}_{t}=\boldsymbol{u}^{\top} \mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t} \\
& =\boldsymbol{u}^{\top}\left(\mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-\int_{0}^{t} \mathrm{e}^{-v \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} v\right)+\boldsymbol{u}^{\top} \int_{0}^{t} \mathrm{e}^{-v \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} v  \tag{3.5}\\
& =\boldsymbol{u}^{\top}\left(\mathrm{e}^{-t \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-\int_{0}^{t} \mathrm{e}^{-v \widetilde{\boldsymbol{B}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} v)+\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v,}\right.
\end{align*}
$$

where the function $\mathbb{R}_{+} \ni t \mapsto\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v \in \mathbb{R}_{+} \quad$ is increasing, since $\boldsymbol{u} \in \mathbb{R}_{++}^{d}$ and $\widetilde{\boldsymbol{\beta}} \in \mathbb{R}_{+}^{d}$. Consequently, $\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right)_{t \in \mathbb{R}_{+}}$is a submartingale with respect to the filtration $\left(\mathcal{F}_{t}^{\boldsymbol{X}}\right)_{t \in \mathbb{R}_{+}}$. Due to Theorem 4.6 in Barczy et al. [5], $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$and hence $\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right)_{t \in \mathbb{R}_{+}}$ have càdlàg sample paths almost surely. Using again $\boldsymbol{u} \in \mathbb{R}_{++}^{d}$ and (3.5), we get

$$
\begin{aligned}
\mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right|\right) & =\mathbb{E}\left(\mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right)=\mathbb{E}\left(\left\langle\boldsymbol{u}, \boldsymbol{X}_{0}\right\rangle\right)+\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v \\
& \leqslant\|\boldsymbol{u}\| \mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)+\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{\infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v=\|\boldsymbol{u}\| \mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\widetilde{\boldsymbol{B}})}
\end{aligned}
$$

for all $t \in \mathbb{R}_{+}$, thus we conclude $\sup _{t \in \mathbb{R}_{+}} \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right|\right)<\infty$. Hence, by the submartingale convergence theorem, there exists a non-negative random variable $w_{\boldsymbol{u}, \boldsymbol{X}_{0}}$ with $\mathbb{E}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right)<\infty$ such that (3.1) holds.

If $\lambda \in \sigma(\widetilde{\boldsymbol{B}})$ such that $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$ and the moment condition (3.2) holds, and $\boldsymbol{v} \in \mathbb{C}^{d}$ is a left eigenvector of $\boldsymbol{B}$ corresponding to the eigenvalue $\lambda$, then first we show the $L_{1}$-convergence of $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$ towards the right hand side of (3.4) together with the $L_{1}$-convergence of the improper integrals in (3.4). Note that the condition $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$ yields $\operatorname{Re}(\lambda)>0$, so $\lambda \neq 0$. For each $t \in \mathbb{R}_{+}$, by Lemma 2.7, we have the representation

$$
\begin{equation*}
\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle=\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+Z_{t}^{(1)}+Z_{t}^{(2)}+Z_{t}^{(3)}+Z_{t}^{(4)}+Z_{t}^{(5)} \tag{3.6}
\end{equation*}
$$

with

$$
\begin{aligned}
& Z_{t}^{(1)}:=\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-\lambda u} \mathrm{~d} u \\
& Z_{t}^{(2)}:=\sum_{\ell=1}^{d}\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle \int_{0}^{t} \mathrm{e}^{-\lambda u} \sqrt{2 c_{\ell} X_{u, \ell}} \mathrm{~d} W_{u, \ell}, \\
& Z_{t}^{(3)}:=\sum_{\ell=1}^{d} \int_{0}^{t} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{z}\rangle \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathbb{1}_{\left\{w \leqslant X_{u-,\}}\right\}} \widetilde{N}_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w),
\end{aligned}
$$

$$
\begin{aligned}
Z_{t}^{(4)} & :=\sum_{\ell=1}^{d} \int_{0}^{t} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{z}\rangle \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathbb{1}_{\left\{w \leqslant X_{u-, \ell}\right\}} \widetilde{N}_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w), \\
Z_{t}^{(5)} & :=\int_{0}^{t} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{r}\rangle \widetilde{M}(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r}) .
\end{aligned}
$$

Hence the $L_{1}$-convergence of $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$ towards the right hand side of (3.4) together with the $L_{1}$-convergence of the improper integrals in (3.4) will follow from the convergences $D_{t}^{(j)} \xrightarrow{L_{1}} 0$ as $t \rightarrow \infty$ for every $j \in\{1,2,3,4,5\}$ with

$$
\begin{aligned}
D_{t}^{(1)} & :=\frac{\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle}{\lambda}-\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-\lambda u} \mathrm{~d} u, \\
D_{t}^{(2)} & :=\sum_{\ell=1}^{d}\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle \int_{t}^{\infty} \mathrm{e}^{-\lambda u} \sqrt{2 c_{\ell} X_{u, \ell}} \mathrm{~d} W_{u, \ell}, \\
D_{t}^{(3)} & :=\sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{z}\rangle \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathbb{1}_{\left\{w \leqslant X_{u-, \ell\}}\right.} \widetilde{N}_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w), \\
D_{t}^{(4)} & :=\sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{z}\rangle \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathbb{1}_{\left\{w \leqslant X_{u-, \ell\}}\right.} \widetilde{N}_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w), \\
D_{t}^{(5)} & :=\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{r}\rangle \widetilde{M}(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r})
\end{aligned}
$$

for $t \in \mathbb{R}_{+}$. We have

$$
\begin{equation*}
D_{t}^{(1)}=\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{t}^{\infty} \mathrm{e}^{-\lambda u} \mathrm{~d} u \rightarrow 0 \quad \text { as } t \rightarrow \infty \tag{3.7}
\end{equation*}
$$

Moreover, for each $t \in \mathbb{R}_{+}$, we have

$$
\mathbb{E}\left(\int_{t}^{\infty}\left|\mathrm{e}^{-\lambda u}\right|^{2} 2 c_{\ell} X_{u, \ell} \mathrm{~d} u\right)=2 c_{\ell} \int_{t}^{\infty} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u
$$

By formulae (2.4) and (2.5), for each $v \in \mathbb{R}_{+}$and $\ell \in\{1, \ldots, d\}$, we get

$$
\begin{align*}
\mathbb{E}\left(X_{v, \ell}\right) & =\mathbb{E}\left(\boldsymbol{e}_{\ell}^{\top} \boldsymbol{X}_{v}\right)=\mathbb{E}\left(\boldsymbol{e}_{\ell}^{\top} \mathrm{e}^{v \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{0}+\boldsymbol{e}_{\ell}^{\top} \int_{0}^{v} \mathrm{e}^{u \widetilde{\boldsymbol{B}} \widetilde{\boldsymbol{\beta}}} \mathrm{~d} u\right) \\
& \leqslant\left\|\mathrm{e}^{v \widetilde{\boldsymbol{B}}}\right\| \mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)+\|\widetilde{\boldsymbol{\beta}}\| \int_{0}^{v}\left\|\mathrm{e}^{u \widetilde{\boldsymbol{B}}}\right\| \mathrm{d} u \\
& \leqslant C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v} \mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)+C_{3}\|\widetilde{\boldsymbol{\beta}}\| \int_{0}^{v} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) u} \mathrm{~d} u  \tag{3.8}\\
& =C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v} \mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)+C_{3}\|\widetilde{\boldsymbol{\beta}}\| \frac{\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}-1}{s(\widetilde{\boldsymbol{B}})} \leqslant C_{4} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}
\end{align*}
$$

with $C_{4}:=C_{3} \mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)+\frac{C_{3}\|\widetilde{\|}\|}{s(\tilde{\boldsymbol{B}})}$. By ((3.8), for each $t \in \mathbb{R}_{+}$and $\ell \in\{1, \ldots, d\}$, we obtain

$$
\begin{aligned}
& \mathbb{E}\left(\int_{t}^{\infty}\left|\mathrm{e}^{-\lambda u}\right|^{2} 2 c_{\ell} X_{u, \ell} \mathrm{~d} u\right) \leqslant 2 C_{4} c_{\ell} \int_{t}^{\infty} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) u} \mathrm{~d} u \\
& =2 C_{4} c_{\ell} \int_{t}^{\infty} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) u} \mathrm{~d} u=\frac{2 C_{4} c_{\ell}}{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) t}<\infty
\end{aligned}
$$

thus, by the independence of $\left(W_{t, 1}\right)_{t \in \mathbb{R}_{+}}, \ldots,\left(W_{t, d}\right)_{t \in \mathbb{R}_{+}}$and Itô's isometry for Itô's integrals (see, e.g., Ikeda and Watanabe [11, Chapter II, Proposition 2.2]),

$$
\begin{aligned}
& \mathbb{E}\left(\left|D_{t}^{(2)}\right|^{2}\right)= \mathbb{E}\left(\sum_{\ell=1}^{d} \int_{t}^{\infty} \operatorname{Re}\left(\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle \mathrm{e}^{-\lambda u}\right) \sqrt{2 c_{\ell} X_{u, \ell}} \mathrm{~d} W_{u, \ell}\right)^{2} \\
&+\mathbb{E}\left(\sum_{\ell=1}^{d} \int_{t}^{\infty} \operatorname{Im}\left(\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle \mathrm{e}^{-\lambda u}\right) \sqrt{2 c_{\ell} X_{u, \ell}} \mathrm{~d} W_{u, \ell}\right)^{2} \\
&=\sum_{\ell=1}^{d} \int_{t}^{\infty}\left(\operatorname{Re}\left(\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle \mathrm{e}^{-\lambda u}\right)\right)^{2} 2 c_{\ell} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u+\sum_{\ell=1}^{d} \int_{t}^{\infty}\left(\operatorname{Im}\left(\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle \mathrm{e}^{-\lambda u}\right)\right)^{2} 2 c_{\ell} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u \\
&= \sum_{\ell=1}^{d} \int_{t}^{\infty}\left|\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle \mathrm{e}^{-\lambda u}\right|^{2} 2 c_{\ell} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u=2 \sum_{\ell=1}^{d}\left|\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle\right|^{2} c_{\ell} \int_{t}^{\infty}\left|\mathrm{e}^{-\lambda u}\right|^{2} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u \\
& \leqslant 2 C_{4}\|\boldsymbol{v}\|^{2} \sum_{\ell=1}^{d} c_{\ell} \int_{t}^{\infty} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) u} \mathrm{~d} u=\frac{2 C_{4}\|\boldsymbol{v}\|^{2}}{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})}\left(\sum_{\ell=1}^{d} c_{\ell}\right) \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) t} .
\end{aligned}
$$

Consequently, we have

$$
\begin{equation*}
D_{t}^{(2)} \xrightarrow{L_{2}} 0 \quad \text { as } t \rightarrow \infty \tag{3.9}
\end{equation*}
$$

hence we conclude

$$
\begin{equation*}
D_{t}^{(2)} \xrightarrow{L_{1}} 0 \quad \text { as } t \rightarrow \infty . \tag{3.10}
\end{equation*}
$$

By (3.8), for each $t \in \mathbb{R}_{+}$, we have

$$
\begin{aligned}
& \sum_{\ell=1}^{d} \mathbb{E}\left(\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}}\left|\mathrm{e}^{-\lambda u}\right|^{2}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathbb{1}_{\left\{w \leqslant X_{u, \ell\}}\right.} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
& =\sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\mathrm{Re}(\lambda) u}\right\}} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \leqslant C_{4}\|\boldsymbol{v}\|^{2} K_{t}^{(3)}
\end{aligned}
$$

with

$$
\begin{equation*}
K_{t}^{(3)}:=\sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\tilde{\boldsymbol{B}})) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\operatorname{Re}(\lambda) u\}}\right.} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \leqslant K_{0}^{(3)} \tag{3.11}
\end{equation*}
$$

We show that $K_{0}^{(3)}<\infty$. For each $\ell \in\{1, \ldots, d\}$, using Fubini's theorem, we obtain

$$
\begin{aligned}
& \int_{0}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\|<1\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\int_{0}^{\infty} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) u} \mathrm{~d} u \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\|<1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\frac{1}{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\|<1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty
\end{aligned}
$$

by Definition 2.2, and

$$
\begin{aligned}
& \int_{0}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{1 \leqslant\|\boldsymbol{z}\|<\mathrm{e} \operatorname{Re}(\lambda) u\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\int_{\mathcal{U}_{d}}\left(\int_{\frac{1}{\operatorname{Re}(\lambda)} \log (\|\boldsymbol{z}\|)}^{\infty} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) u} \mathrm{~d} u\right)\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\int_{\mathcal{U}_{d}} \frac{1}{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})}\|\boldsymbol{z}\|^{-\frac{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})}{\operatorname{Re}(\lambda)}}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\frac{1}{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{\frac{s(\widetilde{\boldsymbol{B}})}{\operatorname{Re}(\lambda)}} \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty
\end{aligned}
$$

by the moment condition (3.2) (in case of $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)$ ) and by Definition 2.2 (in case of $\operatorname{Re}(\lambda)=s(\widetilde{\boldsymbol{B}})$ or equivalently $\lambda=s(\widetilde{\boldsymbol{B}}))$. Thus we obtain $K_{0}^{(3)}<\infty$. Consequently, by page 63 in Ikeda and Watanabe [11, for each $t \in \mathbb{R}_{+}$, we conclude

$$
\begin{aligned}
\mathbb{E}\left(\left|D_{t}^{(3)}\right|^{2}\right) & =\sum_{\ell=1}^{d} \mathbb{E}\left(\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}}\left|\mathrm{e}^{-\lambda u}\right|^{2}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\mathrm{Re}(\lambda) u}\right\}} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
& \leqslant C_{4}\|\boldsymbol{v}\|^{2} K_{t}^{(3)}<\infty
\end{aligned}
$$

We have

$$
K_{t}^{(3)}=K_{0}^{(3)}-\sum_{\ell=1}^{d} \int_{0}^{t} \int_{\mathcal{U}_{d}} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\mathrm{Re}(\lambda) u}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z})
$$

yielding

$$
\begin{equation*}
K_{t}^{(3)} \rightarrow 0 \quad \text { as } t \rightarrow \infty \tag{3.12}
\end{equation*}
$$

thus $\mathbb{E}\left(\left|D_{t}^{(3)}\right|^{2}\right) \rightarrow 0$ as $t \rightarrow \infty$. This implies $D_{t}^{(3)} \xrightarrow{L_{2}} 0$ as $t \rightarrow \infty$, thus we conclude

$$
\begin{equation*}
D_{t}^{(3)} \xrightarrow{L_{1}} 0 \quad \text { as } t \rightarrow \infty \tag{3.13}
\end{equation*}
$$

Further, for each $t \in \mathbb{R}_{+}$, we get

$$
\begin{aligned}
D_{t}^{(4)}= & \sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{z}\rangle \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant \mathrm{e} \mathrm{ere}(\lambda) u\}} \mathbb{1}_{\left\{w \leqslant X_{u-, \ell\}}\right.} N_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w) \\
& -\sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{z}\rangle \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u}\right\}} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w .
\end{aligned}
$$

Indeed, for each $t \in \mathbb{R}_{+}$and $\ell \in\{1, \ldots, d\}$, we have

$$
\begin{aligned}
& \left|\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{z}\rangle \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right| \\
& \leqslant\|\boldsymbol{v}\| \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w \in \mathbb{R}_{+}
\end{aligned}
$$

almost surely, since, by Fubini's theorem, (3.8) and (3.2), we get

$$
\begin{aligned}
& \mathbb{E}\left(\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u}\right\}} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
& =\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\operatorname{Re}(\lambda) u\}}\right.} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \leqslant C_{4} \int_{\mathcal{U}_{d}}\left(\int_{t}^{\infty} \mathrm{e}^{(s(\widetilde{\boldsymbol{B}})-\operatorname{Re}(\lambda)) u} \mathbb{1}_{\left\{u \leqslant \frac{1}{\operatorname{Re}(\lambda)} \log (\|\boldsymbol{z}\|)\right\}} \mathrm{d} u\right)\|\boldsymbol{z}\| \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \begin{cases}=C_{4} \int_{\mathcal{U}_{d}}\left(\frac{\log (\|\boldsymbol{z}\|)}{s(\widetilde{\boldsymbol{B}})}-t\right) \mathbb{1}_{\left\{\frac{\log (\|\boldsymbol{z}\|)}{s(\boldsymbol{B})} \geqslant t\right\}}\|\boldsymbol{z}\| \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
\leqslant \frac{C_{4}}{s(\widetilde{\boldsymbol{B}})} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\| \log (\|\boldsymbol{z}\|) \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\tilde{\boldsymbol{B}}) t}\right\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty & \text { if } \operatorname{Re}(\lambda)=s(\widetilde{\boldsymbol{B}}), \\
\text { and } \\
=C_{4} \int_{\mathcal{U}_{d}} \frac{\|\boldsymbol{z}\|}{\frac{s(\widetilde{\boldsymbol{B}})-\operatorname{Re}(\lambda)}{\operatorname{Re}(\lambda)}-\mathrm{e}^{(s(\widetilde{\boldsymbol{B}})-\operatorname{Re}(\lambda)) t}}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\frac{\log (\|\boldsymbol{B}\|)}{\operatorname{Re}(\lambda)} \geqslant t\right\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) & \\
\leqslant \frac{C_{4}}{s(\widetilde{\boldsymbol{B}})-\operatorname{Re}(\lambda)} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{\frac{s(\widetilde{B})}{\operatorname{Re}(\lambda)}} \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\operatorname{Re}(\lambda) t\}}\right\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty & \text { if } \operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)\end{cases}
\end{aligned}
$$

by the moment condition (3.2) and part (vi) of Definition 2.2. Consequently, we obtain

$$
\begin{align*}
\left|D_{t}^{(4)}\right| \leqslant & \|\boldsymbol{v}\| \sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u\}} \mathbb{1}_{\left\{w \leqslant X_{u-, \ell\}}\right.} N_{\ell}(\mathrm{d} u, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w)\right.}  \tag{3.14}\\
& +\|\boldsymbol{v}\| \sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u}\right\}} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w=:\|\boldsymbol{v}\| K_{t}^{(4)}
\end{align*}
$$

almost surely, where, by page 62 in Ikeda and Watanabe [11], for each $t \in \mathbb{R}_{+}$, we have

$$
\begin{aligned}
\mathbb{E}\left(K_{t}^{(4)}\right) & =2 \sum_{\ell=1}^{d} \mathbb{E}\left(\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\operatorname{Re}(\lambda) u}\right\}} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
& \leqslant \begin{cases}\frac{2 C_{4}}{s(\widetilde{\boldsymbol{B}})} \sum_{\ell=1}^{d} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\| \log (\|\boldsymbol{z}\|) \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) t}\right\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) & \text { if } \operatorname{Re}(\lambda)=s(\widetilde{\boldsymbol{B}}), \\
\frac{2 C_{4}}{s(\widetilde{\boldsymbol{B}})-\operatorname{Re}(\lambda)} \sum_{\ell=1}^{d} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{\frac{s(\widetilde{\boldsymbol{B}})}{\operatorname{Re}(\lambda)}} \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\operatorname{Re}(\lambda) t}\right\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) & \text { if } \\
\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right) .\end{cases}
\end{aligned}
$$

Thus the moment condition (3.2) yields that $\mathbb{E}\left(K_{t}^{(4)}\right)<\infty$ for all $t \in \mathbb{R}_{+}$and, by the dominated convergence theorem,

$$
\begin{equation*}
K_{t}^{(4)} \xrightarrow{L_{1}} 0 \quad \text { as } t \rightarrow \infty, \tag{3.15}
\end{equation*}
$$

and hence

$$
\begin{equation*}
D_{t}^{(4)} \xrightarrow{L_{1}} 0 \quad \text { as } t \rightarrow \infty . \tag{3.16}
\end{equation*}
$$

In a similar way, for each $t \in \mathbb{R}_{+}$, we get

$$
D_{t}^{(5)}=\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{r}\rangle M(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r})-\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{r}\rangle \mathrm{d} u \nu(\mathrm{~d} \boldsymbol{r}),
$$

since

$$
\begin{aligned}
\left|\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\lambda u}\langle\boldsymbol{v}, \boldsymbol{r}\rangle \mathrm{d} u \nu(\mathrm{~d} \boldsymbol{r})\right| & \leqslant\|\boldsymbol{v}\| \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{r}\| \mathrm{d} u \nu(\mathrm{~d} \boldsymbol{r}) \\
& =\frac{\|\boldsymbol{v}\|}{\operatorname{Re}(\lambda)} \mathrm{e}^{-\operatorname{Re}(\lambda) t} \int_{\mathcal{U}_{d}}\|\boldsymbol{r}\| \nu(\mathrm{d} \boldsymbol{r})<\infty
\end{aligned}
$$

by the moment condition (2.3) and by Definition 2.2. Consequently, we obtain

$$
\begin{align*}
\left|D_{t}^{(5)}\right| \leqslant & \|\boldsymbol{v}\| \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{r}\| M(\mathrm{~d} u, \mathrm{~d} \boldsymbol{r}) \\
& +\|\boldsymbol{v}\| \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{r}\| \mathrm{d} u \nu(\mathrm{~d} \boldsymbol{r})=:\|\boldsymbol{v}\| K_{t}^{(5)} \tag{3.17}
\end{align*}
$$

almost surely, where, by page 62 in Ikeda and Watanabe [11], for each $t \in \mathbb{R}_{+}$, we obtain

$$
\mathbb{E}\left(K_{t}^{(5)}\right)=\frac{2}{\operatorname{Re}(\lambda)} \mathrm{e}^{-\operatorname{Re}(\lambda) t} \int_{\mathcal{U}_{d}}\|\boldsymbol{r}\| \nu(\mathrm{d} \boldsymbol{r})<\infty .
$$

Hence we conclude

$$
\begin{equation*}
K_{t}^{(5)} \xrightarrow{L_{1}} 0 \quad \text { as } t \rightarrow \infty, \tag{3.18}
\end{equation*}
$$

implying

$$
\begin{equation*}
D_{t}^{(5)} \xrightarrow{L_{1}} 0 \quad \text { as } t \rightarrow \infty \tag{3.19}
\end{equation*}
$$

The convergences (3.7), (3.10), (3.13), (3.16) and (3.19) yield the $L_{1}$-convergence of $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ towards the right hand side of (3.4) as $t \rightarrow \infty$ together with the $L_{1}$-convergence of the improper integrals in (3.4). In fact, it turned out that $D_{t}^{(2)}$ and $D_{t}^{(3)}$ converge to 0 in $L_{2}$ as $t \rightarrow \infty$, but $D_{t}^{(4)}$ and $D_{t}^{(5)}$ converge to 0 only in $L_{1}$ as $t \rightarrow \infty$.

Next we show the almost sure convergence of $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$ together with the almost sure convergence of the improper integrals in (3.4). For each $t \in \mathbb{R}_{+}$, we use the representation (3.6). We have

$$
\begin{equation*}
Z_{t}^{(1)} \rightarrow\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{\infty} \mathrm{e}^{-\lambda u} \mathrm{~d} u=\frac{\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle}{\lambda} \quad \text { as } t \rightarrow \infty \tag{3.20}
\end{equation*}
$$

As in case of $\left(D_{t}^{(2)}\right)_{t \in \mathbb{R}_{+}}$, for each $t \in \mathbb{R}_{+}$, one can derive

$$
\begin{aligned}
\mathbb{E}\left(\left|Z_{t}^{(2)}\right|^{2}\right) & =2 \sum_{\ell=1}^{d}\left|\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle\right|^{2} c_{\ell} \int_{0}^{t}\left|\mathrm{e}^{-\lambda u}\right|^{2} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u \\
& \leqslant 2 C_{4}\|\boldsymbol{v}\|^{2} \sum_{\ell=1}^{d} c_{\ell} \int_{0}^{\infty} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) u} \mathrm{~d} u=\frac{2 C_{4}\|\boldsymbol{v}\|^{2}}{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})} \sum_{\ell=1}^{d} c_{\ell}<\infty
\end{aligned}
$$

hence the real and imaginary parts of $\left(Z_{t}^{(2)}\right)_{t \in \mathbb{R}_{+}}$are $L_{2}$-bounded martingales. As in case of $\left(D_{t}^{(3)}\right)_{t \in \mathbb{R}_{+}}$, for each $t \in \mathbb{R}_{+}$, one can derive

$$
\begin{aligned}
& \mathbb{E}\left(\left|Z_{t}^{(3)}\right|^{2}\right)=\sum_{\ell=1}^{d} \int_{0}^{t} \int_{\mathcal{U}_{d}} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\operatorname{Re}(\lambda) u\}}\right.} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \leqslant C_{4}\|\boldsymbol{v}\|^{2} \sum_{\ell=1}^{d} \int_{0}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\tilde{\boldsymbol{B}})) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z})=C_{4}\|\boldsymbol{v}\|^{2} K_{0}^{(3)}<\infty .
\end{aligned}
$$

Consequently, the real and imaginary parts of $\left(Z_{t}^{(3)}\right)_{t \in \mathbb{R}_{+}}$are $L_{2}$-bounded martingales. As in case of $\left(D_{t}^{(4)}\right)_{t \in \mathbb{R}_{+}}$, for each $t \in \mathbb{R}_{+}$, one can derive

$$
\mathbb{E}\left(\left|Z_{t}^{(4)}\right|\right) \leqslant 2\|\boldsymbol{v}\| \sum_{\ell=1}^{d} \int_{0}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{\mathrm{Re}(\lambda) u\}}\right.} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z})=\|\boldsymbol{v}\| \mathbb{E}\left(K_{0}^{(4)}\right)<\infty
$$

hence the real and imaginary parts of $\left(Z_{t}^{(4)}\right)_{t \in \mathbb{R}_{+}}$are $L_{1}$-bounded martingales. As in case of $\left(D_{t}^{(5)}\right)_{t \in \mathbb{R}_{+}}$, for each $t \in \mathbb{R}_{+}$, one can derive

$$
\mathbb{E}\left(\left|Z_{t}^{(5)}\right|\right) \leqslant 2\|\boldsymbol{v}\| \int_{0}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-\operatorname{Re}(\lambda) u}\|\boldsymbol{r}\| \mathrm{d} u \nu(\mathrm{~d} \boldsymbol{r})=\|\boldsymbol{v}\| \mathbb{E}\left(K_{0}^{(5)}\right)<\infty
$$

hence the real and imaginary parts of $\left(Z_{t}^{(5)}\right)_{t \in \mathbb{R}_{+}}$are $L_{1}$-bounded martingales. Consequently, by the martingale convergence theorem, we conclude that the real and imaginary parts of the martingales $\left(Z_{t}^{(j)}\right)_{t \in \mathbb{R}_{+}}, \quad j \in\{2,3,4,5\}$, are almost sure convergent as $t \rightarrow \infty$, hence, by
(3.20), we conclude the almost sure convergence of $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$ together with the almost sure convergence of the improper integrals in (3.4).

We have already showed the $L_{1}$-convergence of $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$ towards the right hand side of (3.4), so the almost sure convergence of $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$ yields the almost sure convergence in (3.3) as well.

By the convergence $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle \xrightarrow{L_{1}} w_{\boldsymbol{v}, \boldsymbol{X}_{0}}$ as $t \rightarrow \infty$ in (3.3), we obtain $\mathbb{E}\left(\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right) \rightarrow$ $\mathbb{E}\left(w_{\boldsymbol{v}, \boldsymbol{X}_{0}}\right)$ as $t \rightarrow \infty$. On the other hand, for each $t \in \mathbb{R}_{+}$, using the representation (3.6) and the martingale property of the processes $\left(Z_{t}^{(j)}\right)_{t \in \mathbb{R}_{+}}, \quad j \in\{2,3,4,5\}$, we have $\mathbb{E}\left(\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right)=\left\langle\boldsymbol{v}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+Z_{t}^{(1)}=\left\langle\boldsymbol{v}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-\lambda u} \mathrm{~d} u$, hence we obtain $\mathbb{E}\left(\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right) \rightarrow\left\langle\boldsymbol{v}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle}{\lambda}$ as $t \rightarrow \infty$. Consequently, $\mathbb{E}\left(w_{\boldsymbol{v}, \boldsymbol{X}_{0}}\right)=\left\langle\boldsymbol{v}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle}{\lambda}$. From here, we can see that if $\left\langle\boldsymbol{v}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle}{\lambda} \neq 0$, then $\mathbb{P}\left(w_{\boldsymbol{v}, \boldsymbol{X}_{0}}=0\right)<1$.

Next, we prove that $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \stackrel{\text { a.s. }}{=} 0$ if and only if $\boldsymbol{X}_{0} \stackrel{\text { a.s. }}{=} \mathbf{0}$ and $\widetilde{\boldsymbol{\beta}}=\mathbf{0}$. Since $w_{\boldsymbol{u}, \boldsymbol{X}_{0}}$ is nonnegative, we have $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \stackrel{\text { a.s. }}{=} 0$ if and only if $\mathbb{E}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right)=0$. Since $\boldsymbol{u} \in \mathbb{R}_{++}^{d}, \mathbb{P}\left(\boldsymbol{X}_{0} \in \mathbb{R}_{+}^{d}\right)=1$ and $\widetilde{\boldsymbol{\beta}} \in \mathbb{R}_{+}^{d}$, we have $\mathbb{E}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right)=\left\langle\boldsymbol{u}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\widetilde{\boldsymbol{B}})}=0$ if and only if $\mathbb{E}\left(\boldsymbol{X}_{0}\right)=\mathbf{0}$ and $\widetilde{\boldsymbol{\beta}}=\mathbf{0}$, yielding the assertion in question.

Next, we prove that if the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $\mathbb{P}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}=0\right)=1$. For each $t, T \in \mathbb{R}_{+}$, by Lemma A.1, we have $\boldsymbol{X}_{t+T} \stackrel{\mathcal{D}}{=} \boldsymbol{X}_{t}^{(1)}+\boldsymbol{X}_{t}^{(2, T)}$, where $\left(\boldsymbol{X}_{s}^{(1)}\right)_{s \in \mathbb{R}_{+}}$and $\left(\boldsymbol{X}_{s}^{(2, T)}\right)_{s \in \mathbb{R}_{+}}$are independent multi-type CBI processes with $\mathbb{P}\left(\boldsymbol{X}_{0}^{(1)}=\mathbf{0}\right)=1$, $\boldsymbol{X}_{0}^{(2, T)} \stackrel{\mathcal{D}}{=} \boldsymbol{X}_{T}, \quad$ and with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ and $(d, \boldsymbol{c}, \mathbf{0}, \boldsymbol{B}, 0, \boldsymbol{\mu})$, respectively. Taking the inner product with $\boldsymbol{u}$ and multiplying by $\mathrm{e}^{-s(\tilde{\boldsymbol{B}})(t+T)}$, we obtain

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t+T}\right\rangle \stackrel{\mathcal{D}}{=} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}^{(1)}\right\rangle\right)+\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}^{(2, T)}\right\rangle\right), \quad t, T \in \mathbb{R}_{+} .
$$

Letting $t \rightarrow \infty$, by (3.1), we obtain

$$
w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \stackrel{\mathcal{D}}{=} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) T} w_{\boldsymbol{u}, \mathbf{0}}^{(1)}+\mathrm{e}^{-s(\tilde{\boldsymbol{B}}) T} w_{\boldsymbol{u}, \boldsymbol{X}_{0}^{(2, T)}}^{(2, T)}, \quad T \in \mathbb{R}_{+},
$$

where $w_{\boldsymbol{u}, \mathbf{0}}^{(1)}:=\lim _{t \rightarrow \infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}^{(1)}\right\rangle \quad$ and $w_{\boldsymbol{u}, \boldsymbol{X}_{0}^{(2, T)}}^{(2, T)}:=\lim _{t \rightarrow \infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}^{(2, T)}\right\rangle \quad$ almost surely. Due to Kyprianou et al. [14, Theorem 1.3/(ii)] and the law of total probability, we have $\mathbb{P}\left(w_{u, \boldsymbol{X}_{0}^{(2, T)}}^{(2, T)}=0\right)=1$, since $\left(\boldsymbol{X}_{s}^{(2, T)}\right)_{s \in \mathbb{R}_{+}}$is a CB process and the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$. Consequently, for each $T \in \mathbb{R}_{+}$, we have $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \stackrel{\mathcal{D}}{=} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T} w_{\boldsymbol{u}, \mathbf{0}}^{(1)}$. The convergence $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T} w_{\boldsymbol{u}, \mathbf{0}}^{(1)} \xrightarrow{\text { a.s. }} 0$ as $T \rightarrow \infty$ yields $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T} w_{\boldsymbol{u}, \mathbf{0}}^{(1)} \xrightarrow{\mathcal{D}} 0$ as $T \rightarrow \infty$, thus we conclude that $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \stackrel{\text { a.s. }}{=} 0$, hence the proof is complete.

Finally, we prove that if the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle$ does not converge in $L_{1}$ as $t \rightarrow \infty$, provided that $\mathbb{P}\left(\boldsymbol{X}_{0}=\mathbf{0}\right)<1$ or $\widetilde{\boldsymbol{\beta}} \neq \mathbf{0}$. On the contrary, let us suppose that $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle$ converges in $L_{1}$ as $t \rightarrow \infty$. Recall that if the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle \xrightarrow{\text { a.s. }} 0$ as $t \rightarrow \infty$, which yields that $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle$ could converge only to 0 in $L_{1}$ as $t \rightarrow \infty$.

Especially, $\mathbb{E}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right)$ would converge to 0 as $t \rightarrow \infty$. Using (3.5), we have
$\mathbb{E}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right)=\left\langle\boldsymbol{u}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v \rightarrow\left\langle\boldsymbol{u}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\widetilde{\boldsymbol{B}})} \quad$ as $t \rightarrow \infty$,
so $\boldsymbol{u}^{\top} \mathbb{E}\left(\boldsymbol{X}_{0}\right)+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\boldsymbol{B})}=0$ should hold. Since $\boldsymbol{u} \in \mathbb{R}_{++}^{d}, \mathbb{P}\left(\boldsymbol{X}_{0} \in \mathbb{R}_{+}^{d}\right)=1$ and $\widetilde{\boldsymbol{\beta}} \in \mathbb{R}_{+}^{d}$, this would imply that $\boldsymbol{X}_{0} \stackrel{\text { a.s. }}{=} \mathbf{0}$ and $\widetilde{\boldsymbol{\beta}}=\mathbf{0}$ (equivalently, $\boldsymbol{X}_{0} \stackrel{\text { a.s. }}{=} \mathbf{0}, \boldsymbol{\beta}=\mathbf{0}$ and $\nu=0$ ) leading us to a contradiction.

In the next remark we explain why we do not have any result in the case when the moment condition (3.2) does not hold for $\lambda \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)$.
3.2 Remark. If the moment condition (3.2) does not hold for $\lambda \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)$, then we do not know whether $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ converges almost surely or not as $t \rightarrow \infty$, where $\boldsymbol{v} \in \mathbb{C}^{d}$ is a left eigenvector of $\widetilde{\boldsymbol{B}}$ corresponding to the eigenvalue $\lambda$. Provided that it converges almost surely to a complex random variable $w_{\boldsymbol{v}, \boldsymbol{X}_{0}}$, then, similarly as in the proof of Theorem 3.1, we would have

$$
w_{\boldsymbol{v}, \boldsymbol{X}_{0}} \stackrel{\mathcal{D}}{=} \mathrm{e}^{-\lambda T} w_{\boldsymbol{v}, \mathbf{0}}^{(1)}+\mathrm{e}^{-\lambda T} w_{\boldsymbol{v}, \boldsymbol{X}_{0}^{(2, T)}}^{(2, T)}, \quad T \in \mathbb{R}_{+}
$$

where $w_{\boldsymbol{v}, \mathbf{0}}^{(1)}:=\lim _{t \rightarrow \infty} \mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}^{(1)}\right\rangle$ and $w_{\boldsymbol{v}, \boldsymbol{X}_{0}^{(2, T)}}^{(2, T)}:=\lim _{t \rightarrow \infty} \mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}^{(2, T)}\right\rangle$ almost surely, and $\left(\boldsymbol{X}_{s}^{(1)}\right)_{s \in \mathbb{R}_{+}}$and $\left(\boldsymbol{X}_{s}^{(2, T)}\right)_{s \in \mathbb{R}_{+}}$are independent multi-type CBI processes with $\mathbb{P}\left(\boldsymbol{X}_{0}^{(1)}=\mathbf{0}\right)=1$, $\boldsymbol{X}_{0}^{(2, T)} \stackrel{\mathcal{D}}{=} \boldsymbol{X}_{T}$, and with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ and $(d, \boldsymbol{c}, \mathbf{0}, \boldsymbol{B}, 0, \boldsymbol{\mu})$, respectively. However, contrary to the proof of Theorem 3.1, we do not know whether $\mathbb{P}\left(w_{\boldsymbol{v}, \boldsymbol{X}}^{(2, T)}=0\right)=1$ holds or not. In the proof of Theorem 3.1 we used that the corresponding result for $\lambda=s(\widetilde{\boldsymbol{B}})$ holds for CB processes due to Kyprianou et al. [14, Theorem 1.3/(ii)] which is based on a socalled spine decomposition technique. Unfortunately, we do not know whether this technique could be adapted to the case of $\lambda \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)$ or not. We also do not know if the moment condition (3.2) does not hold for $\lambda \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right)$, then $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle$ converges in $L_{1}$ or not as $t \rightarrow \infty$. In the proof of Theorem 3.1, the corresponding $L_{1}$-convergence in the case of $\lambda=s(\widetilde{\boldsymbol{B}})$ is based on the almost sure convergence of $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle$ as $t \rightarrow \infty$. The above mentioned questions remain open problems.
3.3 Theorem. Let $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$be a supercritical and irreducible multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ such that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|\right)<\infty$ and the moment condition (2.3) holds. Then

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \rightarrow w_{u, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}} \quad \text { as } t \rightarrow \infty \quad \text { almost surely, }
$$

where $w_{\boldsymbol{u}, \boldsymbol{X}_{0}}$ is introduced in (3.3). If the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \stackrel{\text { a.s. }}{=} 0$. If the moment condition (3.2) holds for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \rightarrow$ $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}}$ as $t \rightarrow \infty$ in $L_{1}$ as well. If the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t}$ does not converge in $L_{1}$ as $t \rightarrow \infty$, provided that $\mathbb{P}\left(\boldsymbol{X}_{0}=\mathbf{0}\right)<1$ or $\widetilde{\boldsymbol{\beta}} \neq \mathbf{0}$. If $\mathbb{P}\left(\boldsymbol{X}_{0}=\mathbf{0}\right)=1$ and $\widetilde{\boldsymbol{\beta}}=\mathbf{0}$, then $\mathbb{P}\left(\boldsymbol{X}_{t}=\mathbf{0}\right)=1$ for all $t \in \mathbb{R}_{+}$.

Proof. First, let us suppose that the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$. Then for each $\ell \in\{1, \ldots, d\}$, by (3.1), we have

$$
\limsup _{t \rightarrow \infty} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t} \boldsymbol{e}_{\ell}^{\top} \boldsymbol{X}_{t} \leqslant \limsup _{t \rightarrow \infty} \frac{1}{u_{\ell}} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle=0
$$

almost surely, yielding that $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \rightarrow \mathbf{0}$ as $t \rightarrow \infty$ almost surely.
In what follows, let us suppose that the moment condition (3.2) holds for $\lambda=s(\widetilde{\boldsymbol{B}})$. For each $t, T \in \mathbb{R}_{+}$, put

$$
\begin{equation*}
\boldsymbol{\Delta}_{t, t+T}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \boldsymbol{X}_{t+T}-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \mathrm{e}^{T \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)}\left(\boldsymbol{X}_{t+T}-\mathrm{e}^{T \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}\right) . \tag{3.21}
\end{equation*}
$$

We are going to carry out the proof in several steps. As an initial step, we show that for each $T \in \mathbb{R}_{+}$, we have

$$
\begin{equation*}
\boldsymbol{\Delta}_{t, t+T} \xrightarrow{L_{1}} \mathbf{0} \quad \text { as } t \rightarrow \infty . \tag{3.22}
\end{equation*}
$$

Using (3.22) we prove the $L_{1}$ convergence of $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t}$ towards $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}}$ as $t \rightarrow \infty$ (see (3.34)). Then we show the almost sure convergence of $\boldsymbol{\Delta}_{t, t+T}$ and that of $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t}$ along lattice times (see (3.36) and (3.46)). Finally, we prove almost sure convergence of $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t}$ towards $w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}}$ as $t \rightarrow \infty$.

By Lemma 2.7, we obtain the representation

$$
\begin{equation*}
\Delta_{t, t+T}=\boldsymbol{J}_{t, t+T}^{(1)}+\boldsymbol{J}_{t, t+T}^{(2)}+\boldsymbol{J}_{t, t+T}^{(3)}+\boldsymbol{J}_{t, t+T}^{(4)}+\boldsymbol{J}_{t, t+T}^{(5)} \tag{3.23}
\end{equation*}
$$

for all $t, T \in \mathbb{R}_{+}$with

$$
\begin{aligned}
& \boldsymbol{J}_{t, t+T}^{(1)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{t+T} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{d} v, \\
& \boldsymbol{J}_{t, t+T}^{(2)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \int_{t}^{t+T} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell} \sqrt{2 c_{\ell} X_{v, \ell}} \mathrm{~d} W_{v, \ell}, \\
& \boldsymbol{J}_{t, t+T}^{(3)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\|\left\langle\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}\right.} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell}\right\}} \widetilde{N}_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w), \\
& \boldsymbol{J}_{t, t+T}^{(4)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell}\right\}} \widetilde{N}_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w), \\
& \boldsymbol{J}_{t, t+T}^{(5)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{r} \widetilde{M}(\mathrm{~d} v, \mathrm{~d} \boldsymbol{r})
\end{aligned}
$$

for all $t, T \in \mathbb{R}_{+}$. For each $t, T \in \mathbb{R}_{+}$, we obtain

$$
\boldsymbol{J}_{t, t+T}^{(1)}=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \int_{0}^{T} \mathrm{e}^{(T-u) \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{d} u
$$

hence, for each $T \in \mathbb{R}_{+}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{t, t+T}^{(1)} \rightarrow \mathbf{0} \quad \text { as } t \rightarrow \infty \tag{3.24}
\end{equation*}
$$

By (2.5) and (3.8), for each $t, T \in \mathbb{R}_{+}$and $\ell \in\{1, \ldots, d\}$, we obtain

$$
\begin{aligned}
& \mathbb{E}\left(\int_{t}^{t+T}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell}\right\|^{2} 2 c_{\ell} X_{v, \ell} \mathrm{~d} v\right)=2 c_{\ell} \int_{t}^{t+T}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell}\right\|^{2} \mathbb{E}\left(X_{v, \ell}\right) \mathrm{d} v \\
& \leqslant 2 C_{3}^{2} C_{4} c_{\ell} \int_{t}^{t+T} \mathrm{e}^{2 s(\widetilde{\boldsymbol{B}})(t+T-v)} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v \\
& \leqslant 2 C_{3}^{2} C_{4} c_{\ell} \mathrm{e}^{2 s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{\infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v=\frac{2 C_{3}^{2} C_{4} c_{\ell}}{s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})(t+2 T)}<\infty
\end{aligned}
$$

thus

$$
\begin{align*}
\mathbb{E}\left(\left\|\boldsymbol{J}_{t, t+T}^{(2)}\right\|^{2}\right) & =\mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \mathbb{E}\left(\int_{t}^{t+T}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell}\right\|^{2} 2 c_{\ell} X_{v, \ell} \mathrm{~d} v\right)  \tag{3.25}\\
& \leqslant \frac{2 C_{3}^{2} C_{4}}{s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} c_{\ell} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})(t+2 T)}=\frac{2 C_{3}^{2} C_{4}}{s(\widetilde{\boldsymbol{B}})}\left(\sum_{\ell=1}^{d} c_{\ell}\right) \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}
\end{align*}
$$

Consequently, for each $T \in \mathbb{R}_{+}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{t, t+T}^{(2)} \xrightarrow{L_{2}} \mathbf{0} \quad \text { as } \quad t \rightarrow \infty \tag{3.26}
\end{equation*}
$$

and hence

$$
\begin{equation*}
\boldsymbol{J}_{t, t+T}^{(2)} \xrightarrow{L_{1}} \mathbf{0} \quad \text { as } t \rightarrow \infty . \tag{3.27}
\end{equation*}
$$

By (2.5), (3.8) and (3.11), for each $t, T \in \mathbb{R}_{+}$and $\ell \in\{1, \ldots, d\}$, we get

$$
\begin{aligned}
& \mathbb{E}\left(\int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z}\right\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v, \ell\}}\right.} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
& =\int_{t}^{t+T} \int_{\mathcal{U}_{d}}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z}\right\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\tilde{\boldsymbol{B}}) v}\right\}} \mathbb{E}\left(X_{v, \ell} \mathrm{~d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z})\right. \\
& \leqslant C_{3}^{2} C_{4} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) 2(t+T-v)}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v\}}\right.} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \leqslant C_{3}^{2} C_{4} \mathrm{e}^{2 s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty,
\end{aligned}
$$

hence, by page 62 in Ikeda and Watanabe [11] and (3.11), we have

$$
\begin{align*}
& \mathbb{E}\left(\left\|\boldsymbol{J}_{t, t+T}^{(3)}\right\|^{2}\right) \\
& =\mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \mathbb{E}\left(\int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z}\right\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v, \ell}\right\}} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right)  \tag{3.28}\\
& \leqslant C_{3}^{2} C_{4} \sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z})=C_{3}^{2} C_{4} K_{t}^{(3)} .
\end{align*}
$$

Thus, by (3.12), we obtain $\boldsymbol{J}_{t, t+T}^{(3)} \xrightarrow{L_{2}} \mathbf{0}$ as $t \rightarrow \infty$ for each $T \in \mathbb{R}_{+}$. Consequently, for each $T \in \mathbb{R}_{+}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{t, t+T}^{(3)} \xrightarrow{L_{1}} \mathbf{0} \quad \text { as } \quad t \rightarrow \infty . \tag{3.29}
\end{equation*}
$$

Further, similarly as in case of $\left(D_{t}^{(4)}\right)_{t \in \mathbb{R}_{+}}$, for each $t, T \in \mathbb{R}_{+}$, we have

$$
\begin{aligned}
\boldsymbol{J}_{t, t+T}^{(4)}= & \mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\tilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell\}}\right.} N_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w) \\
& -\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v, \ell}\right\}} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w
\end{aligned}
$$

almost surely, thus for each $i \in\{1, \ldots, d\}$, by (2.5), we obtain

$$
\begin{aligned}
\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{J}_{t, t+T}^{(4)}\right| \leqslant & \mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\tilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell\}}\right.} N_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w) \\
& +\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \boldsymbol{e}_{i}^{T} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v, \ell\}}\right.} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w \\
\leqslant & C_{3} \sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell\}}\right.} N_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w) \\
& +C_{3} \sum_{\ell=1}^{d} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v, \ell\}}\right.} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w,
\end{aligned}
$$

hence, by (3.14), we get

$$
\begin{equation*}
\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{J}_{t, t+T}^{(4)}\right| \leqslant C_{3} K_{t}^{(4)} \tag{3.30}
\end{equation*}
$$

almost surely. Consequently, by (3.15), for each $T \in \mathbb{R}_{+}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{t, t+T}^{(4)} \xrightarrow{L_{1}} \mathbf{0} \quad \text { as } t \rightarrow \infty . \tag{3.31}
\end{equation*}
$$

In a similar way, for each $t, T \in \mathbb{R}_{+}$, we have

$$
\boldsymbol{J}_{t, t+T}^{(5)}=\mathrm{e}^{-s(\tilde{\boldsymbol{B}})(t+T)} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{r} M(\mathrm{~d} v, \mathrm{~d} \boldsymbol{r})-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{r} \mathrm{d} v \nu(\mathrm{~d} \boldsymbol{r})
$$

almost surely, thus for each $i \in\{1, \ldots, d\}$, by (2.5), we obtain

$$
\begin{aligned}
& \left|\boldsymbol{e}_{i}^{\top} \boldsymbol{J}_{t, t+T}^{(5)}\right| \\
& \leqslant \mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \boldsymbol{e}_{i}^{T} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{r} M(\mathrm{~d} v, \mathrm{~d} \boldsymbol{r})+\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \boldsymbol{e}_{i}^{T} \mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{r} \mathrm{d} v \nu(\mathrm{~d} \boldsymbol{r}) \\
& \leqslant C_{3} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{r}\| M(\mathrm{~d} v, \mathrm{~d} \boldsymbol{r})+C_{3} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{r}\| \mathrm{d} v \nu(\mathrm{~d} \boldsymbol{r}),
\end{aligned}
$$

thus, by (3.17),

$$
\begin{equation*}
\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{J}_{t, t+T}^{(5)}\right| \leqslant C_{3} K_{t}^{(5)} \tag{3.32}
\end{equation*}
$$

almost surely. Consequently, by (3.18), for each $T \in \mathbb{R}_{+}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{t, t+T}^{(5)} \xrightarrow{L_{1}} \mathbf{0} \quad \text { as } \quad t \rightarrow \infty . \tag{3.33}
\end{equation*}
$$

The convergences (3.24), (3.27), (3.29), (3.31) and (3.33) yield (3.22). In fact, it turned out that, for each $T \in \mathbb{R}_{+}$, we have $\boldsymbol{J}_{t, t+T}^{(2)} \xrightarrow{L_{2}} \mathbf{0}$ and $\boldsymbol{J}_{t, t+T}^{(3)} \xrightarrow{L_{2}} \mathbf{0}$ as $t \rightarrow \infty$, but only $\boldsymbol{J}_{t, t+T}^{(4)} \xrightarrow{L_{1}} \mathbf{0}$ and $\boldsymbol{J}_{t, t+T}^{(5)} \xrightarrow{L_{1}} \mathbf{0}$ as $t \rightarrow \infty$.

Next we prove

$$
\begin{equation*}
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \xrightarrow{L_{1}} w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}} \quad \text { as } t \rightarrow \infty \tag{3.34}
\end{equation*}
$$

by (3.22) and (3.3) with $\lambda=s(\widetilde{\boldsymbol{B}})$ and $\boldsymbol{v}=\boldsymbol{u}$. For each $t, T \in \mathbb{R}_{+}$and $i \in\{1, \ldots, d\}$, using (3.21) and the identity $\boldsymbol{I}_{d}=\sum_{j=1}^{d} \boldsymbol{e}_{j} \boldsymbol{e}_{j}^{\top}$, we have

$$
\begin{aligned}
& \mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t+T}=\boldsymbol{e}_{i}^{\top} \boldsymbol{\Delta}_{t, t+T}+\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{T \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t} \\
& =\boldsymbol{e}_{i}^{\top} \boldsymbol{\Delta}_{t, t+T}+\sum_{j=1}^{d}\left(\boldsymbol{e}_{i}^{\top} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T} \mathrm{e}^{T \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{j}\right) \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{j}^{\top} \boldsymbol{X}_{t} \\
& =\boldsymbol{e}_{i}^{\top} \boldsymbol{\Delta}_{t, t+T}+\sum_{j=1}^{d}\left(\boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \boldsymbol{e}_{j}\right) \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{j}^{\top} \boldsymbol{X}_{t}+\sum_{j=1}^{d}\left[\boldsymbol{e}_{i}^{\top}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T} \mathrm{e}^{T \widetilde{\boldsymbol{B}}}-\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right) \boldsymbol{e}_{j}\right] \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{j}^{\top} \boldsymbol{X}_{t} \\
& =\boldsymbol{e}_{i}^{\top} \boldsymbol{\Delta}_{t, t+T}+\left(\boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right) \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{u}^{\top} \boldsymbol{X}_{t}+\sum_{j=1}^{d}\left[\boldsymbol{e}_{i}^{\top}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T} \mathrm{e}^{T \widetilde{\boldsymbol{B}}}-\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right) \boldsymbol{e}_{j}\right] \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{j}^{\top} \boldsymbol{X}_{t} .
\end{aligned}
$$

By (2.5), for each $T \in \mathbb{R}_{+}$and $i, j \in\{1, \ldots, d\}$, we have

$$
\left|\boldsymbol{e}_{i}^{\top}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T} \mathrm{e}^{T \widetilde{\boldsymbol{B}}}-\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right) \boldsymbol{e}_{j}\right| \leqslant\left\|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) T} \mathrm{e}^{T \widetilde{\boldsymbol{B}}}-\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right\| \leqslant C_{1} \mathrm{e}^{-C_{2} T} \leqslant C_{5} \mathrm{e}^{-C_{2} T} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \boldsymbol{e}_{j},
$$

where

$$
C_{5}:=\max _{i, j \in\{1, \ldots, d\}} \frac{C_{1}}{\boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \boldsymbol{e}_{j}}=\frac{C_{1}}{\min _{i, j \in\{1, \ldots, d\}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \boldsymbol{e}_{j}}=\frac{C_{1}}{\min _{i \in\{1, \ldots, d\}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}} \min _{j \in\{1, \ldots, d\}} \boldsymbol{u}^{\top} \boldsymbol{e}_{j}} \in \mathbb{R}_{++},
$$

since $\boldsymbol{u}, \widetilde{\boldsymbol{u}} \in \mathbb{R}_{++}^{d}$. Hence for each $t, T \in \mathbb{R}_{+}$and $i \in\{1, \ldots, d\}$, we have

$$
\begin{align*}
& \left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t+T}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right| \\
& \leqslant\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{\Delta}_{t, t+T}\right|+\left|\left(\boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right)\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{u}^{\top} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right)\right| \\
& \quad+\sum_{j=1}^{d} C_{5} \mathrm{e}^{-C_{2} T}\left(\boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \boldsymbol{e}_{j}\right) \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{j}^{\top} \boldsymbol{X}_{t}  \tag{3.35}\\
& \leqslant\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{\Delta}_{t, t+T}\right|+\|\widetilde{\boldsymbol{u}}\|\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle-w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right|+C_{5}\|\widetilde{\boldsymbol{u}}\| \mathrm{e}^{-C_{2} T-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle .
\end{align*}
$$

For each $t, T \in \mathbb{R}_{+}$and $i \in\{1, \ldots, d\}$, by (3.35), we obtain

$$
\begin{aligned}
& \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t+T}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right|\right) \\
& \leqslant \mathbb{E}\left(\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{\Delta}_{t, t+T}\right|\right)+\|\widetilde{\boldsymbol{u}}\| \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle-w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right|\right)+C_{5}\|\widetilde{\boldsymbol{u}}\| \mathrm{e}^{-C_{2} T-s(\widetilde{\boldsymbol{B}}) t} \mathbb{E}\left(\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right) .
\end{aligned}
$$

By (3.22) and (3.3) with $\lambda=s(\widetilde{\boldsymbol{B}})$ and $\boldsymbol{v}=\boldsymbol{u}$, for each $T \in \mathbb{R}_{+}$and $i \in\{1, \ldots, d\}$, we obtain

$$
\begin{aligned}
& \limsup _{t \rightarrow \infty} \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right|\right)=\limsup _{t \rightarrow \infty} \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t+T}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right|\right) \\
& \leqslant C_{5}\|\widetilde{\boldsymbol{u}}\| \mathrm{e}^{-C_{2} T} \limsup _{t \rightarrow \infty} \mathbb{E}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right)=C_{5}\|\widetilde{\boldsymbol{u}}\| \mathrm{e}^{-C_{2} T} \mathbb{E}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right)
\end{aligned}
$$

hence, by $T \rightarrow \infty$, we conclude (3.34).
Next we show that for each $m \in \mathbb{N}$ and $\delta \in \mathbb{R}_{++}$, we have

$$
\begin{equation*}
\boldsymbol{\Delta}_{n \delta,(n+m) \delta} \xrightarrow{\text { a.s. }} \mathbf{0} \quad \text { as } n \rightarrow \infty . \tag{3.36}
\end{equation*}
$$

For each $m \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $\varepsilon \in \mathbb{R}_{++}$, by (3.25), we obtain

$$
\begin{align*}
\sum_{n=1}^{\infty} \mathbb{P}\left(\left\|\boldsymbol{J}_{n \delta,(n+m) \delta}^{(2)}\right\|>\varepsilon\right) & \leqslant \frac{1}{\varepsilon^{2}} \sum_{n=1}^{\infty} \mathbb{E}\left(\left\|\boldsymbol{J}_{n \delta,(n+m) \delta}^{(2)}\right\|^{2}\right) \\
& \leqslant \frac{2 C_{3}^{2} C_{4}}{\varepsilon^{2} s(\widetilde{\boldsymbol{B}})}\left(\sum_{\ell=1}^{d} c_{\ell}\right) \sum_{n=1}^{\infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta}<\infty \tag{3.37}
\end{align*}
$$

hence, by the Borel-Cantelli lemma, for each $m \in \mathbb{N}$ and $\delta \in \mathbb{R}_{++}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{n \delta,(n+m) \delta}^{(2)} \xrightarrow{\text { a.s. }} \mathbf{0} \quad \text { as } n \rightarrow \infty . \tag{3.38}
\end{equation*}
$$

For each $m \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $\varepsilon \in \mathbb{R}_{++}$, by (3.28), we obtain

$$
\begin{equation*}
\sum_{n=1}^{\infty} \mathbb{P}\left(\left\|\boldsymbol{J}_{n \delta,(n+m) \delta}^{(3)}\right\|>\varepsilon\right) \leqslant \frac{1}{\varepsilon^{2}} \sum_{n=1}^{\infty} \mathbb{E}\left(\left\|\boldsymbol{J}_{n \delta,(n+m) \delta}^{(3)}\right\|^{2}\right) \leqslant \frac{C_{3}^{2} C_{4}}{\varepsilon^{2}} \sum_{n=1}^{\infty} K_{n \delta}^{(3)} \tag{3.39}
\end{equation*}
$$

We will show

$$
\begin{equation*}
\sum_{n=1}^{\infty} K_{n \delta}^{(3)}<\infty . \tag{3.40}
\end{equation*}
$$

By Fubini's theorem, we have

$$
\begin{aligned}
\sum_{n=1}^{\infty} K_{n \delta}^{(3)} & =\sum_{n=1}^{\infty} \sum_{\ell=1}^{d} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\tilde{\boldsymbol{B}}) u\}}\right.} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\sum_{\ell=1}^{d} \sum_{n=1}^{\infty} \sum_{k=n}^{\infty} \int_{k \delta}^{(k+1) \delta} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) u}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z})
\end{aligned}
$$

$$
\begin{aligned}
& =\sum_{\ell=1}^{d} \sum_{k=1}^{\infty} k \int_{k \delta}^{(k+1) \delta} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\tilde{\boldsymbol{B}}) u\}}\right.} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \leqslant \frac{1}{\delta} \sum_{\ell=1}^{d} \int_{0}^{\infty} \int_{\mathcal{U}_{d}} u \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) u}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z})
\end{aligned}
$$

Here, for each $\ell \in\{1, \ldots, d\}$, using Fubini's theorem, we have

$$
\begin{aligned}
\int_{0}^{\infty} \int_{\mathcal{U}_{d}} u \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\|<1\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) & =\int_{0}^{\infty} u \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u} \mathrm{~d} u \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\|<1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\frac{1}{s(\widetilde{\boldsymbol{B}})^{2}} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\|<1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty
\end{aligned}
$$

by Definition [2.2, and

$$
\begin{aligned}
& \int_{0}^{\infty} \int_{\mathcal{U}_{d}} u \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\left\{1 \leqslant\|\boldsymbol{z}\|<\mathrm{e}^{s(\tilde{\boldsymbol{B}}) u\}}\right.} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\int_{\mathcal{U}_{d}}\left(\int_{\frac{1}{s(\widetilde{\boldsymbol{B}})} \log (\|\boldsymbol{z}\|)}^{\infty} u \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u} \mathrm{~d} u\right)\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\int_{\mathcal{U}_{d}} \frac{\|\boldsymbol{z}\|^{-1}}{s(\widetilde{\boldsymbol{B}})^{2}}(1+\log (\|\boldsymbol{z}\|))\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \leqslant \frac{1}{s(\widetilde{\boldsymbol{B}})^{2}} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\| \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})+\frac{1}{s(\widetilde{\boldsymbol{B}})^{2}} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\| \log (\|\boldsymbol{z}\|) \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty
\end{aligned}
$$

by the moment condition (3.2) with $\lambda=s(\widetilde{\boldsymbol{B}})$. Thus, for each $\delta \in \mathbb{R}_{++}$, we obtain (3.40). Hence, for each $m \in \mathbb{N}$ and $\delta \in \mathbb{R}_{++}$, by (3.39) and by the Borel-Cantelli lemma, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{n \delta,(n+m) \delta}^{(3)} \xrightarrow{\text { a.s. }} \mathbf{0} \quad \text { as } n \rightarrow \infty \tag{3.41}
\end{equation*}
$$

By (3.30), for each $m, n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $i \in\{1, \ldots, d\}$, we have $\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{J}_{n \delta,(n+m) \delta}^{(4)}\right| \leqslant C_{3} K_{n \delta}^{(4)}$. For each $\delta \in \mathbb{R}_{++}$, the function $\mathbb{N} \ni n \mapsto K_{n \delta}^{(4)}$ is decreasing almost surely, hence, by (3.15), we obtain

$$
\begin{equation*}
K_{n \delta}^{(4)} \xrightarrow{\text { a.s. }} 0 \quad \text { as } n \rightarrow \infty . \tag{3.42}
\end{equation*}
$$

Consequently, for each $m \in \mathbb{N}$ and $\delta \in \mathbb{R}_{++}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{n \delta,(n+m) \delta}^{(4)} \xrightarrow{\text { a.s. }} \mathbf{0} \quad \text { as } n \rightarrow \infty \tag{3.43}
\end{equation*}
$$

By (3.32), for each $m, n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $i \in\{1, \ldots, d\}$, we have $\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{J}_{n \delta,(n+m) \delta}^{(5)}\right| \leqslant C_{3} K_{n \delta}^{(5)}$. For each $\delta \in \mathbb{R}_{++}$, the function $\mathbb{N} \ni n \mapsto K_{n \delta}^{(5)}$ is decreasing almost surely, hence, by (3.18), we obtain

$$
\begin{equation*}
K_{n \delta}^{(5)} \xrightarrow{\text { a.s. }} 0 \quad \text { as } n \rightarrow \infty . \tag{3.44}
\end{equation*}
$$

Consequently, for each $m \in \mathbb{N}$ and $\delta \in \mathbb{R}_{++}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{n \delta,(n+m) \delta}^{(5)} \xrightarrow{\text { a.s. }} \mathbf{0} \quad \text { as } n \rightarrow \infty . \tag{3.45}
\end{equation*}
$$

The representation (3.23) and the convergences (3.24), (3.38), (3.41), (3.43) and (3.45) yield (3.36).

Next, using the almost sure convergences (3.36) and (3.1), we will show the almost sure convergence of $\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$along lattice times, i.e., we will prove that for each $\delta \in \mathbb{R}_{++}$, we have

$$
\begin{equation*}
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta} \boldsymbol{X}_{n \delta} \xrightarrow{\text { a.s. }} w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}} \quad \text { as } n \rightarrow \infty . \tag{3.46}
\end{equation*}
$$

By (3.35), (3.36) and (3.1), for each $m \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $i \in\{1, \ldots, d\}$, we obtain

$$
\begin{aligned}
& \limsup _{n \rightarrow \infty}\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{n \delta}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right|=\underset{n \rightarrow \infty}{\limsup }\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(n+m) \delta} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{(n+m) \delta}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right| \\
& \leqslant C_{5}\|\widetilde{\boldsymbol{u}}\| \mathrm{e}^{-C_{2} m \delta} \limsup _{n \rightarrow \infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta}\left\langle\boldsymbol{u}, \boldsymbol{X}_{n \delta}\right\rangle=C_{5}\|\widetilde{\boldsymbol{u}}\| \mathrm{e}^{-C_{2} m \delta} w_{\boldsymbol{u}, \boldsymbol{X}_{0}}
\end{aligned}
$$

almost surely, hence, by $m \rightarrow \infty$, we conclude (3.46).
The aim of the following discussion is to derive

$$
\begin{equation*}
\mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \xrightarrow{\text { a.s. }} w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}} \quad \text { as } t \rightarrow \infty \tag{3.47}
\end{equation*}
$$

by the help of the almost sure convergence (3.46) of $\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta} \boldsymbol{X}_{n \delta}\right)_{n \in \mathbb{N}}$ for all $\delta \in \mathbb{R}_{++}$ together with the almost sure convergence (3.1) of $\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right)_{t \in \mathbb{R}_{+}}$. For each $t \in \mathbb{R}_{+}$, $i \in\{1, \ldots, d\}, n \in \mathbb{N}$ and $\delta \in \mathbb{R}_{++}$, we have

$$
\begin{aligned}
\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right| \leqslant & \left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t}-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-t) \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}\right| \\
& +\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-t) \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right|,
\end{aligned}
$$

hence for each $i \in\{1, \ldots, d\}$, we get

$$
\begin{aligned}
& \limsup _{t \rightarrow \infty}\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right| \\
& \leqslant \limsup \limsup _{\delta \downarrow 0} \sup _{n \rightarrow \infty}\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right| \leqslant I_{i}^{(1)}+I_{i}^{(2)}
\end{aligned}
$$

with

$$
\begin{aligned}
I_{i}^{(1)} & :=\underset{\delta \downarrow 0}{\limsup } \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t}-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-t) \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}\right| \\
& =\underset{\delta \downarrow 0}{\limsup } \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top}\left(\boldsymbol{I}_{d}-\mathrm{e}^{((n+1) \delta-t) \widetilde{\boldsymbol{B}}}\right) \boldsymbol{X}_{t}\right|, \\
I_{i}^{(2)} & :=\underset{\delta \downarrow 0}{\limsup } \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-t) \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right| .
\end{aligned}
$$

For each $i \in\{1, \ldots, d\}$, we have $I_{i}^{(1)} \leqslant I^{(1,1)} I^{(1,2)}$ with

$$
\begin{aligned}
I^{(1,1)} & :=\underset{\delta \downarrow 0}{\lim \sup } \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left\|\boldsymbol{I}_{d}-\mathrm{e}^{((n+1) \delta-t) \widetilde{\boldsymbol{B}}}\right\| \\
& =\underset{\delta \downarrow 0}{\lim \sup } \limsup _{n \rightarrow \infty} \sup _{v \in(0, \delta]}\left\|\boldsymbol{I}_{d}-\mathrm{e}^{v \widetilde{\boldsymbol{B}}}\right\|=\limsup _{\delta \downarrow 0} \sup _{v \in(0, \delta]}\left\|\boldsymbol{I}_{d}-\mathrm{e}^{v \widetilde{\boldsymbol{B}}}\right\|=0, \\
I^{(1,2)} & :=\limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\|\boldsymbol{X}_{t}\right\| .
\end{aligned}
$$

For each $\boldsymbol{x} \in \mathbb{R}_{+}^{d}$, we have

$$
\begin{equation*}
\|\boldsymbol{x}\|^{2}=\sum_{j=1}^{d}\left(\boldsymbol{e}_{j}^{\top} \boldsymbol{x}\right)^{2}=\sum_{j=1}^{d} \frac{\left[\left(\boldsymbol{e}_{j}^{\top} \boldsymbol{u}\right)\left(\boldsymbol{e}_{j}^{\top} \boldsymbol{x}\right)\right]^{2}}{\left(\boldsymbol{e}_{j}^{\top} \boldsymbol{u}\right)^{2}} \leqslant \sum_{j=1}^{d} \frac{\langle\boldsymbol{u}, \boldsymbol{x}\rangle^{2}}{\left(\boldsymbol{e}_{j}^{\top} \boldsymbol{u}\right)^{2}}=C_{6}\langle\boldsymbol{u}, \boldsymbol{x}\rangle^{2} \tag{3.48}
\end{equation*}
$$

with $C_{6}:=\sum_{j=1}^{d} \frac{1}{\left(e_{j}^{1} u\right)^{2}} \in \mathbb{R}_{++}$, since $\boldsymbol{u} \in \mathbb{R}_{++}^{d}$. Thus, by the almost sure convergence (3.1), we obtain

$$
I^{(1,2)} \leqslant \sqrt{C_{6}} \limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle=\sqrt{C_{6}} w_{\boldsymbol{u}, \boldsymbol{X}_{0}}<\infty
$$

almost surely, and hence, we conclude $I_{i}^{(1)}=0$ almost surely for each $i \in\{1, \ldots, d\}$.
In order to show $I_{i}^{(2)}=0$ almost surely for each $i \in\{1, \ldots, d\}$, by Lemma 2.7, we consider the representation

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-t) \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}=\sum_{j=0}^{5} J_{t, i}^{(n, \delta, j)}
$$

with

$$
\begin{aligned}
& J_{t, i}^{(n, \delta, 0)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{\delta \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{n \delta}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}, \\
& J_{t, i}^{(n, \delta, 1)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{n \delta}^{t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{d} v, \\
& J_{t, i}^{(n, \delta, 2)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \sum_{\ell=1}^{d} \int_{n \delta}^{t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell} \sqrt{2 c_{\ell} X_{v, \ell}} \mathrm{~d} W_{v, \ell}, \\
& J_{t, i}^{(n, \delta, 3)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \sum_{\ell=1}^{d} \int_{n \delta}^{t} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\tilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell\}}\right.} \widetilde{N}_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w), \\
& J_{t, i}^{(n, \delta, 4)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \sum_{\ell=1}^{d} \int_{n \delta}^{t} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s}(\widetilde{\boldsymbol{B}}) v\right\}} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell\}}\right.} \widetilde{N}_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w), \\
& J_{t, i}^{(n, \delta, 5)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{n \delta}^{t} \int_{\mathcal{U}_{d}} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{r} \widetilde{M}(\mathrm{~d} v, \mathrm{~d} \boldsymbol{r})
\end{aligned}
$$

for all $i \in\{1, \ldots, d\}, \quad n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $t \in[n \delta,(n+1) \delta)$. For each $i \in\{1, \ldots, d\}$, $n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $t \in[n \delta,(n+1) \delta)$, we have

$$
J_{t, i}^{(n, \delta, 0)} \leqslant J_{t, i}^{(n, \delta, 0,1)}+J_{t, i}^{(n, \delta, 0,2)}+J_{i}^{(n, \delta, 0,3)}
$$

with

$$
\begin{aligned}
J_{t, i}^{(n, \delta, 0,1)} & :=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{\delta \widetilde{\boldsymbol{B}}} \boldsymbol{X}_{n \delta}-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{n \delta}=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top}\left(\mathrm{e}^{\delta \widetilde{\boldsymbol{B}}}-\boldsymbol{I}_{d}\right) \boldsymbol{X}_{n \delta}, \\
J_{t, i}^{(n, \delta, 0,2)} & :=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{n \delta}-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{n \delta}=\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t-n \delta)}-1\right) \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{n \delta}, \\
J_{i}^{(n, \delta, 0,3)} & :=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{n \delta}-w_{u, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}} .
\end{aligned}
$$

For each $i \in\{1, \ldots, d\}$, we have

$$
\begin{aligned}
& \limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 0,1)}\right| \\
& \leqslant\left(\limsup _{\delta \downarrow 0}\left\|\mathrm{e}^{\delta \widetilde{\boldsymbol{B}}}-\boldsymbol{I}_{d}\right\|\right)\left(\limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)} \mathrm{e}^{-\delta(\widetilde{\boldsymbol{B}}) t}\left\|\boldsymbol{X}_{n \delta}\right\|\right),
\end{aligned}
$$

where $\lim \sup _{\delta \downarrow 0}\left\|e^{\delta \widetilde{B}}-\boldsymbol{I}_{d}\right\|=0$, and, by (3.48) and (3.1),

$$
\begin{aligned}
\limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t}\left\|\boldsymbol{X}_{n \delta}\right\| & \leqslant \sqrt{C_{6}} \limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta}\left\langle\boldsymbol{u}, \boldsymbol{X}_{n \delta}\right\rangle \\
& =\sqrt{C_{6}} w_{\boldsymbol{u}, \boldsymbol{X}_{0}}<\infty
\end{aligned}
$$

almost surely, hence we get $\lim \sup _{\delta \downarrow 0} \lim \sup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 0,1)}\right|=0 \quad$ almost surely. Moreover, for each $i \in\{1, \ldots, d\}$, we have

$$
\begin{aligned}
& \limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 0,2)}\right| \\
& \leqslant\left(\limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t-n \delta)}-1\right|\right)\left(\limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta}\left\|\boldsymbol{X}_{n \delta}\right\|\right),
\end{aligned}
$$

where

$$
\limsup \limsup _{\delta \downarrow 0} \sup _{n \rightarrow \infty}\left|\mathrm{e}^{-s(\tilde{\boldsymbol{B}})(t-n \delta,(n+1) \delta)}-1\right|=\limsup _{\delta \downarrow 0}\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) \delta}-1\right|=0,
$$

and, by (3.48) and (3.1),

$$
\underset{\delta \downarrow 0}{\limsup } \limsup _{n \rightarrow \infty} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) n \delta}\left\|\boldsymbol{X}_{n \delta}\right\| \leqslant \sqrt{C_{6}} \limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) n \delta}\left\langle\boldsymbol{u}, \boldsymbol{X}_{n \delta}\right\rangle=\sqrt{C_{6}} w_{\boldsymbol{u}, \boldsymbol{X}_{0}}<\infty
$$

almost surely, hence we obtain $\lim \sup _{\delta \downarrow 0} \lim \sup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 0,2)}\right|=0$ almost surely. Further, for each $i \in\{1, \ldots, d\}$, by (3.46), we get

$$
\underset{\delta \downarrow 0}{\lim \sup } \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{i}^{(n, \delta, 0,3)}\right|=\limsup \limsup _{\delta \downarrow 0}\left|J_{i \rightarrow \infty}^{(n, \delta, 0,3)}\right|=0 \quad \text { almost surely, }
$$

and we conclude

$$
\begin{equation*}
\underset{\delta \downarrow 0}{\lim \sup } \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 0)}\right|=0 \quad \text { almost surely. } \tag{3.49}
\end{equation*}
$$

Moreover, for all $i \in\{1, \ldots, d\}, n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $t \in[n \delta,(n+1) \delta)$, by (2.5),

$$
\begin{aligned}
\left|J_{t, i}^{(n, \delta, 1)}\right| & \leqslant \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{n \delta}^{t}\left|\boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}}\right| \mathrm{d} v \leqslant\|\widetilde{\boldsymbol{\beta}}\| \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{n \delta}^{t}\left\|\mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}}\right\| \mathrm{d} v \\
& \leqslant C_{3}\|\widetilde{\boldsymbol{\beta}}\| \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{n \delta}^{t} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})((n+1) \delta-v)} \mathrm{d} v \leqslant C_{3}\|\widetilde{\boldsymbol{\beta}}\| \mathrm{e}^{s(\widetilde{\boldsymbol{B}})((n+1) \delta-t)} \int_{n \delta}^{\infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v \\
& =\frac{C_{3}\|\widetilde{\boldsymbol{\beta}}\|}{s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})((n+1) \delta-t)} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta}=\frac{C_{3}\|\widetilde{\boldsymbol{\beta}}\|}{s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})(\delta-t)}
\end{aligned}
$$

thus

$$
\begin{align*}
\underset{\delta \downarrow 0}{\lim \sup } \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 1)}\right| & \leqslant \frac{C_{3}\|\widetilde{\boldsymbol{\beta}}\|}{s(\widetilde{\boldsymbol{B}})} \limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})(\delta-t)} \\
& =\frac{C_{3}\|\widetilde{\boldsymbol{\beta}}\|}{s(\widetilde{\boldsymbol{B}})} \limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(n-1) \delta}=0 . \tag{3.50}
\end{align*}
$$

Further, for each $i \in\{1, \ldots, d\}, \quad n \in \mathbb{N}, \quad \delta \in \mathbb{R}_{++}$and $t \in[n \delta,(n+1) \delta)$, we have $\left|J_{t, i}^{(n, \delta, 2)}\right| \leqslant\left|M_{t, i}^{(n, \delta, 2)}\right|$, where

$$
M_{t, i}^{(n, \delta, 2)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta} \sum_{\ell=1}^{d} \int_{n \delta}^{t} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell} \sqrt{2 c_{\ell} X_{v, \ell}} \mathrm{~d} W_{v, \ell}, \quad t \in \mathbb{R}_{+},
$$

is a square integrable martingale (which can be checked as in case of $\left(Z_{t}^{(2)}\right)_{t \in \mathbb{R}_{+}}$in the proof of Theorem (3.1). By the maximal inequality for submartingales and by (3.25), for each $i \in$ $\{1, \ldots, d\}, n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $\varepsilon \in \mathbb{R}_{++}$, we obtain

$$
\begin{aligned}
& \mathbb{P}\left(\sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 2)}\right|>\varepsilon\right) \leqslant \mathbb{P}\left(\sup _{t \in[n \delta,(n+1) \delta)}\left|M_{t, i}^{(n, \delta, 2)}\right|>\varepsilon\right) \leqslant \frac{1}{\varepsilon^{2}} \mathbb{E}\left[\left(M_{(n+1) \delta, i}^{(n, \delta, 2)}\right)^{2}\right] \\
& =\frac{1}{\varepsilon^{2}} \mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}}) n \delta} \sum_{\ell=1}^{d} \mathbb{E}\left(\int_{n \delta}^{(n+1) \delta}\left|\boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell}\right|^{2} 2 c_{\ell} X_{v, \ell} \mathrm{~d} v\right) \\
& =\frac{2}{\varepsilon^{2}} \mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}}) n \delta} \sum_{\ell=1}^{d} c_{\ell} \int_{n \delta}^{(n+1) \delta}\left|\boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell}\right|^{2} \mathbb{E}\left(X_{v, \ell}\right) \mathrm{d} v \\
& \leqslant \frac{2}{\varepsilon^{2}} \mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}}) n \delta} \sum_{\ell=1}^{d} c_{\ell} \int_{n \delta}^{(n+1) \delta}\left\|\mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{e}_{\ell}\right\|^{2} \mathbb{E}\left(X_{v, \ell}\right) \mathrm{d} v=\frac{1}{\varepsilon^{2}} \mathrm{e}^{2 s(\widetilde{\boldsymbol{B}}) \delta} \mathbb{E}\left(\left\|\boldsymbol{J}_{n \delta,(n+1) \delta}^{(2)}\right\|^{2}\right),
\end{aligned}
$$

hence, for each $i \in\{1, \ldots, d\}, \delta \in \mathbb{R}_{++}$and $\varepsilon \in \mathbb{R}_{++}$, we obtain

$$
\sum_{n=1}^{\infty} \mathbb{P}\left(\sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 2)}\right|>\varepsilon\right) \leqslant \frac{1}{\varepsilon^{2}} \mathrm{e}^{2 s(\tilde{\boldsymbol{B}}) \delta} \sum_{n=1}^{\infty} \mathbb{E}\left(\left\|\boldsymbol{J}_{n \delta,(n+1) \delta}^{(2)}\right\|^{2}\right)<\infty
$$

by (3.37). By the Borel-Cantelli lemma, for each $i \in\{1, \ldots, d\}$ and $\delta \in \mathbb{R}_{++}$, we conclude

$$
\sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 2)}\right| \xrightarrow{\text { a.s. }} 0 \quad \text { as } n \rightarrow \infty
$$

and hence

$$
\begin{equation*}
\limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 2)}\right|=0 \quad \text { almost surely. } \tag{3.51}
\end{equation*}
$$

In a similar way, for each $i \in\{1, \ldots, d\}, n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $t \in[n \delta,(n+1) \delta)$, we have $\left|J_{t, i}^{(n, \delta, 3)}\right| \leqslant\left|M_{t, i}^{(n, \delta, 3)}\right|$, where

$$
M_{t, i}^{(n, \delta, 3)}:=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) n \delta} \sum_{\ell=1}^{d} \int_{n \delta}^{t} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \boldsymbol{e}_{i}^{\top} \mathrm{e}^{((n+1) \delta-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z} \mathbb{1}_{\left\{\|\boldsymbol{z}\|<\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v\}}\right.} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell}\right\}} \widetilde{N}_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w)
$$

for $t \in \mathbb{R}_{+}$defines a square integrable martingale (which can be checked as in case of $\left(Z_{t}^{(3)}\right)_{t \in \mathbb{R}_{+}}$ in the proof of Theorem (3.1). By the maximal inequality for submartingales and by (3.28), (3.39) and (3.40), for each $i \in\{1, \ldots, d\}, \delta \in \mathbb{R}_{++}$and $\varepsilon \in \mathbb{R}_{++}$, we obtain

$$
\begin{aligned}
\sum_{n=1}^{\infty} \mathbb{P}\left(\sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 3)}\right|>\varepsilon\right) & \leqslant \sum_{n=1}^{\infty} \mathbb{P}\left(\sup _{t \in[n \delta,(n+1) \delta)}\left|M_{t, i}^{(n, \delta, 3)}\right|>\varepsilon\right) \leqslant \frac{1}{\varepsilon^{2}} \sum_{n=1}^{\infty} \mathbb{E}\left[\left(M_{(n+1) \delta, i}^{(n, \delta, 3)}\right)^{2}\right] \\
& \leqslant \frac{1}{\varepsilon^{2}} \mathrm{e}^{2 s(\widetilde{\boldsymbol{B}}) \delta} \sum_{n=1}^{\infty} \mathbb{E}\left(\left\|\boldsymbol{J}_{n \delta,(n+1) \delta}^{(3)}\right\|^{2}\right) \leqslant \frac{C_{3}^{2} C_{4}}{\varepsilon^{2}} \mathrm{e}^{2 s(\widetilde{\boldsymbol{B}}) \delta} \sum_{n=1}^{\infty} K_{n \delta}^{(3)}<\infty,
\end{aligned}
$$

thus, by the Borel-Cantelli lemma, for each $i \in\{1, \ldots, d\}$ and $\delta \in \mathbb{R}_{++}$, we conclude

$$
\sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 3)}\right| \xrightarrow{\text { a.s. }} 0 \quad \text { as } n \rightarrow \infty
$$

and hence

$$
\begin{equation*}
\underset{\delta \downarrow 0}{\lim \sup } \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 3)}\right|=0 \quad \text { almost surely. } \tag{3.52}
\end{equation*}
$$

Next, for each $i \in\{1, \ldots, d\}, n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $t \in[n \delta,(n+1) \delta)$, as in case of $\boldsymbol{J}_{t, t+T}^{(4)}$, we obtain

$$
\begin{aligned}
\left|J_{t, i}^{(n, \delta, 4)}\right| \leqslant & C_{3} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t} \sum_{\ell=1}^{d} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})((n+1) \delta-v)}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\tilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell\}}\right.} N_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w) \\
& +C_{3} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \sum_{\ell=1}^{d} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})((n+1) \delta-v)}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v, \ell\}}\right.} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w
\end{aligned}
$$

almost surely, hence, by (3.14),

$$
\sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 4)}\right| \leqslant C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) \delta} \sum_{\ell=1}^{d} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v-, \ell\}}\right.} N_{\ell}(\mathrm{d} v, \mathrm{~d} \boldsymbol{z}, \mathrm{~d} w)
$$

$$
+C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) \delta} \sum_{\ell=1}^{d} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{z}\| \mathbb{1}_{\left\{\|\boldsymbol{z}\| \geqslant \mathrm{e}^{s(\tilde{\boldsymbol{B}}) v}\right\}} \mathbb{1}_{\left\{w \leqslant X_{v, \ell\}}\right.} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w=C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) \delta} K_{n \delta}^{(4)}
$$

almost surely. By (3.42), for each $i \in\{1, \ldots, d\}$, we conclude

$$
\begin{equation*}
\limsup \limsup _{\delta \downarrow 0} \sup _{n \rightarrow \infty}\left|J_{t \in[n \delta,(n+1) \delta)}^{(n, \delta, 4)}\right| \leqslant \limsup \limsup _{\delta \downarrow 0} C_{3 \rightarrow \infty} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) \delta} K_{n \delta}^{(4)}=0 \quad \text { almost surely. } \tag{3.53}
\end{equation*}
$$

Finally, for each $i \in\{1, \ldots, d\}, n \in \mathbb{N}, \delta \in \mathbb{R}_{++}$and $t \in[n \delta,(n+1) \delta)$, as in case of $\boldsymbol{J}_{t, t+T}^{(5)}$, we obtain

$$
\begin{aligned}
\left|J_{t, i}^{(n, \delta, 5)}\right| \leqslant & C_{3} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})((n+1) \delta-v)}\|\boldsymbol{r}\| M(\mathrm{~d} v, \mathrm{~d} \boldsymbol{r}) \\
& +C_{3} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})((n+1) \delta-v)}\|\boldsymbol{r}\| \mathrm{d} v \nu(\mathrm{~d} \boldsymbol{r})
\end{aligned}
$$

almost surely, hence, by (3.17),

$$
\begin{aligned}
\sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 5)}\right| \leqslant & C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) \delta} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{r}\| M(\mathrm{~d} v, \mathrm{~d} \boldsymbol{r}) \\
& +C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) \delta} \int_{n \delta}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{r}\| \mathrm{d} v \nu(\mathrm{~d} \boldsymbol{z})=C_{3} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) \delta} K_{n \delta}^{(5)}
\end{aligned}
$$

almost surely. By (3.44), for each $i \in\{1, \ldots, d\}$, we conclude

$$
\begin{equation*}
\limsup _{\delta \downarrow 0} \limsup _{n \rightarrow \infty} \sup _{t \in[n \delta,(n+1) \delta)}\left|J_{t, i}^{(n, \delta, 5)}\right| \leqslant \limsup \limsup _{\delta \downarrow 0} C_{3 \rightarrow \infty} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) \delta} K_{n \delta}^{(5)}=0 \quad \text { almost surely. } \tag{3.54}
\end{equation*}
$$

The convergences (3.49), (3.50), (3.51), (3.52), (3.53) and (3.54) yield (3.47).
Finally, we check that if the moment condition (3.2) does not hold for $\lambda=s(\widetilde{\boldsymbol{B}})$, then $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t}$ does not converge in $L_{1}$ as $t \rightarrow \infty$, provided that $\mathbb{P}\left(\boldsymbol{X}_{0}=\mathbf{0}\right)<1$ or $\widetilde{\boldsymbol{\beta}} \neq \mathbf{0}$. On the contrary, let us suppose that $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t}$ converges in $L_{1}$ as $t \rightarrow \infty$. Then, especially, $\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle$ converges in $L_{1}$ as $t \rightarrow \infty$, which leads us to a contradiction by Theorem 3.1.

Next we present $L_{2}$-convergence results for supercritical and irreducible multi-type CBI processes.
3.4 Theorem. Let $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$be a supercritical and irreducible multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ such that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|^{2}\right)<\infty$ and the moment conditions

$$
\begin{equation*}
\sum_{\ell=1}^{d} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty, \quad \int_{\mathcal{U}_{d}}\|\boldsymbol{r}\|^{2} \mathbb{1}_{\{\|\boldsymbol{r}\| \geqslant 1\}} \nu(\mathrm{d} \boldsymbol{r})<\infty \tag{3.55}
\end{equation*}
$$

hold. Then for each $\lambda \in \sigma(\widetilde{\boldsymbol{B}})$ with $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$ and for each left eigenvector $\boldsymbol{v} \in \mathbb{C}^{d}$ of $\widetilde{\boldsymbol{B}}$ corresponding to the eigenvalue $\lambda$, we have

$$
\begin{equation*}
\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle \xrightarrow{L_{2}} w_{\boldsymbol{v}, \boldsymbol{X}_{0}} \quad \text { as } t \rightarrow \infty \tag{3.56}
\end{equation*}
$$

(especially, $\left.\mathbb{E}\left(\left|w_{\boldsymbol{v}, \boldsymbol{X}_{0}}\right|^{2}\right)<\infty\right)$, where $w_{\boldsymbol{v}, \boldsymbol{X}_{0}}$ is introduced in (3.3), and the improper integrals in (3.4) are convergent in $L_{2}$.

Moreover,

$$
\begin{equation*}
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{X}_{t} \xrightarrow{L_{2}} w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \widetilde{\boldsymbol{u}} \quad \text { as } t \rightarrow \infty . \tag{3.57}
\end{equation*}
$$

Proof. First, note that the moment conditions of Theorem 3.1 hold, so, especially, we have the representation (3.4), where the improper integrals are convergent almost surely and in $L_{1}$ as well. In order to show (3.56), we consider the representation

$$
w_{\boldsymbol{v}, \boldsymbol{X}_{0}}-\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle=D_{t}^{(1)}+D_{t}^{(2)}+D_{t}^{(3)}+D_{t}^{(4)}+D_{t}^{(5)}, \quad t \in \mathbb{R}_{+},
$$

where $D_{t}^{(j)}, j \in\{1,2,3,4,5\}, t \in \mathbb{R}_{+}$, are defined in the proof of Theorem 3.1. Recall that by (3.7) and (3.9), we have $D_{t}^{(1)} \rightarrow 0$ and $D_{t}^{(2)} \xrightarrow{L_{2}} 0$ as $t \rightarrow \infty$. By (3.8), for each $t \in \mathbb{R}_{+}$ and $\ell \in\{1, \ldots, d\}$, we have

$$
\begin{aligned}
& \mathbb{E}\left(\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}}\left|\mathrm{e}^{-\lambda u}\right|^{2}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
& =\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \leqslant C_{4}\|\boldsymbol{v}\|^{2} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) u}\|\boldsymbol{z}\|^{2} \mathrm{~d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =C_{4}\|\boldsymbol{v}\|^{2} \int_{t}^{\infty} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) u} \mathrm{~d} u \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\frac{C_{4}\|\boldsymbol{v}\|^{2}}{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) t} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty,
\end{aligned}
$$

since $\int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z})=\int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\|<1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})+\int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mathbb{1}_{\{\|\boldsymbol{z}\| \geqslant 1\}} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty$ by Definition 2.2 and by the moment condition (3.55). Hence, by page 63 in Ikeda and Watanabe [11], for each $t \in \mathbb{R}_{+}$, we get

$$
\begin{aligned}
\mathbb{E}\left(\left|D_{t}^{(3)}+D_{t}^{(4)}\right|^{2}\right) & =\sum_{\ell=1}^{d} \mathbb{E}\left(\int_{t}^{\infty} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}}\left|\mathrm{e}^{-\lambda u}\right|^{2}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mathbb{1}_{\left\{w \leqslant X_{u, \ell}\right\}} \mathrm{d} u \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
& \leqslant \frac{C_{4}\|\boldsymbol{v}\|^{2}}{2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) t} \sum_{\ell=1}^{d} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z}),
\end{aligned}
$$

yielding

$$
\begin{equation*}
D_{t}^{(3)}+D_{t}^{(4)} \xrightarrow{L_{2}} 0 \quad \text { as } t \rightarrow \infty . \tag{3.58}
\end{equation*}
$$

Moreover, again by page 63 in Ikeda and Watanabe [11], for each $t \in \mathbb{R}_{+}$, we have

$$
\mathbb{E}\left(\left|D_{t}^{(5)}\right|^{2}\right)=\int_{t}^{\infty} \int_{\mathcal{U}_{d}}\left|\mathrm{e}^{-\lambda u}\right|^{2}|\langle\boldsymbol{v}, \boldsymbol{r}\rangle|^{2} \mathrm{~d} u \nu(\mathrm{~d} \boldsymbol{r}) \leqslant \frac{\|\boldsymbol{v}\|^{2}}{2 \operatorname{Re}(\lambda)} \mathrm{e}^{-2 \operatorname{Re}(\lambda) t} \int_{\mathcal{U}_{d}}\|\boldsymbol{r}\|^{2} \nu(\mathrm{~d} \boldsymbol{r})<\infty
$$

since $\operatorname{Re}(\lambda)>0$ and $\int_{\mathcal{U}_{d}}\|\boldsymbol{r}\|^{2} \nu(\mathrm{~d} \boldsymbol{r})=\int_{\mathcal{U}_{d}}\|\boldsymbol{r}\|^{2} \mathbb{1}_{\{\|\boldsymbol{r}\|<1\}} \nu(\mathrm{d} \boldsymbol{r})+\int_{\mathcal{U}_{d}}\|\boldsymbol{r}\|^{2} \mathbb{1}_{\{\|\boldsymbol{r}\| \geqslant 1\}} \nu(\mathrm{d} \boldsymbol{r})<\infty$ by Definition 2.2 and by the moment condition (3.55), implying

$$
\begin{equation*}
D_{t}^{(5)} \xrightarrow{L_{2}} 0 \quad \text { as } t \rightarrow \infty \tag{3.59}
\end{equation*}
$$

The convergences (3.7), (3.9), (3.58) and (3.59) yield (3.56).
In order to show (3.57), using the moment condition (3.55), first we prove that for each $T \in \mathbb{R}_{+}$, we have

$$
\begin{equation*}
\boldsymbol{\Delta}_{t, t+T} \xrightarrow{L_{2}} \mathbf{0} \quad \text { as } t \rightarrow \infty \tag{3.60}
\end{equation*}
$$

where $\boldsymbol{\Delta}_{t, t+T}, \quad t, T \in \mathbb{R}_{+}$, are introduced in (3.21). We use the representation (3.23) for $\boldsymbol{\Delta}_{t, t+T}, t, T \in \mathbb{R}_{+}$. Recall that by (3.24) and (3.26), we have $\boldsymbol{J}_{t, t+T}^{(1)} \rightarrow \mathbf{0}$ and $\boldsymbol{J}_{t, t+T}^{(2)} \xrightarrow{L_{2}} \mathbf{0}$ as $t \rightarrow \infty$ for all $T \in \mathbb{R}_{+}$. By (2.5) and (3.8), for each $t, T \in \mathbb{R}_{+}$and $\ell \in\{1, \ldots, d\}$, we get

$$
\begin{aligned}
& \mathbb{E}\left(\int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z}\right\|^{2} \mathbb{1}_{\left\{w \leqslant X_{v, \ell}\right\}} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
& =\int_{t}^{t+T} \int_{\mathcal{U}_{d}}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z}\right\|^{2} \mathbb{E}\left(X_{v, \ell}\right) \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \leqslant C_{3}^{2} C_{4} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) 2(t+T-v)}\|\boldsymbol{z}\|^{2} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& \leqslant C_{3}^{2} C_{4} \mathrm{e}^{2 s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{\infty} \int_{\mathcal{U}_{d}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v}\|\boldsymbol{z}\|^{2} \mathrm{~d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =C_{3}^{2} C_{4} \mathrm{e}^{2 s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{\infty} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) v} \mathrm{~d} v \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \\
& =\frac{C_{3}^{2} C_{4}}{s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{s(\widetilde{\boldsymbol{B}})(t+2 T)} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z})<\infty
\end{aligned}
$$

hence, by page 62 in Ikeda and Watanabe [11], we have

$$
\begin{gathered}
\mathbb{E}\left(\left\|\boldsymbol{J}_{t, t+T}^{(3)}+\boldsymbol{J}_{t, t+T}^{(4)}\right\|^{2}\right)=\mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}})(t+T)} \sum_{\ell=1}^{d} \mathbb{E}\left(\int_{t}^{t+T} \int_{\mathcal{U}_{d}} \int_{\mathcal{U}_{1}}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{z}\right\|^{2} \mathbb{1}_{\left\{w \leqslant X_{v, \ell}\right\}} \mathrm{d} v \mu_{\ell}(\mathrm{d} \boldsymbol{z}) \mathrm{d} w\right) \\
\leqslant \frac{C_{3}^{2} C_{4}}{s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \sum_{\ell=1}^{d} \int_{\mathcal{U}_{d}}\|\boldsymbol{z}\|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z})
\end{gathered}
$$

Consequently, for each $T \in \mathbb{R}_{+}$, we conclude

$$
\begin{equation*}
\boldsymbol{J}_{t, t+T}^{(3)}+\boldsymbol{J}_{t, t+T}^{(4)} \xrightarrow{L_{2}} \mathbf{0} \quad \text { as } \quad t \rightarrow \infty \tag{3.61}
\end{equation*}
$$

Moreover, by page 62 in Ikeda and Watanabe [11] and (2.5), for each $t, T \in \mathbb{R}_{+}$, we have

$$
\begin{aligned}
& \mathbb{E}\left(\left\|\boldsymbol{J}_{t, t+T}^{(5)}\right\|^{2}\right)=\mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{t+T} \int_{\mathcal{U}_{d}}\left\|\mathrm{e}^{(t+T-v) \widetilde{\boldsymbol{B}}} \boldsymbol{r}\right\|^{2} \mathrm{~d} v \nu(\mathrm{~d} \boldsymbol{r}) \\
& \leqslant C_{3}^{2} \mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}})(t+T)} \int_{t}^{t+T} \int_{\mathcal{U}_{d}} \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) 2(t+T-v)}\|\boldsymbol{r}\|^{2} \mathrm{~d} v \nu(\mathrm{~d} \boldsymbol{r}) \leqslant \frac{C_{3}^{2}}{2 s(\widetilde{\boldsymbol{B}})} \mathrm{e}^{-2 s(\widetilde{\boldsymbol{B}}) t} \int_{\mathcal{U}_{d}}\|\boldsymbol{r}\|^{2} \nu(\mathrm{~d} \boldsymbol{r}) .
\end{aligned}
$$

Consequently, for each $T \in \mathbb{R}_{+}$, we get

$$
\begin{equation*}
\boldsymbol{J}_{t, t+T}^{(5)} \xrightarrow{L_{2}} \mathbf{0} \quad \text { as } t \rightarrow \infty \tag{3.62}
\end{equation*}
$$

The convergences (3.24), (3.26), (3.61), and (3.62) yield (3.60).
Finally we prove (3.57). For each $t, T \in \mathbb{R}_{+}$and $i \in\{1, \ldots, d\}$, by (3.35), we obtain

$$
\begin{aligned}
& \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t+T}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right|^{2}\right) \\
& \leqslant 3 \mathbb{E}\left(\left|\boldsymbol{e}_{i}^{\top} \boldsymbol{\Delta}_{t, t+T}\right|^{2}\right)+3\|\widetilde{\boldsymbol{u}}\|^{2} \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle-w_{\boldsymbol{u}, \boldsymbol{X}_{0}}\right|^{2}\right)+3 C_{5}^{2}\|\widetilde{\boldsymbol{u}}\|^{2} \mathrm{e}^{-2 C_{2} T-2 s(\widetilde{\boldsymbol{B}}) t} \mathbb{E}\left(\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle^{2}\right)
\end{aligned}
$$

By (3.60) and (3.56) with $\lambda=s(\widetilde{\boldsymbol{B}})$ and $\boldsymbol{v}=\boldsymbol{u}$, for each $T \in \mathbb{R}_{+}$and $i \in\{1, \ldots, d\}$, we obtain

$$
\begin{aligned}
& \limsup _{t \rightarrow \infty} \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right|^{2}\right)=\limsup _{t \rightarrow \infty} \mathbb{E}\left(\left|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}})(t+T)} \boldsymbol{e}_{i}^{\top} \boldsymbol{X}_{t+T}-w_{\boldsymbol{u}, \boldsymbol{X}_{0}} \boldsymbol{e}_{i}^{\top} \widetilde{\boldsymbol{u}}\right|^{2}\right) \\
& \leqslant 3 C_{5}^{2}\|\widetilde{\boldsymbol{u}}\|^{2} \mathrm{e}^{-2 C_{2} T} \limsup _{t \rightarrow \infty} \mathbb{E}\left(\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left\langle\boldsymbol{u}, \boldsymbol{X}_{t}\right\rangle\right)^{2}\right)=3 C_{5}^{2}\|\widetilde{\boldsymbol{u}}\|^{2} \mathrm{e}^{-2 C_{2} T} \mathbb{E}\left(w_{\boldsymbol{u}, \boldsymbol{X}_{0}}^{2}\right)
\end{aligned}
$$

hence, by $T \rightarrow \infty$, we conclude (3.57).

## Appendix

## A On a decomposition of CBI processes

A. 1 Lemma. If $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is a multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$, then for each $t, T \in \mathbb{R}_{+}$, we have $\boldsymbol{X}_{t+T} \stackrel{\mathcal{D}}{=} \boldsymbol{X}_{t}^{(1)}+\boldsymbol{X}_{t}^{(2, T)}$, where $\left(\boldsymbol{X}_{s}^{(1)}\right)_{s \in \mathbb{R}_{+}}$and $\left(\boldsymbol{X}_{s}^{(2, T)}\right)_{s \in \mathbb{R}_{+}}$ are independent multi-type CBI processes with $\mathbb{P}\left(\boldsymbol{X}_{0}^{(1)}=\mathbf{0}\right)=1, \quad \boldsymbol{X}_{0}^{(2, T)} \stackrel{\mathcal{D}}{=} \boldsymbol{X}_{T}$, and with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ and $(d, \boldsymbol{c}, \mathbf{0}, \boldsymbol{B}, 0, \boldsymbol{\mu})$, respectively.

Proof. It is known that $\boldsymbol{v}(r, \boldsymbol{v}(s, \boldsymbol{\lambda}))=\boldsymbol{v}(r+s, \boldsymbol{\lambda})$ for all $r, s \in \mathbb{R}_{+}$and $\boldsymbol{\lambda} \in \mathbb{R}_{+}^{d}$, see, e.g., Li [15, page 58]. By the independence of $\left(\boldsymbol{X}_{s}^{(1)}\right)_{s \in \mathbb{R}_{+}}$and $\left(\boldsymbol{X}_{s}^{(2, T)}\right)_{s \in \mathbb{R}_{+}}$, by (2.2) and by the
law of total probability, for each $t, T \in \mathbb{R}_{+}$and $\boldsymbol{\lambda} \in \mathbb{R}_{+}^{d}$, we have

$$
\begin{aligned}
& \mathbb{E}\left(\mathrm{e}^{-\left\langle\boldsymbol{\lambda}, \boldsymbol{X}_{t}^{(1)}+\boldsymbol{X}_{t}^{(2, T)}\right\rangle}\right)=\mathbb{E}\left(\mathrm{e}^{-\left\langle\boldsymbol{\lambda}, \boldsymbol{X}_{t}^{(1)}\right\rangle}\right) \mathbb{E}\left(\mathrm{e}^{-\left\langle\boldsymbol{\lambda}, \boldsymbol{X}_{t}^{(2, T)}\right\rangle}\right) \\
& =\exp \left\{-\langle\mathbf{0}, \boldsymbol{v}(t, \boldsymbol{\lambda})\rangle-\int_{0}^{t} \psi(\boldsymbol{v}(s, \boldsymbol{\lambda})) \mathrm{d} s\right\} \mathbb{E}\left(\mathrm{e}^{-\left\langle\boldsymbol{X}_{0}^{(2, T)}, \boldsymbol{v}(t, \boldsymbol{\lambda})\right\rangle}\right) \\
& =\exp \left\{-\int_{0}^{t} \psi(\boldsymbol{v}(s, \boldsymbol{\lambda})) \mathrm{d} s\right\} \mathbb{E}\left(\mathrm{e}^{-\left\langle\boldsymbol{X}_{T}, \boldsymbol{v}(t, \boldsymbol{\lambda})\right\rangle}\right) \\
& =\exp \left\{-\int_{0}^{t} \psi(\boldsymbol{v}(s, \boldsymbol{\lambda})) \mathrm{d} s\right\} \mathbb{E}\left(\exp \left\{-\left\langle\boldsymbol{X}_{0}, \boldsymbol{v}(T, \boldsymbol{v}(t, \boldsymbol{\lambda}))\right\rangle-\int_{0}^{T} \psi(\boldsymbol{v}(s, \boldsymbol{v}(t, \boldsymbol{\lambda}))) \mathrm{d} s\right\}\right) \\
& =\mathbb{E}\left(\exp \left\{-\left\langle\boldsymbol{X}_{0}, \boldsymbol{v}(t+T, \boldsymbol{\lambda})\right\rangle-\int_{0}^{t} \psi(\boldsymbol{v}(s, \boldsymbol{\lambda})) \mathrm{d} s-\int_{0}^{T} \psi(\boldsymbol{v}(s+t, \boldsymbol{\lambda})) \mathrm{d} s\right\}\right) \\
& =\mathbb{E}\left(\exp \left\{-\left\langle\boldsymbol{X}_{0}, \boldsymbol{v}(t+T, \boldsymbol{\lambda})\right\rangle-\int_{0}^{t} \psi(\boldsymbol{v}(s, \boldsymbol{\lambda})) \mathrm{d} s-\int_{t}^{t+T} \psi(\boldsymbol{v}(u, \boldsymbol{\lambda})) \mathrm{d} u\right\}\right) \\
& =\mathbb{E}\left(\exp \left\{-\left\langle\boldsymbol{X}_{0}, \boldsymbol{v}(t+T, \boldsymbol{\lambda})\right\rangle-\int_{0}^{t+T} \psi(\boldsymbol{v}(s, \boldsymbol{\lambda})) \mathrm{d} s\right\}\right)=\mathbb{E}\left(\mathrm{e}^{-\left\langle\boldsymbol{\lambda}, \boldsymbol{X}_{t+T}\right\rangle}\right),
\end{aligned}
$$

hence we obtain the assertion.

## B On the second moment of projections of multi-type CBI processes

An explicit formula for the second moment of the projection of a multi-type CBI process on the left eigenvectors of its branching mean matrix is presented together with its asymptotic behavior in the supercritical and irreducible case.
B. 1 Proposition. If $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is a multi-type CBI process with parameters $(d, \boldsymbol{c}, \boldsymbol{\beta}, \boldsymbol{B}, \nu, \boldsymbol{\mu})$ such that $\mathbb{E}\left(\left\|\boldsymbol{X}_{0}\right\|^{2}\right)<\infty$ and the moment condition (3.55) holds, then for each left eigenvector $\boldsymbol{v} \in \mathbb{C}^{d}$ of $\widetilde{\boldsymbol{B}}$ corresponding to an arbitrary eigenvalue $\lambda \in \sigma(\widetilde{\boldsymbol{B}})$, we have

$$
\mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right|^{2}\right)=E_{\boldsymbol{v}, \lambda}(t)+\sum_{\ell=1}^{d} C_{\boldsymbol{v}, \ell} I_{\lambda, \ell}(t)+I_{\lambda}(t) \int_{\mathcal{U}_{d}}|\langle\boldsymbol{v}, \boldsymbol{r}\rangle|^{2} \nu(\mathrm{~d} \boldsymbol{r}), \quad t \in \mathbb{R}_{+},
$$

with

$$
\begin{aligned}
E_{\boldsymbol{v}, \lambda}(t) & :=\mathbb{E}\left(\left|\mathrm{e}^{\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{\lambda(t-u)} \mathrm{d} u\right|^{2}\right), \\
I_{\lambda, \ell}(t) & :=\int_{0}^{t} \mathrm{e}^{2 \operatorname{Re}(\lambda)(t-u)} \mathbb{E}\left(X_{u, \ell}\right) \mathrm{d} u, \quad \ell \in\{1, \ldots, d\},
\end{aligned}
$$

$$
\begin{aligned}
I_{\lambda}(t) & :=\int_{0}^{t} \mathrm{e}^{2 \operatorname{Re}(\lambda)(t-u)} \mathrm{d} u \\
C_{\boldsymbol{v}, \ell} & :=2\left|\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle\right|^{2} c_{\ell}+\int_{\mathcal{U}_{d}}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z}), \quad \ell \in\{1, \ldots, d\} .
\end{aligned}
$$

If, in addition, $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is supercritical and irreducible, then we have

$$
\lim _{t \rightarrow \infty} h(t) \mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right|^{2}\right)=M_{2}
$$

where

$$
h(t):= \begin{cases}\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} & \text { if } \operatorname{Re}(\lambda) \in\left(-\infty, \frac{1}{2} s(\widetilde{\boldsymbol{B}})\right) \\ t^{-1} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} & \text { if } \operatorname{Re}(\lambda)=\frac{1}{2} s(\widetilde{\boldsymbol{B}}) \\ \mathrm{e}^{-2 \operatorname{Re}(\lambda) t} & \text { if } \operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]\end{cases}
$$

and

$$
M_{2}:= \begin{cases}\frac{1}{s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)}\left(\left\langle\boldsymbol{u}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\widetilde{\boldsymbol{B}})}\right) \sum_{\ell=1}^{d} C_{\boldsymbol{v}, \ell}\left\langle\boldsymbol{e}_{\ell}, \widetilde{\boldsymbol{u}}\right\rangle & \text { if } \operatorname{Re}(\lambda) \in\left(-\infty, \frac{1}{2} s(\widetilde{\boldsymbol{B}})\right), \\ \left(\left\langle\boldsymbol{u}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\widetilde{\boldsymbol{B}})}\right) \sum_{\ell=1}^{d} C_{\boldsymbol{v}, \ell}\left\langle\boldsymbol{e}_{\ell}, \widetilde{\boldsymbol{u}}\right\rangle & \text { if } \operatorname{Re}(\lambda)=\frac{1}{2} s(\widetilde{\boldsymbol{B}}), \\ \mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+\frac{\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle}{\lambda}\right|^{2}\right)+\frac{1}{2 \operatorname{Re}(\lambda)} \int_{\mathcal{U}_{d}}|\langle\boldsymbol{v}, \boldsymbol{r}\rangle|^{2} \nu(\mathrm{~d} \boldsymbol{r}) \\ +\sum_{\ell=1}^{d} C_{\boldsymbol{v}, \ell} \boldsymbol{e}_{\ell}^{\top}\left(2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}-\widetilde{\boldsymbol{B}}\right)^{-1}\left(\mathbb{E}\left(\boldsymbol{X}_{0}\right)+\frac{\widetilde{\boldsymbol{\beta}}}{2 \operatorname{Re}(\lambda)}\right) & \text { if } \operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right] .\end{cases}
$$

Note that Proposition B. 1 can be considered as a counterpart of Theorem 1 in Section 7 in Chapter V in Athreya and Ney 3].
Proof of Proposition B.1. We use the representation (3.6) for $\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle, t \in \mathbb{R}_{+}$, where $Z_{t}^{(1)}$ is deterministic for all $t \in \mathbb{R}_{+}$. By the independence of $\boldsymbol{X}_{0},\left(W_{t, 1}\right)_{t \geqslant 0}, \ldots,\left(W_{t, d}\right)_{t \geqslant 0}$, $N_{1}, \ldots, N_{d}$ and $M$, for each $t \in \mathbb{R}_{+}$, the random variables $\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle$ and $Z_{t}^{(j)}, j \in\{2,3,4,5\}$ are conditionally independent with respect to $\left(\boldsymbol{X}_{u}\right)_{u \in[0, t]}$, hence for each $t \in \mathbb{R}_{+}$,

$$
\mathbb{E}\left(\left|\mathrm{e}^{-\lambda t}\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right|^{2} \mid\left(\boldsymbol{X}_{u}\right)_{u \in[0, t]}\right)=\mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+Z_{t}^{(1)}\right|^{2} \mid\left(\boldsymbol{X}_{u}\right)_{u \in[0, t]}\right)+\sum_{j=2}^{5} \mathbb{E}\left(\left|Z_{t}^{(j)}\right|^{2} \mid\left(\boldsymbol{X}_{u}\right)_{u \in[0, t]}\right)
$$

almost surely. We have, for each $t \in \mathbb{R}_{+}$,

$$
\begin{aligned}
\mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+Z_{t}^{(1)}\right|^{2} \mid\left(\boldsymbol{X}_{u}\right)_{u \in[0, t]}\right) & =\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-\lambda u} \mathrm{~d} u\right|^{2} \\
\mathbb{E}\left(\left|Z_{t}^{(2)}\right|^{2} \mid\left(\boldsymbol{X}_{u}\right)_{u \in[0, t]}\right) & =2 \sum_{\ell=1}^{d}\left|\left\langle\boldsymbol{v}, \boldsymbol{e}_{\ell}\right\rangle\right|^{2} c_{\ell} \int_{0}^{t} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} X_{u, \ell} \mathrm{~d} u, \\
\mathbb{E}\left(\left|Z_{t}^{(3)}\right|^{2}+\left|Z_{t}^{(4)}\right|^{2} \mid\left(\boldsymbol{X}_{u}\right)_{u \in[0, t]}\right) & =\sum_{\ell=1}^{d} \int_{0}^{t} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} X_{u, \ell} \mathrm{~d} u \int_{\mathcal{U}_{d}}|\langle\boldsymbol{v}, \boldsymbol{z}\rangle|^{2} \mu_{\ell}(\mathrm{d} \boldsymbol{z}), \\
\mathbb{E}\left(\left|Z_{t}^{(5)}\right|^{2} \mid\left(\boldsymbol{X}_{u}\right)_{u \in[0, t]}\right) & =\int_{0}^{t} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} \mathrm{~d} u \int_{\mathcal{U}_{d}}|\langle\boldsymbol{v}, \boldsymbol{r}\rangle|^{2} \nu(\mathrm{~d} \boldsymbol{r})
\end{aligned}
$$

almost surely. Taking the expectation and multiplying by $\left|\mathrm{e}^{\lambda t}\right|^{2}=\mathrm{e}^{2 \operatorname{Re}(\lambda) t}, t \in \mathbb{R}_{+}$, we obtain the formula for $\mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{t}\right\rangle\right|^{2}\right), t \in \mathbb{R}_{+}$.

Assume now that, in addition, $\left(\boldsymbol{X}_{t}\right)_{t \in \mathbb{R}_{+}}$is supercritical and irreducible. For each $t \in \mathbb{R}_{+}$, we have

$$
E_{\boldsymbol{v}, \lambda}(t) \leqslant 2 \mathrm{e}^{2 \operatorname{Re}(\lambda) t} \mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle\right|^{2}\right)+2|\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle|^{2}\left(\int_{0}^{t} \mathrm{e}^{\operatorname{Re}(\lambda) w} \mathrm{~d} w\right)^{2}
$$

If $\operatorname{Re}(\lambda) \in(-\infty, 0)$, then we have

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} E_{\boldsymbol{v}, \lambda}(t) \leqslant 2 \mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t} \mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle\right|^{2}\right)+\frac{2|\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle|^{2}}{\operatorname{Re}(\lambda)^{2}} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\left(\mathrm{e}^{\operatorname{Re}(\lambda) t}-1\right)^{2} \rightarrow 0
$$

as $t \rightarrow \infty$. If $\operatorname{Re}(\lambda)=0$, then we have

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} E_{\boldsymbol{v}, \lambda}(t) \leqslant 2 \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle\right|^{2}\right)+2|\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle|^{2} t^{2} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \rightarrow 0 \quad \text { as } t \rightarrow \infty .
$$

If $\operatorname{Re}(\lambda) \in\left(0, \frac{1}{2} s(\widetilde{\boldsymbol{B}})\right)$, then we have

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} E_{\boldsymbol{v}, \lambda}(t) \leqslant 2 \mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t} \mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle\right|^{2}\right)+\frac{2|\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle|^{2}}{\operatorname{Re}(\lambda)^{2}} \mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t}\left(1-\mathrm{e}^{-\operatorname{Re}(\lambda) t}\right)^{2} \rightarrow 0
$$

as $t \rightarrow \infty$. If $\operatorname{Re}(\lambda)=\frac{1}{2} s(\widetilde{\boldsymbol{B}})$ then

$$
t^{-1} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} E_{\boldsymbol{v}, \lambda}(t) \leqslant 2 t^{-1} \mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle\right|^{2}\right)+\frac{2|\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle|^{2}}{\operatorname{Re}(\lambda)^{2}} t^{-1}\left(1-\mathrm{e}^{-\operatorname{Re}(\lambda) t}\right)^{2} \rightarrow 0
$$

as $t \rightarrow \infty$. If $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$, then, by the dominated convergence theorem, we have

$$
\begin{aligned}
\mathrm{e}^{-2 \operatorname{Re}(\lambda) t} E_{\boldsymbol{v}, \lambda}(t) & =\mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{t} \mathrm{e}^{-\lambda u} \mathrm{~d} u\right|^{2}\right) \\
& \rightarrow \mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle \int_{0}^{\infty} \mathrm{e}^{-\lambda u} \mathrm{~d} u\right|^{2}\right)=\mathbb{E}\left(\left|\left\langle\boldsymbol{v}, \boldsymbol{X}_{0}\right\rangle+\frac{\langle\boldsymbol{v}, \widetilde{\boldsymbol{\beta}}\rangle}{\lambda}\right|^{2}\right)
\end{aligned}
$$

as $t \rightarrow \infty$.
Moreover, for each $u \in \mathbb{R}_{+}$and $\ell \in\{1, \ldots, d\}$, by formula (2.4), we have

$$
\mathbb{E}\left(X_{u, \ell}\right)=\boldsymbol{e}_{\ell}^{\top} \mathrm{e}^{u \widetilde{\boldsymbol{B}}} \mathbb{E}\left(\boldsymbol{X}_{0}\right)+\int_{0}^{u} \boldsymbol{e}_{\ell}^{\top} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \widetilde{\boldsymbol{\beta}} \mathrm{~d} w
$$

thus, for each $t \in \mathbb{R}_{+}$and $\ell \in\{1, \ldots, d\}$, we get

$$
I_{\lambda, \ell}(t)=\boldsymbol{e}_{\ell}^{\top} \boldsymbol{A}_{\lambda, 1}(t) \mathbb{E}\left(\boldsymbol{X}_{0}\right)+\boldsymbol{e}_{\ell}^{\top} \boldsymbol{A}_{\lambda, 2}(t) \widetilde{\boldsymbol{\beta}}
$$

with

$$
\boldsymbol{A}_{\lambda, 1}(t):=\int_{0}^{t} \mathrm{e}^{2 \operatorname{Re}(\lambda)(t-u)} \mathrm{e}^{u \widetilde{\boldsymbol{B}}} \mathrm{~d} u, \quad \boldsymbol{A}_{\lambda, 2}(t):=\int_{0}^{t} \mathrm{e}^{2 \operatorname{Re}(\lambda)(t-u)}\left(\int_{0}^{u} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w\right) \mathrm{d} u .
$$

If $\operatorname{Re}(\lambda) \in\left(-\infty, \frac{1}{2} s(\widetilde{\boldsymbol{B}})\right)$, then we have

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{A}_{\lambda, 1}(t)=\mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t} \int_{0}^{t} \mathrm{e}^{u\left(\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}\right)} \mathrm{d} u \rightarrow \frac{1}{s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \quad \text { as } t \rightarrow \infty,
$$

since $\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d} \in \mathbb{R}_{(+)}^{d \times d}$ is irreducible, $s\left(\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}\right)=s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)>0$, and the left (right) Perron vectors of the matrices $\mathrm{e}^{\widetilde{\boldsymbol{B}}}$ and $\mathrm{e}^{\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}}$ coincide, see, e.g., the proof of formula (B.2) in Barczy et al. [7]. If $\operatorname{Re}(\lambda)=0$, then, by Fubini's theorem, we obtain

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{A}_{\lambda, 2}(t)=\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t}(t-w) \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \rightarrow \frac{1}{s(\widetilde{\boldsymbol{B}})^{2}} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \quad \text { as } t \rightarrow \infty
$$

since, by (2.5), we have

$$
\begin{aligned}
& \left\|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t}(t-w) \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w-\frac{1}{s(\widetilde{\boldsymbol{B}})^{2}} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right\| \\
& \leqslant\left\|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t}(t-w) \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) w}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) w} \mathrm{e}^{w \widetilde{\boldsymbol{B}}}-\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right) \mathrm{d} w\right\| \\
& +\left\|\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t}(t-w) \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) w} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \mathrm{d} w-\frac{1}{s(\widetilde{\boldsymbol{B}})^{2}} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right\| \\
& \leqslant C_{1} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t}(t-w) \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) w} \mathrm{e}^{-C_{2} w} \mathrm{~d} w+\left\|\frac{\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}}{s(\widetilde{\boldsymbol{B}})^{2}}\left(\mathrm{e}^{s(\widetilde{\boldsymbol{B}}) t}-s(\widetilde{\boldsymbol{B}}) t-1\right) \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}-\frac{1}{s(\widetilde{\boldsymbol{B}})^{2}} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right\| \\
& \leqslant C_{1} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t}(t-w) \mathrm{e}^{s(\widetilde{\boldsymbol{B}}) w} \mathrm{e}^{-\widetilde{C}_{2} w} \mathrm{~d} w+\left\|\frac{\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}}{s(\widetilde{\boldsymbol{B}})^{2}}(s(\widetilde{\boldsymbol{B}}) t+1) \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right\| \\
& =\frac{C_{1} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}}{\left(s(\widetilde{\boldsymbol{B}})-\widetilde{C}_{2}\right)^{2}}\left(\mathrm{e}^{\left(s(\widetilde{\boldsymbol{B}})-\widetilde{C}_{2}\right) t}-\left(s(\widetilde{\boldsymbol{B}})-\widetilde{C}_{2}\right) t-1\right)+\frac{\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}}{s(\widetilde{\boldsymbol{B}})^{2}}(s(\widetilde{\boldsymbol{B}}) t+1)\left\|\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\right\| \rightarrow 0
\end{aligned}
$$

as $t \rightarrow \infty$, where $\widetilde{C}_{2} \in\left(0, C_{2} \wedge s(\widetilde{\boldsymbol{B}})\right)$. Hence, if $\operatorname{Re}(\lambda)=0$, then

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} I_{\lambda, \ell}(t) \rightarrow \frac{1}{s(\widetilde{\boldsymbol{B}})} \boldsymbol{e}_{\ell}^{\top} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\left(\mathbb{E}\left(\boldsymbol{X}_{0}\right)+\frac{\widetilde{\boldsymbol{\beta}}}{s(\widetilde{\boldsymbol{B}})}\right)=\frac{\left\langle\boldsymbol{e}_{\ell}, \widetilde{\boldsymbol{u}}\right\rangle}{s(\widetilde{\boldsymbol{B}})}\left(\left\langle\boldsymbol{u}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\widetilde{\boldsymbol{B}})}\right)
$$

as $t \rightarrow \infty$. If $\operatorname{Re}(\lambda) \in\left(-\infty, \frac{1}{2} s(\widetilde{\boldsymbol{B}})\right) \backslash\{0\}$, then, by Fubini's theorem, we obtain

$$
\begin{aligned}
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{A}_{\lambda, 2}(t) & =\mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t} \int_{0}^{t} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u}\left(\int_{0}^{u} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w\right) \mathrm{d} u \\
& =\mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t} \int_{0}^{t}\left(\int_{w}^{t} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} \mathrm{~d} u\right) \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \\
& =\frac{1}{2 \operatorname{Re}(\lambda)} \mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t} \int_{0}^{t}\left(\mathrm{e}^{-2 \operatorname{Re}(\lambda) w}-\mathrm{e}^{-2 \operatorname{Re}(\lambda) t}\right) \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w
\end{aligned}
$$

$$
\begin{aligned}
& =\frac{1}{2 \operatorname{Re}(\lambda)}\left(\mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t} \int_{0}^{t} \mathrm{e}^{\left(\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}\right) w} \mathrm{~d} w-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w\right) \\
& \rightarrow \frac{1}{2 \operatorname{Re}(\lambda)}\left(\frac{\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}}{s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)}-\frac{\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}}{s(\widetilde{\boldsymbol{B}})}\right)=\frac{\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}}{(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) s(\widetilde{\boldsymbol{B}})}
\end{aligned}
$$

as $t \rightarrow \infty$, since $s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)>0$ and $s(\widetilde{\boldsymbol{B}})>0$ imply

$$
\mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t} \int_{0}^{t} \mathrm{e}^{\left(\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}\right) w} \mathrm{~d} w \rightarrow \frac{\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}}{s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)}, \quad \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \rightarrow \frac{\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}}{s(\widetilde{\boldsymbol{B}})}
$$

as $t \rightarrow \infty$, see, e.g., the proof of Proposition B. 1 in Barczy et al. [7]. Hence, if $\operatorname{Re}(\lambda) \in$ $\left(-\infty, \frac{1}{2} s(\widetilde{\boldsymbol{B}})\right) \backslash\{0\}$, then

$$
\begin{aligned}
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} I_{\lambda, \ell}(t) & \rightarrow \frac{1}{s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)} \boldsymbol{e}_{\ell}^{\top} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\left(\mathbb{E}\left(\boldsymbol{X}_{0}\right)+\frac{\widetilde{\boldsymbol{\beta}}}{s(\widetilde{\boldsymbol{B}})}\right) \\
& =\frac{\left\langle\boldsymbol{e}_{\ell}, \widetilde{\boldsymbol{u}}\right\rangle}{s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)}\left(\left\langle\boldsymbol{u}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\widetilde{\boldsymbol{B}})}\right) \quad \text { as } t \rightarrow \infty .
\end{aligned}
$$

If $\operatorname{Re}(\lambda)=\frac{1}{2} s(\widetilde{\boldsymbol{B}})$, then we have

$$
t^{-1} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \boldsymbol{A}_{\lambda, 1}(t)=\frac{1}{t} \int_{0}^{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u} \mathrm{e}^{u \widetilde{\boldsymbol{B}}} \mathrm{~d} u \rightarrow \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top} \quad \text { as } t \rightarrow \infty
$$

see, e.g., part (v) of Lemma A. 2 in Barczy et al. [7]. Moreover, if $\operatorname{Re}(\lambda)=\frac{1}{2} s(\widetilde{\boldsymbol{B}})$, then, by Fubini's theorem, we obtain

$$
\begin{aligned}
t^{-1} \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t} \boldsymbol{A}_{\lambda, 2}(t) & =\frac{1}{t} \int_{0}^{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u}\left(\int_{0}^{u} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w\right) \mathrm{d} u=\frac{1}{t} \int_{0}^{t}\left(\int_{w}^{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) u} \mathrm{~d} u\right) \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \\
& =\frac{1}{s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t}\left(\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) w}-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\right) \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \\
& =\frac{1}{s(\widetilde{\boldsymbol{B}})}\left(\frac{1}{t} \int_{0}^{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) w} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w-\frac{1}{t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w\right) \rightarrow \frac{1}{s(\widetilde{\boldsymbol{B}})} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}
\end{aligned}
$$

as $t \rightarrow \infty$, since

$$
\begin{equation*}
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \rightarrow \frac{\widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}}{s(\widetilde{\boldsymbol{B}})} \quad \text { as } \quad t \rightarrow \infty \tag{B.1}
\end{equation*}
$$

see, e.g., Barczy et al. [7, formula (B.2)]. Hence, if $\operatorname{Re}(\lambda)=\frac{1}{2} s(\widetilde{\boldsymbol{B}})$, then

$$
t^{-1} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} I_{\lambda, \ell}(t) \rightarrow \boldsymbol{e}_{\ell}^{\top} \widetilde{\boldsymbol{u}} \boldsymbol{u}^{\top}\left(\mathbb{E}\left(\boldsymbol{X}_{0}\right)+\frac{\widetilde{\boldsymbol{\beta}}}{s(\widetilde{\boldsymbol{B}})}\right)=\left\langle\boldsymbol{e}_{\ell}, \widetilde{\boldsymbol{u}}\right\rangle\left(\left\langle\boldsymbol{u}, \mathbb{E}\left(\boldsymbol{X}_{0}\right)\right\rangle+\frac{\langle\boldsymbol{u}, \widetilde{\boldsymbol{\beta}}\rangle}{s(\widetilde{\boldsymbol{B}})}\right)
$$

as $t \rightarrow \infty$. If $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$, then we have

$$
\mathrm{e}^{-2 \operatorname{Re}(\lambda) t} \boldsymbol{A}_{\lambda, 1}(t)=\int_{0}^{t} \mathrm{e}^{u\left(\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}\right)} \mathrm{d} u \rightarrow \int_{0}^{\infty} \mathrm{e}^{u\left(\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}\right)} \mathrm{d} u=\left(2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}-\widetilde{\boldsymbol{B}}\right)^{-1}
$$

as $t \rightarrow \infty$, since $\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d} \in \mathbb{R}_{(+)}^{d \times d}$ is irreducible and $s\left(\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}\right)=s(\widetilde{\boldsymbol{B}})-$ $2 \operatorname{Re}(\lambda)<0$, see, e.g., the proof of Proposition B. 1 in Barczy et al. [7]. Moreover, if $\operatorname{Re}(\lambda) \in$ $\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$, then, by Fubini's theorem and (B.1), we obtain

$$
\begin{aligned}
& \mathrm{e}^{-2 \operatorname{Re}(\lambda) t} \boldsymbol{A}_{\lambda, 2}(t)=\int_{0}^{t} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u}\left(\int_{0}^{u} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w\right) \mathrm{d} u=\int_{0}^{t}\left(\int_{w}^{t} \mathrm{e}^{-2 \operatorname{Re}(\lambda) u} \mathrm{~d} u\right) \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \\
& =\frac{1}{2 \operatorname{Re}(\lambda)} \int_{0}^{t}\left(\mathrm{e}^{-2 \operatorname{Re}(\lambda) w}-\mathrm{e}^{-2 \operatorname{Re}(\lambda) t}\right) \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \\
& =\frac{1}{2 \operatorname{Re}(\lambda)} \int_{0}^{t} \mathrm{e}^{w\left(\widetilde{\boldsymbol{B}}-2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}\right)} \mathrm{d} w-\mathrm{e}^{-(2 \operatorname{Re}(\lambda)-s(\widetilde{\boldsymbol{B}})) t} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} \int_{0}^{t} \mathrm{e}^{w \widetilde{\boldsymbol{B}}} \mathrm{~d} w \\
& \rightarrow \frac{1}{2 \operatorname{Re}(\lambda)}\left(2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}-\widetilde{\boldsymbol{B}}\right)^{-1} \quad \text { as } t \rightarrow \infty
\end{aligned}
$$

Hence, if $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$, then

$$
\mathrm{e}^{-2 \operatorname{Re}(\lambda) t} I_{\lambda, \ell}(t) \rightarrow \boldsymbol{e}_{\ell}^{\top}\left(2 \operatorname{Re}(\lambda) \boldsymbol{I}_{d}-\widetilde{\boldsymbol{B}}\right)^{-1}\left(\mathbb{E}\left(\boldsymbol{X}_{0}\right)+\frac{\widetilde{\boldsymbol{\beta}}}{2 \operatorname{Re}(\lambda)}\right) \quad \text { as } t \rightarrow \infty
$$

Further, we have

$$
I_{\lambda}(t)=\int_{0}^{t} \mathrm{e}^{2 \operatorname{Re}(\lambda) w} \mathrm{~d} w= \begin{cases}t & \text { if } \operatorname{Re}(\lambda)=0 \\ \frac{1}{2 \operatorname{Re}(\lambda)}\left(\mathrm{e}^{2 \operatorname{Re}(\lambda) t}-1\right) & \text { if } \operatorname{Re}(\lambda) \in(-\infty, s(\widetilde{\boldsymbol{B}})] \backslash\{0\}\end{cases}
$$

If $\operatorname{Re}(\lambda)=0$, then

$$
\mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t} I_{\lambda}(t)=t \mathrm{e}^{-s(\tilde{\boldsymbol{B}}) t} \rightarrow 0 \quad \text { as } t \rightarrow \infty
$$

If $\operatorname{Re}(\lambda) \in\left(-\infty, \frac{1}{2} s(\widetilde{\boldsymbol{B}})\right) \backslash\{0\}$, then

$$
\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} I_{\lambda}(t)=\frac{1}{2 \operatorname{Re}(\lambda)}\left(\mathrm{e}^{-(s(\widetilde{\boldsymbol{B}})-2 \operatorname{Re}(\lambda)) t}-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\right) \rightarrow 0 \quad \text { as } t \rightarrow \infty
$$

If $\operatorname{Re}(\lambda)=\frac{1}{2} s(\widetilde{\boldsymbol{B}})$, then $\operatorname{Re}(\lambda) \neq 0$ and

$$
t^{-1} \mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t} I_{\lambda}(t)=\frac{1}{s(\widetilde{\boldsymbol{B}})} t^{-1}\left(1-\mathrm{e}^{-s(\widetilde{\boldsymbol{B}}) t}\right) \rightarrow 0 \quad \text { as } \quad t \rightarrow \infty
$$

If $\operatorname{Re}(\lambda) \in\left(\frac{1}{2} s(\widetilde{\boldsymbol{B}}), s(\widetilde{\boldsymbol{B}})\right]$, then $\operatorname{Re}(\lambda) \in \mathbb{R}_{++}$and

$$
\mathrm{e}^{-2 \operatorname{Re}(\lambda) t} I_{\lambda}(t)=\frac{1}{2 \operatorname{Re}(\lambda)}\left(1-\mathrm{e}^{-2 \operatorname{Re}(\lambda) t}\right) \rightarrow \frac{1}{2 \operatorname{Re}(\lambda)} \quad \text { as } t \rightarrow \infty
$$

The proof is complete.

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