



# Level structures of $^{56,58}\text{Ca}$ cast doubt on a doubly magic $^{60}\text{Ca}$

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## ABSTRACT

Gamma decays were observed in  $^{56}\text{Ca}$  and  $^{58}\text{Ca}$  following quasi-free one-proton knockout reactions from  $^{57,59}\text{Sc}$  beams at  $\approx 200$  MeV/nucleon. For  $^{56}\text{Ca}$ , a  $\gamma$  ray transition was measured to be 1456(12) keV, while for  $^{58}\text{Ca}$  an indication for a transition was observed at 1115(34) keV. Both transitions were tentatively assigned as the  $2_1^+ \rightarrow 0_{gs}^+$  decays, and were compared to results from *ab initio* and conventional shell-model approaches. A shell-model calculation in a wide model space with a marginally modified effective nucleon-nucleon interaction depicts excellent agreement with experiment for  $2_1^+$  level energies, two-neutron separation energies, and reaction cross sections, corroborating the formation of a new nuclear shell above the  $N = 34$  shell. Its constituents, the  $0f_{5/2}$  and  $0g_{9/2}$  orbitals, are almost degenerate. This degeneracy precludes the possibility for a doubly magic  $^{60}\text{Ca}$  and potentially drives the dripline of Ca isotopes to  $^{70}\text{Ca}$  or even beyond.

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Understanding properties of atomic nuclei at the extremes, for example those with large proton-to-neutron imbalances, is of paramount importance in nuclear physics. In these systems, often called exotic nuclei, new features emerge [1] including those that can be traced back to facets of nuclear forces. For instance, the tensor force, which has been known for decades [2,3], can modify the spin-orbit energy splitting as a function of the proton number ( $Z$ ) or the neutron number ( $N$ ), resulting in changes of shell structures, i.e., shell evolution [4,5]. Examples have been found in several regions across the Segrè chart (see review papers, [6,7]). Among them, the Ca isotopes provide an exemplary case of shell evolution, with striking appearances of the new magic numbers  $N = 32$  [8–10] and  $N = 34$  [11–14]. The discovery of new magic numbers is usually followed by the exploration of the new nuclear shell lying above them, which may yield precious hints of the whereabouts of the dripline [15–17]. This letter presents a finding along these lines based on state-of-the-art experimental and theoretical studies.

The Ca isotopes correspond to a complete filling of the  $Z = 20$  shell, leading to a high sensitivity of the shell evolution according to the neutron number. Signatures of magicity or sub-shell closures have been observed in the Ca isotopes at  $N = 16, 20, 28, 32$ , and  $34$  based on the steep decrease of the two-neutron separation energies  $S_{2n}$  [10,13] and the enhancement of the excitation energy of the first excited state  $E(2_1^+)$  [8,18,12,19]. The ground state of  $^{54}\text{Ca}$  has been shown, by knockout reactions, to have a closed-shell configuration [14], supporting the  $N = 34$  magicity. Having the  $N = 34$  magic number thus confirmed, the nexus of interest is the shell above it. If the shell is composed only of the  $0f_{5/2}$  orbital, the recently observed  $^{60}\text{Ca}$  [20] may be doubly magic and become a dripline nucleus. However, if the orbitals above  $0f_{5/2}$  contribute substantially, the dripline can be located deep into the *terra incognita* of the Segrè chart. The influence of the *gds* orbitals above the *pf* shell is often discussed in the literature when neutron-rich Ca, Ti, and Ni isotopes are addressed [21–26]. There are, however, no experimental data probing this shell in the Ca isotopes.

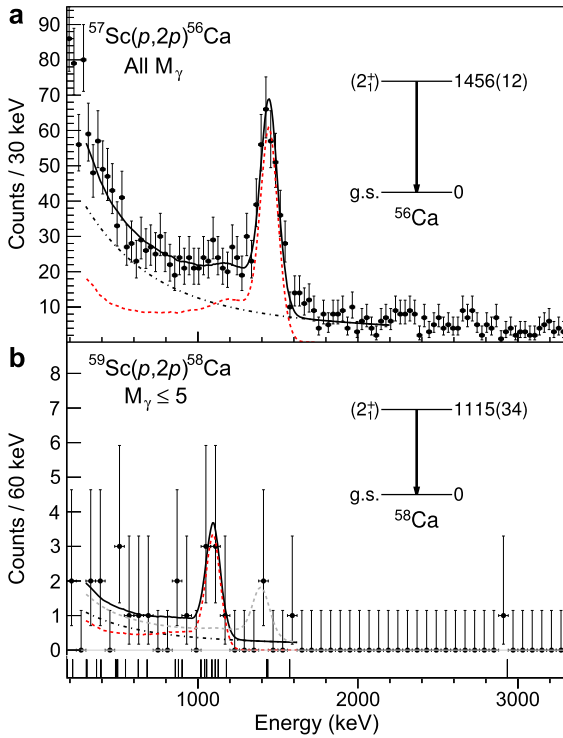
Theoretical predictions of the level structure of Ca isotopes beyond  $^{54}\text{Ca}$  and the location of the dripline have been made by modern shell-model, *ab initio*, beyond mean field calculations, and energy density functionals [21,22,27–34,16,35]. There seems to be no sign of convergence or consistency of such predictions for the level structure of  $^{56,58}\text{Ca}$ , as discussed later. In fact, the predicted values of  $E(2_1^+)$  for  $^{56,58}\text{Ca}$  range from 0.5 to 2 MeV [27,31,32,30,28,33,34]. Such a large variance prevents useful insights or conclusions regarding the shell structure beyond the  $N = 34$  (sub-)shell closure. Recent predictions of a newly developed fitted interaction within the *fp*-model space, tailored for the neutron-rich Ca isotopes, imply  $^{60}\text{Ca}$  being doubly magic at a similar level to  $^{68}\text{Ni}$  [34,36]. This prediction is, however, strongly

dependent on the agreement to experimental data for  $^{55-59}\text{Ca}$  [34,36]. The closest isotope along  $N = 40$  with experimental information,  $^{62}\text{Ti}$ , showed no indication for a new magic number [26], in agreement with the predictions presented in Ref. [37]. This letter reports on the first measurement of excitation energies of  $^{56,58}\text{Ca}$  by means of in-beam  $\gamma$ -ray spectroscopy. Experimental data were confronted with modern shell-model and *ab initio* calculations combined with reaction theory.

The experiment was carried out at the Radioactive Isotope Beam Factory, operated by the RIKEN Nishina Center and the Center for Nuclear Study, the University of Tokyo. Radioactive beams were produced by fragmentation of a  $^{70}\text{Zn}$  beam at 345 MeV/nucleon on a 10-mm-thick  $^9\text{Be}$  target. The  $^{57,59}\text{Sc}$  isotopes were then separated and identified from focal plane F0 to F13 of the BigRIPS separator [38]. Afterwards, the secondary beams with intensities of 13.6 particles/s for  $^{57}\text{Sc}$  and 0.3 particles/s for  $^{59}\text{Sc}$  impinged on the MINOS liquid-hydrogen ( $\text{LH}_2$ ) target [39] to induce proton knockout reactions. Reaction residues,  $^{56,58}\text{Ca}$ , were identified by the SAMURAI spectrometer [40]. Secondary beam energies at the target center were 209 MeV/nucleon for  $^{57}\text{Sc}$  and 199 MeV/nucleon for  $^{59}\text{Sc}$ , inducing considerable Doppler shifts for the emitted  $\gamma$  rays. The DALI2<sup>+</sup> detector array [41] was used to measure the de-excitation  $\gamma$  rays. To overcome the large Doppler broadening partially caused by the long  $\text{LH}_2$  target, Doppler corrections were performed using the reaction vertex information reconstructed by the MINOS time projection chamber. For further experimental details, the interested reader is referred to the supplemental material.

The Doppler-corrected  $\gamma$ -ray spectrum in coincidence with the  $^{57}\text{Sc}(p,2p)^{56}\text{Ca}$  reaction is shown in Fig. 1a. A single peak is observed at 1456(12) keV and tentatively assigned to the  $2_1^+ \rightarrow 0_{gs}^+$  transition. Energy uncertainties are dominated by the fitting error and energy calibration. Lifetime effects on the measured energies were also evaluated. The heaviest Ca isotope with a known  $B(E2)^\uparrow$  is  $^{50}\text{Ca}$  [42]. Assuming the same transition strength,  $37.5(10) e^2 \text{ fm}^4$ , for  $^{56}\text{Ca}$  gives a lifetime of 17 ps. This lifetime value was adopted with an error of 100% and taken into account in the error determination.

Despite low statistics, the Doppler-corrected  $\gamma$ -ray energy spectrum of the  $^{59}\text{Sc}(p,2p)^{58}\text{Ca}$ , shown in Fig. 1b, revealed a peak-like structure in the energy range of 1000–1200 keV. To test the significance level of this peak, a maximum likelihood fit procedure was applied to the unbinned data (see bottom of Fig. 1b). The background in this spectrum was modeled from the  $^{57}\text{Sc}(p,2p)^{56}\text{Ca}$  reaction, with the amplitude normalized according to the event numbers. This procedure was validated with the  $^{55}\text{Sc}(p,2p)^{54}\text{Ca}$  data from the same experiment, which yielded a good description of the background. A significance of  $2.8 \sigma$ , defined as the peak amplitude over the statistical uncertainty from the maximum likelihood fit, was obtained for the tentative 1115(34) keV  $\gamma$ -ray

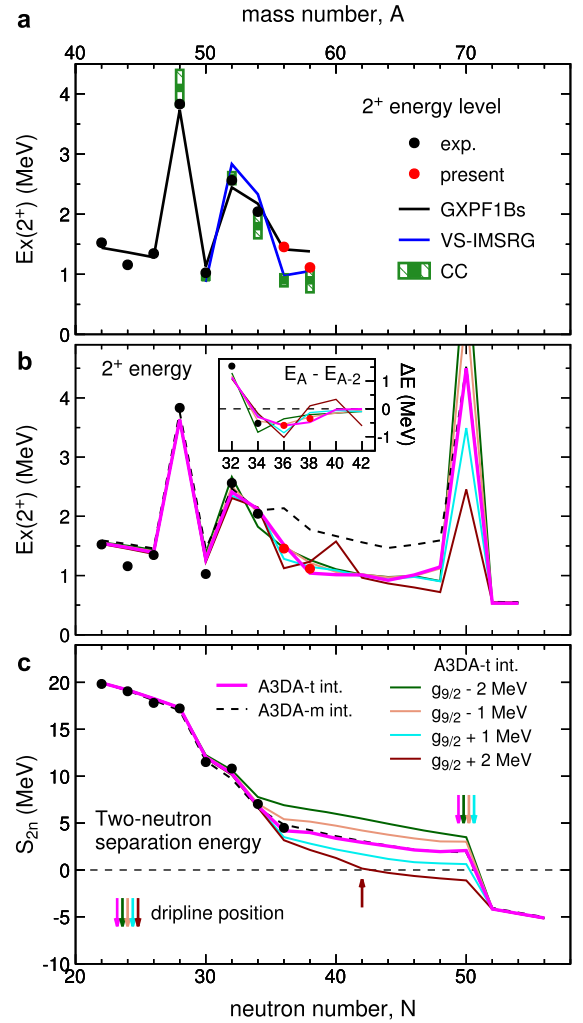


**Fig. 1.** Doppler-corrected  $\gamma$ -ray spectra. **a**, Spectrum in coincidence with the  $^{57}\text{Sc}(p,2p)^{56}\text{Ca}$  reaction. **b**, Spectrum in coincidence with the  $^{59}\text{Sc}(p,2p)^{58}\text{Ca}$  reaction restricted to  $\gamma$  multiplicity  $\leq 5$  (spectra with other multiplicities are shown in Fig. 3 of supplemental material). Spectra were fitted with simulated DALI2<sup>+</sup> response functions (red) with a two-exponential background (black). Poisson-statistics errors were adopted for the data points. Unbinned data are shown in the bottom of panel **b** for the  $^{59}\text{Sc}(p,2p)^{58}\text{Ca}$  channel. The response curve for an assumed 1400 keV  $\gamma$ -ray transition with a theoretical cross section of 0.25 mbarn is indicated for  $^{58}\text{Ca}$  by the gray dashed line.

transition, including systematic errors from lifetime effects, and tentatively assigned to the  $2_1^+ \rightarrow 0_{\text{gs}}^+$  decay of  $^{58}\text{Ca}$ . The assumed lifetime is 66 ps based on the same assumption as  $^{56}\text{Ca}$ . Noteworthy are the two counts observed at  $\sim 1400$  keV, comparable to the  $E(2_1^+)$  of  $^{56}\text{Ca}$ . However, taking into account calculated theoretical cross sections, as discussed below, resulted in a poor overall agreement of the response function with the data, as evidenced by the gray dashed line in Fig. 1b. Further tests for the validity and the impact of the  $^{58}\text{Ca}$  data is discussed in the supplemental material.

The systematics of  $E(2_1^+)$  values as a function of neutron number presented in Fig. 2a evince the expected pattern for magic nuclei at  $N = 28$ : A sharp increase from  $N = 26$  to 28 followed by a large reduction at  $N = 30$ . Similarly, a characteristic sharp increase from  $N = 30$  to 32 exists for the  $N = 32$  magic number, while the enhanced  $E(2_1^+)$  at  $N = 34$  is indicative of magicity. The firmly established data point for  $^{56}\text{Ca}$  and the tentative one for  $^{58}\text{Ca}$  are as low as the  $N = 22, 24, 26$ , and 30 values, with a decrease from  $N = 36$  to 38.

The  $E(2_1^+)$  systematics of Ca isotopes were compared to conventional shell-model calculations with the GXPFI1Bs Hamiltonian in the model space of the full  $pf$  shell [14,22], and two state-of-the-art *ab initio* approaches: The valence-space in-medium similarity renormalization group (VS-IMSRG) [43–46] and the coupled-cluster theory (CC) [47], both employing the two- (NN) and three-nucleon (3N) interaction 1.8/2.0 (EM) [48], derived from chiral effective field theory [49]. Details of these theoretical approaches are provided in the supplemental material. Fig. 2a shows the theoretical calculations well describe the  $E(2_1^+)$  excitation energies up to  $N = 34$ , and the GXPFI1Bs Hamiltonian also provides a good agree-



**Fig. 2.** Comparison of calculated  $E(2_1^+)$  and  $S_{2n}$  values with experimental data. **a**,  $E(2_1^+)$  systematics in even-even Ca isotopes confronted with theoretical approaches: The shell model using the GXPFI1Bs Hamiltonian, the VS-IMSRG method, and CC calculations. **b**,  $E(2_1^+)$  systematics in even-even Ca isotopes, and their differences (inset). Experimental points are the same as **a**. The calculated values are obtained by the original A3DA-m Hamiltonian as well as its revised one (A3DA-t). **c**,  $S_{2n}$  systematics in even-even Ca isotopes. Also shown in **b** and **c** are the effect of shifting the neutron  $0g_{9/2}$  orbital for predictions of A3DA-t.

ment with the present experimental value for  $^{56}\text{Ca}$ . However, all these calculations predict a flat behavior from  $N = 36$  to 38.

A more general discussion provides an instructive viewpoint of the  $E(2_1^+)$  values of  $^{56,58}\text{Ca}$ . If the  $0f_{5/2}$  orbital is isolated from the other orbitals,  $N = 36$  corresponds to a system of two neutrons solely occupying the  $0f_{5/2}$  orbital. Likewise,  $N = 38$  would be four neutrons in the  $0f_{5/2}$  orbital, or, equivalently, two neutron holes of the fully occupied  $0f_{5/2}$  orbital. The two-body interaction is invariant between particle and hole systems, but the single-particle energies can vary with neutron number. Such changes of single-particle energies do not affect excitation level energies, because only one orbital is relevant. Thus, assuming the  $0f_{5/2}$  neutron orbital is marginally modified between  $^{56}\text{Ca}$  and  $^{58}\text{Ca}$ , the  $E(2_1^+)$  value should be identical between  $N = 36$  and  $N = 38$  as a consequence of this particle-hole symmetry. It is emphasized that this consequence is independent of the choice of the two-body interaction.

The present results indicate a decrease of  $E(2_1^+)$  from  $N = 36$  to 38 by several hundred keV. This observation conflicts with the arguments above, implying a non-isolated  $0f_{5/2}$  orbital. Since all experimental evidence supports an  $N = 34$  magic number in the

Ca isotopes, the  $0f_{5/2}$  orbital is considered as isolated from lower-energy orbitals. This points to the other possibility that the  $0f_{5/2}$  orbital is coupled to higher orbitals, suggesting a shell comprising the  $0f_{5/2}$  orbital and at least one higher orbital. This new shell has never been discussed and its appearance excludes the  $N = 40$  magic number in Ca isotopes.

A previous theoretical study has discussed an “sdg” shell built on an inert  $^{60}\text{Ca}$  core [24]. This approach proved valid for the  $^{78}\text{Ni}$  region, but remains untested for the Ca isotopes. In the present work, only the characteristics of the  $0g_{9/2}$  and  $1d_{5/2}$  orbitals can be constrained by experiment. As the protons can be assumed to form a  $Z = 20$  closed shell, only neutrons above  $N = 20$  are treated as valence nucleons in the calculations.

The existing effective A3DA-m [23]  $NN$  interaction, defined for a model space comprising the full  $pf$  shell, the  $0g_{9/2}$ , and  $1d_{5/2}$  orbitals, is used as a starting point. It has been successfully used for the systematic descriptions of Ni ( $Z = 28$ ) [23] and Cu ( $Z = 29$ ) [50] isotopes. Fig. 2b shows that the observed  $E(2_1^+)$  values are well reproduced by the A3DA-m interaction up to  $N = 34$ , and substantial deviations for  $N = 36$  and 38. The excitation-energy lowering from  $N = 36$  to 38 is well reproduced by the A3DA-m interaction, in contrast to trends shown in Fig. 2a. This suggests a minor revision of the interaction may be sufficient to reproduce the experimental data. Fig. 2c shows the  $S_{2n}$  values for the Ca isotopes. There is no notable deviation for the nuclei where experimental data are available, implying the validity of the A3DA-m interaction.

The A3DA-m interaction is revised by varying only the two-body-matrix-elements (TBME) in the linear combination (LC) method [51,52] to better reproduce the  $E(2_1^+)$  values of  $^{54,56,58}\text{Ca}$ . Changes to the TBMEs are small, as expected. The maximum change is 0.198 MeV, while the others are much smaller. The correlation between the original TBMEs and the revised TBMEs are shown in the supplemental material.

The revised interaction is labeled “A3DA-t” hereafter. Fig. 2b depicts the  $E(2_1^+)$  values obtained with the A3DA-m and the A3DA-t interactions from  $^{42}\text{Ca}$  to  $^{74}\text{Ca}$ , with A3DA-t reproducing the  $E(2_1^+)$  values of  $^{56,58}\text{Ca}$ . The  $E(2_1^+)$  value remains almost constant until  $^{68}\text{Ca}$ , where the value for  $^{70}\text{Ca}$  rises due to the filled  $0f_{5/2}$ -plus- $0g_{9/2}$  shell and the necessity of neutron excitations to the high-lying  $1d_{5/2}$  orbital. As orbitals above this are not included, the present work cannot describe excitation energies much beyond  $^{70}\text{Ca}$ .

Fig. 2c shows  $S_{2n}$  values up to  $^{76}\text{Ca}$ , the last possible nucleus in the present model space. A plateau is formed from  $^{56}\text{Ca}$  to  $^{70}\text{Ca}$ . Beyond this,  $S_{2n}$  becomes negative, implying the dripline is located at  $^{70}\text{Ca}$ , close to some predictions [53,15], beyond others [54] or within argued ranges [20,16,35]. Inclusion of higher *sdg* orbitals may slant the dripline even further due to quadrupole correlations [24,25]. As noted in Ref. [24], neutrons in higher *sdg* orbitals can enhance quadrupole collectivity, which may lead to a well-deformed ground state of  $^{70}\text{Ca}$ .

The sensitivity of the neutron  $0g_{9/2}$  single-particle-energy (SPE) was characterized by varying it up to  $\pm 2$  MeV with respect to the original value of the A3DA-t interaction. Results of this can be seen in Fig. 2b for the  $E(2_1^+)$  and in Fig. 2c for the  $S_{2n}$ . Of particular interest is the difference of  $E(2_1^+)$ , defined as  $\Delta E = E(2_1^+) - E(2_{A-2}^+)$  and shown in the inset of Fig. 2b. The larger the neutron  $0g_{9/2}$  SPE, the larger the drop from  $^{54}\text{Ca}$  to  $^{56}\text{Ca}$ , producing a local  $E(2_1^+)$  maximum for  $^{60}\text{Ca}$  and shifting the dripline to  $^{62}\text{Ca}$ . A positive shift of +1 or +2 MeV can be excluded from the experimental  $E(2_1^+)$  and  $S_{2n}$  of  $^{56}\text{Ca}$ . Conversely, a too low  $0g_{9/2}$  SPE value quenches the experimentally established magicity at  $N = 34$  [12–14], resulting in a high neutron  $0g_{9/2}$  occupation number not observed in  $^{54}\text{Ca}$  [14]. Our results obtained from  $^{56}\text{Ca}$  challenge

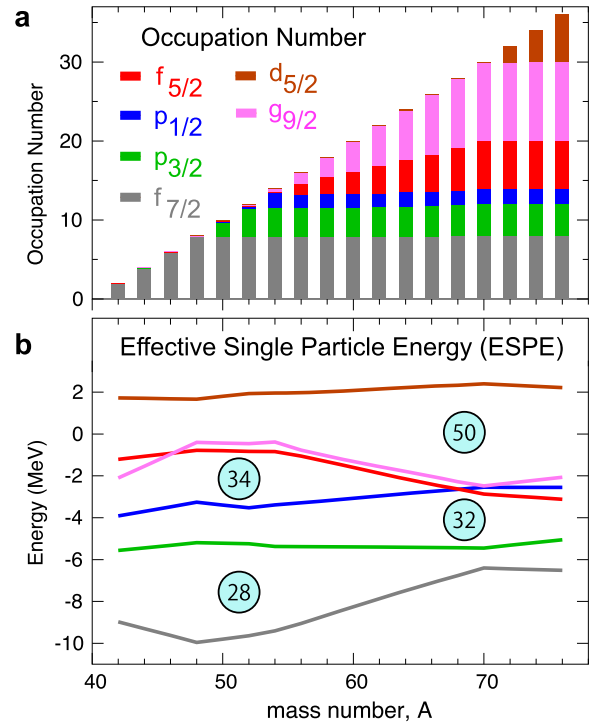


Fig. 3. a, Occupation numbers of the ground states of even-even Ca isotopes obtained from the A3DA-t interaction. b, Effective single-particle energies for the same states as panel a. The numbers in the circles are the neutron numbers corresponding to the magic gaps.

the notion of an  $N = 40$  magicity at  $^{60}\text{Ca}$  and are reinforced by the tentative experimental value for  $^{58}\text{Ca}$ , as all  $\Delta E$  remain negative except for the  $0g_{9/2}$  SPE shifted by +2 MeV.

The occupation number of each single-particle orbital is shown in Fig. 3a. Likewise, the ESPE [7] are displayed in Fig. 3b. The shell structure above  $N = 34$  is clearly characterized by two orbitals,  $0f_{5/2}$  and  $0g_{9/2}$ , which remain almost degenerate across the range shown in Fig. 3b. Thus, the emergence of a new shell above  $N = 34$  comprising  $0f_{5/2}$ ,  $0g_{9/2}$  orbitals and some others, like  $1d_{5/2}$ , is evident. The  $N = 34$  magic gap decreases beyond  $A = 60$ , but has little affect because the  $1p_{1/2}$  orbital remains almost completely occupied. This degeneracy is lifted for isotopes with  $Z > 20$  due to the strong monopole attraction between a proton in  $0f_{7/2}$  and a neutron in  $0f_{5/2}$  [7].

Further discussions are concentrated on the cross sections. Inclusive cross sections for the  $^{57}\text{Sc}(p,2p)^{56}\text{Ca}$  and  $^{59}\text{Sc}(p,2p)^{58}\text{Ca}$  reactions were measured to be 1.23(5) and 1.14(15) mb, respectively. Partial cross sections of the excited states were extracted using the efficiency-corrected  $\gamma$ -ray intensities, and those to the ground-states deduced by subtraction. All measured cross sections are summarized in Table 1. Inclusive and partial cross sections were comparable between both nuclei, lending support to the assignment of a peak in  $^{58}\text{Ca}$ .

Theoretical cross sections were obtained by combining single-particle cross sections  $\sigma_{sp}$  calculated from the distorted-wave impulse approximation (DWIA) and the spectroscopic factors  $C^2S$  from the GXPF1Bs and A3DA-t Hamiltonians, and VS-IMSFG approach described above [55]. They are listed in Table 1. The beams of ground-state  $^{57,59}\text{Sc}$  have assumed  $J^\pi = 7/2^-$ . Only removal from the proton  $0f_{7/2}$  orbital was considered, as higher-lying proton orbitals contributed only a few percent to the final states. Negligible cross sections were calculated to states other than the listed  $0_{g.s.}^+$ ,  $2_1^+$ , and  $4_1^+$  states.

Similar inclusive cross sections for both reaction channels were predicted, as observed experimentally, and with  $\sigma_{\text{exp-to-}\sigma_{\text{th}}}$  ra-



**Table 1**

Observed excitation energies ( $E_{\text{exp}}$ ) in keV and cross sections ( $\sigma_{\text{exp}}$ ) in mbarn from the  $^{57}\text{Sc}(p,2p)^{56}\text{Ca}$  and  $^{59}\text{Sc}(p,2p)^{58}\text{Ca}$  reaction channels compared to theoretical values ( $\sigma_{\text{th}}$ ) using the DWIA calculated single-particle cross sections ( $\sigma_{\text{sp}}$ ) and spectroscopic factors ( $C^2S$ ) from VS-IMSRG, GXPF1Bs, and A3DA-t. Predicted spin-parities ( $J^\pi$ ), associated proton-removal orbitals ( $nl_j$ ), and excitation energies ( $E_x$ ) are also provided.

	Experiment		$J^\pi$	$nl_j$	DWIA	VS-IMSRG			GXPF1Bs			A3DA-t		
	$E_{\text{exp}}$	$\sigma_{\text{exp}}$			$\sigma_{\text{sp}}$	$E_x$	$C^2S_{\text{th}}$	$\sigma_{\text{th}}$	$E_x$	$C^2S_{\text{th}}$	$\sigma_{\text{th}}$	$E_x$	$C^2S_{\text{th}}$	$\sigma_{\text{th}}$
$^{56}\text{Ca}$	0	0.80(6)	$0^+_{\text{g.s.}}$	$0f_{7/2}$	1.80	0	0.61	1.10	0	0.69	1.24	0	0.62	1.11
	1456(12)	0.43(4)	$2^+_{\text{1}}$	$0f_{7/2}$	1.74	1002	0.29	0.50	1416	0.25	0.44	1519	0.27	0.47
			$4^+_{\text{1}}$	$0f_{7/2}$	1.73	1307	0.05	0.09	1776	0.02	0.04	2339	0.01	0.02
	Inclusive	1.23(5)						1.69			1.72			1.60
$^{58}\text{Ca}$	0	0.66(24)	$0^+_{\text{g.s.}}$	$0f_{7/2}$	1.58	0	0.80	1.26	0	0.83	1.31	0	0.46	0.73
	1115(34)	0.47(19)	$2^+_{\text{1}}$	$0f_{7/2}$	1.54	1075	0.16	0.25	1382	0.15	0.23	1040	0.42	0.65
			$4^+_{\text{1}}$	$0f_{7/2}$	1.52	1423	0.001	0.002	1772	0.001	0.002	2084	0.05	0.08
	Inclusive	1.14(15)						1.51			1.54			1.46

tios  $\sim 0.75$ , agreeing with previous values obtained in the region [26,56,57] and for stable nuclei [58]. This signifies a low occupation number of protons across the  $Z = 20$  shell in the ground states of  $^{57,59}\text{Sc}$ , hence a good proton shell closure. A different picture is observed for the partial cross sections to the  $2^+_{\text{g.s.}}$  states. While the  $\sigma_{\text{exp-to-}\sigma_{\text{th}}$  ratio holds for  $^{56}\text{Ca}$ , despite considerable uncertainties, the experimental partial cross section for  $^{58}\text{Ca}$  is two times larger than the value predicted by the GXPF1Bs and VS-IMSRG calculations. In contrast, the A3DA-t Hamiltonian gives results in good agreement with experiment: Partial cross sections change from  $N = 36$  to 38 in a consistent manner with experiment.

In conclusion, the first spectroscopy measurements for  $^{56,58}\text{Ca}$  following the  $1p$ -knockout reactions from scandium isotopes were carried out. A  $\gamma$  ray transition associated with the  $2^+_{\text{g.s.}} \rightarrow 0^+_{\text{g.s.}}$  decay was assigned for  $^{56}\text{Ca}$  and an indication of this transition was observed for  $^{58}\text{Ca}$ . A comparison with standard shell-model and *ab initio* theoretical calculations exhibits a notable deficiency in their descriptions of nuclear structure around  $N = 40$ . The particle-hole symmetry argument robustly leads to a new shell comprising at least the  $0f_{5/2}$  and  $0g_{9/2}$  orbitals above  $N = 34$ . The fitted A3DA-t interaction, introduced in this work, shows an excellent description of so far known experimental data, and predicts the Ca dripline at  $N = 50$ , because of substantial correlation energies from the pairing between the  $0f_{5/2}$  and  $0g_{9/2}$  orbitals. Thus, the picture of the new magic number  $N = 34$  [11–14] becomes more complete with a shell built atop of it. More detailed structure information of  $^{56,58}\text{Ca}$  can be obtained by future measurements with Ge detector arrays at next-generation facilities like FRIB [59], notably validation of the tentative transition at 1115(34) keV. Additional experimental investigations of neutron-rich Ca isotopes, such as particle states in  $^{55}\text{Ca}$  from neutron pickup reactions and the spectroscopy of  $^{60}\text{Ca}$ , not believed to be doubly magic from the assessment of results presented here and Ref. [26], would provide key information to further characterize the new shell.

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### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Data availability

Data will be made available on request.

## Appendix A. Supplementary material

Supplementary material related to this article can be found online at <https://doi.org/10.1016/j.physletb.2023.138025>.

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