

SIMULTANEOUS MULTI-BAND RADIO & X-RAY OBSERVATIONS OF THE GALACTIC CENTER MAGNETAR SGR 1745–2900

T. T. PENNUCCI¹, A. POSSENTI², P. ESPOSITO^{3,4}, N. REA^{5,6}, D. HAGGARD⁷,
F. K. BAGANOFF⁸, M. BURGAY², F. COTI ZELATI^{9,5,10}, G. L. ISRAEL¹¹, A. MINTER¹²

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ABSTRACT

We report on multi-frequency, wideband radio observations of the Galactic Center magnetar (SGR 1745–2900) with the Green Bank Telescope for ~ 100 days immediately following its initial X-ray outburst in April 2013. We made multiple simultaneous observations at 1.5, 2.0, and 8.9 GHz, allowing us to examine the magnetar’s flux evolution, radio spectrum, and interstellar medium parameters (such as the dispersion measure (DM), the scattering timescale and its index). During two epochs, we have simultaneous observations from the *Chandra* X-ray Observatory, which permitted the absolute alignment of the radio and X-ray profiles. As with the two other radio magnetars with published alignments, the radio profile lies within the broad peak of the X-ray profile, preceding the X-ray profile maximum by ~ 0.2 rotations. We also find that the radio spectral index γ is significantly negative between ~ 2 and 9 GHz; during the final ~ 30 days of our observations $\gamma \sim -1.4$, which is typical of canonical pulsars. The radio flux has not decreased during this outburst, whereas the long-term trends in the other radio magnetars show concomitant fading of the radio and X-ray fluxes. Finally, our wideband measurements of the DMs taken in adjacent frequency bands in tandem are stochastically inconsistent with one another. Based on recent theoretical predictions, we consider the possibility that the dispersion measure is frequency-dependent. Despite having several properties in common with the other radio magnetars, such as $L_{X,\text{qui}}/L_{\text{rot}} \lesssim 1$, an increase in the radio flux during the X-ray flux decay has not been observed thus far in other systems.

Subject headings: Galaxy: center — pulsars: individual (PSR J1745–2900, SGR 1745–2900) — stars: magnetars

1. INTRODUCTION

Magnetars are exotica among the exotic: whereas other pulsars are sustained by their stored angular momentum, the primary energy source that powers this special class of objects is likely the neutron star’s immense magnetic field (Mereghetti et al. 2015). The field strengths take on the highest values ever inferred, typically $> 10^{12}$ G and even up to $\sim 10^{15}$ G. According to the McGill Online

Magnetar Catalog¹³ (Olausen & Kaspi 2014), there are 28 known magnetars, of which only four have displayed pulsed radio emission.

SGR 1745–2900 (J1745–2900, hereafter) is the most recent addition to the small collection of magnetars with observed pulsed radio emission (the “radio magnetars”, to which we will refer by their PSR names: J1809–1943 (XTE 1810–197), J1550–5418 (1E 1547.0–5408), & J1622–4950 (Camilo et al. 2006, 2007a; Levin et al. 2010)). On 25 April 2013, one day after the XRT aboard the *Swift* satellite detected flaring activity coincident with the Galactic Center (Degenaar et al. 2013), a short X-ray burst was observed by *Swift*/BAT showing characteristics similar to those usually observed from soft gamma-ray repeaters (Kennea et al. 2013c). Shortly thereafter, observations from the *NuSTAR* satellite identified the source as a magnetar with a $P_s = 3.76$ s spin period, and its radio pulsations were subsequently seen by the Effelsberg 100-m Telescope (Mori et al. 2013a,b; Eatough et al. 2013a). J1745–2900 was soon physically associated with the Galactic Center, located only $\sim 2.5''$ away from Sagittarius A* (Sgr A*) with a neutral hydrogen column density and dispersion measure (DM) consistent with being within ~ 2 pc of the Milky Way’s central black hole (Eatough et al. 2013b; Rea et al. 2013).

Early determinations of its spin-down \dot{P}_s put J1745–2900 squarely within the magnetar population, having an inferred magnetic field strength at the equator

¹University of Virginia, Department of Astronomy, PO Box 400325 Charlottesville, VA 22904-4325, USA, pennucci@virginia.edu

²Osservatorio Astronomico di Cagliari, via della Scienza 5, 09047, Cagliari, Italy

³Istituto di Astrofisica Spaziale e Fisica Cosmica - Milano, INAF, via E. Bassini 15, I-20133, Milano, Italy

⁴Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138, USA

⁵Anton Pannekoek Institute for Astronomy, University of Amsterdam, Postbus 94249, NL-1090-GE Amsterdam, The Netherlands

⁶Institute of Space Sciences (ICE, CSIC-IEEC), Carrer de Can Magrans, S/N, 08193, Barcelona, Spain

⁷Department of Physics and Astronomy, Amherst College, Amherst, MA 01002-5000, USA

⁸Kavli Institute for Astrophysics and Space Research, Massachusetts Institute of Technology, Cambridge, MA 02139, USA

⁹Università dell’Insubria, via Valleggio 11, I-22100 Como, Italy

¹⁰INAF-Osservatorio Astronomico di Brera, via Bianchi 46, I-23807, Merate (LC), Italy

¹¹INAF-Osservatorio Astronomico di Roma, via Frascati 33, I-00040, Monteporzio Catone, Roma, Italy

¹²National Radio Astronomy Observatory, PO Box 2, 155 Observatory Road, Green Bank, WV 24944, USA

¹³ <http://www.physics.mcgill.ca/~pulsar/magnetar/main.html>

$B_s \sim 3.2 \times 10^{19} \text{ G} \sqrt{P_s \dot{P}_s} \sim 1.6 \times 10^{14} \text{ G}$, a characteristic age $\tau_c \sim P_s / (2\dot{P}_s) \sim 9 \text{ kyr}$, and a spin-down luminosity of $\dot{E} = L_{\text{rot}} = 3.95 \times 10^{46} \text{ erg s}^{-1} (P_s^{-3} \dot{P}_s) \sim 4.9 \times 10^{33} \text{ erg s}^{-1}$ (Rea et al. 2013). However, its estimated quiescent X-ray luminosity of $L_{X,\text{qui}} < 10^{34} \text{ erg s}^{-1}$ (Coti Zelati et al. 2015) may place J1745–2900 on the side of $L_{X,\text{qui}}/L_{\text{rot}} < 1$, opposite the “classic magnetars” but alongside the other three radio magnetars, high- B pulsars, and radio pulsars with X-ray emission (Rea et al. 2012).

Given the unique environment in which J1745–2900 resides, the detection of its radio pulses is somewhat surprising. Indeed, numerous surveys of the Galactic Center region covering $\sim 1\text{--}20 \text{ GHz}$ have failed to find a pulsar within the central parsec (most recently, Johnston et al. 2006; Deneva et al. 2009; Macquart et al. 2010; Bates et al. 2011; Siemion et al. 2013). The discovery of this single magnetar has led to a windfall of implications for future discoveries (Chennamangalam & Lorimer 2014; Dexter & O’Leary 2014; Macquart & Kanekar 2015). Because of its proximity to the Galactic Center, J1745–2900 has the largest DM ($1778 \text{ cm}^{-3} \text{ pc}$) and rotation measure ($-6.696 \times 10^4 \text{ rad m}^{-2}$) of any known pulsar (Eatough et al. 2013b). The predicted value for the scattering timescale at 1 GHz, based on empirical relationships given its DM, is $\sim 1000 \text{ s}$ (Krishnakumar et al. 2015; Lewandowski et al. 2015a), meaning that J1745–2900 would be undetectable at frequencies less than $\sim 5 \text{ GHz}$. The situation is exacerbated by the presence of an additional scattering screen in the Galactic Center (Cordes & Lazio 1997). Normally, the prospect of detecting distant radio pulsars above several GHz is bleak, since their average spectral index is ~ -1.4 (Bates et al. 2013). However, because the other radio magnetars have flat/inverted spectra, one might expect to detect J1745–2900’s unscattered pulse profile at high frequencies. In the analyses that follow, we will reiterate the finding that J1745–2900 has a significantly smaller scattering timescale than predicted (Spitler et al. 2014), and will show that J1745–2900 was much brighter at lower frequencies, having a very negative spectral index some 100 days after the onset of its outburst, even though more recent observations by Torne et al. (2015) showed the spectral index has since flattened.

In this paper, we analyze multi-frequency radio data over the first ~ 100 days after J1745–2900’s discovery, during which time there were two additional *Swift*/BAT-detected bursts on 7 June 2013 and 5 August 2013 (Kennea et al. 2013a,b). For two of our epochs, which bracket the third burst by ~ 1 week on either side, we have simultaneous *Chandra* observations. These observations allow us to find the absolute alignment of the radio and X-ray profiles, and to look for correlated events. We comment on the spin evolution and timing, and examine the profile stability, the radio flux evolution, and the radio spectrum. Finally, we make global models of the profile evolution across the low frequency bands in order to examine the temporal and frequency dependencies of the scattering timescale and dispersion measure. We then discuss characteristics of this source in comparison with other radio-loud magnetars.

2. OBSERVATIONS

2.1. Radio

We made early detections of J1745–2900 during fourteen observing epochs with the 100-m Robert C. Byrd Green Bank Telescope (GBT) in three different frequency bands with various overlap: 1.1–1.9 GHz (5 epochs), 1.6–2.4 GHz (7 epochs), and 8.5–9.3 GHz (11 epochs) (PI: A. Possenti). Because each observation covers a large bandwidth, we refer to each set of data based on the IEEE radio band for which each of the receiver systems is named (“L-band”, “S-band”, or “X-band”, respectively), instead of referring to specific (central) frequencies. Table 1 contains details of the observations. In all cases, we observed using the Green Bank Ultimate Pulsar Processing Instrument (GUPPI¹⁴, DuPlain et al. 2008) in “incoherent search mode”, recording dual-polarization time-series data in 2048 frequency channels with a temporal resolution of 0.65536 ms.

Each epoch’s data were folded with the pulsar software library DSPSR¹⁵ using a nominal ephemeris with a constant spin frequency (see §3.3) and the *Chandra*-determined position $\alpha_{J2000.0} = 17^{\text{h}}45^{\text{m}}40^{\text{s}}.169$, $\delta_{J2000.0} = 29^{\circ}00'29''.84$ (Rea et al. 2013). The data were initially folded into 1 min subintegrations, with 2048 profile phase bins across 128 frequency channels. We adopted the published dispersion measure value of $1778 \text{ cm}^{-3} \text{ pc}$ for averaging frequency channels together (Eatough et al. 2013b). Persistent, narrow-band radio frequency interference (RFI) was excised automatically; any remaining significantly corrupted channels or subintegrations were removed from the data by hand.

Calibration scans were taken for each observation using the local noise diode, pulsed at 25 Hz while on source. We recorded on- and off-source scans of a standard flux calibrator (QSO B1442+101) in each frequency band only during the final epoch (MJD 56516). We have used this one set of flux calibration scans to calibrate the whole data set. Standard programs from the PSRCHIVE¹⁶ pulsar software library (Hotan et al. 2004; van Straten et al. 2012) were used to calibrate the absolute flux density scale of the noise diode, which is then used to determine the magnetar’s flux density¹⁷.

The combination of the large amount of observed scattering (§3.5.1), the pulsar’s spectrum (§3.4.3, Figure 7), receiver roll-off, and the presence of gain variations (see below) rendered significant portions of the ends of L-band useless. Namely, there was no pulsed signal in the lower 300 MHz portion of L-band, which we masked from further analysis, along with the top 50 MHz (which is part of the overlap with S-band). In combination with

¹⁴ www.safe.nrao.edu/wiki/bin/view/CICADA/NGNPP

¹⁵ <http://dspsr.sourceforge.net/>

¹⁶ <http://psrchive.sourceforge.net/>

¹⁷ The PSRCHIVE calibration process produced unphysical results for the earliest S-band detection (MJD 56424); we have calibrated it by using an approximation based on the measured S-band system equivalent flux density and the radiometer equation (cf. §7.3.2 of Lorimer & Kramer (2004)). The result is reasonable, given that the next S-band observation five days later has a comparable flux density (see Figure 6).

the narrow-band RFI, this left less than ~ 400 MHz of clean, usable bandwidth. Similarly, at S-band we had to remove the lower ~ 100 MHz and the upper ~ 25 MHz, and in total ~ 625 MHz of usable band remained¹⁸. Only 3% of the data was clipped from either end of X-band, with a total of 10% removed. We took these seemingly draconian measures to offset the original data quality and to ensure that the time- and frequency-averaged profiles were of reasonably high quality (e.g., see Figure 1). This was enabled by the source’s relatively large flux density.

The data quality situation at X-band was still more complicated. As also noted by Lynch et al. (2014) in their investigation of this magnetar, large gain variations on timescales from a fraction of a pulse period to several seconds (visible in the time-series data) are prevalent in X-band at the GBT, when pointed at the Galactic Center. The variations did not (necessarily) integrate away over hour-long observations and are representative of a stochastic red-noise process. We attribute these variations to changes in atmospheric opacity (Lynch et al. 2014) and/or small pointing errors, noting a strong resonance in the GBT X-band pointing very near 0.3 Hz ¹⁹. The gain variations would be manifested by the relatively small beam of X-band ($\sim 1.4'$, compared to $\sim 6'$ and $\sim 9'$ for S- and L-band) oscillating over the crowded, bright Galactic Center (the central parsec extends $\sim 0.4'$, and the separation of J1745–2900 from Sgr A* is only $\sim 0.04'$ (Rea et al. 2013)). Additionally, it is likely that the baseline variations are much less prominent at low frequencies because they act as “zero-DM” signals that get smeared out when the pulsar’s signal is dedispersed. Lynch et al. (2014) also state that the effect may be a function of elevation, which fits with our pointing-resonance hypothesis, since the influence of variable elements like the wind will be a function of elevation. The persistence and variability of these variations can be seen in Figure 2.

The analyses that follow utilized these folded profiles in a variety of reduced forms. Unless otherwise noted, the reduced radio data have 2048 profile bins (~ 7.2 ms per bin), 32 frequency channels (25 MHz per channel), and 5 min subintegrations; in this work, we only consider the total intensity profiles.

2.2. X-ray

During two of our radio epochs, MJD 56500 and MJD 56516, we obtained simultaneous observations of J1745–2900 with the *Chandra* X-ray Observatory (Obs. IDs 15041 & 15042; PI: D. Haggard). Table 2 contains details of the X-ray observations (for further details see Coti Zelati et al. 2015). The field of the first observation is shown in Figure 3; the second observation was essentially the same. In each observation, J1745–2900 was positioned on the back-illuminated chip S3 of the Advanced CCD Imaging Spectrometer (ACIS, Garmire et al. 2003) instrument. The data were reprocessed with the *Chan-*

¹⁸ In two epochs, however, instrument problems left only half of S-band viable. See Table 1.

¹⁹ Even though the average pointing errors at X-band are only on the order of several arcseconds at mid-elevations and mild wind conditions, the power spectrum in elevation offset shows resonances overlapping with the magnetar’s spin frequency ($0.27\dots \text{Hz}$). See http://www.gb.nrao.edu/~rmaddale/GBT/Commissioning/Pointing_Gregorian_HighFreq/PntStabilityXBand.pdf for details.

Table 1
Summary of GBT Observations

UTC Epoch	MJD	Bands Observed	Approx. Length [min]
2013-05-04	56416	X	20
2013-05-12	56424	S*,X	122,200
2013-05-13	56425	X	60
2013-05-14	56426	X	49
2013-05-17	56429	S,X	70,53
2013-05-23	56435	X	50
2013-05-30	56442	X	58
2013-06-21	56464	X	54
2013-07-14	56487	X	71
2013-07-15	56488	L,S	120,132
2013-07-27	56500	L,S,X	186,108,68
2013-07-28	56501	L,S*	133,117
2013-08-03	56507	L,S	112,75
2013-08-12	56516	L,S,X	120,60,56

Note. — The listed dates and MJDs for the epochs are representative of the majority of the epoch, not the start time; observations on the same day were taken in tandem. The two boldfaced epochs are those for which we have simultaneous observations with *Chandra*. The lower half (400 MHz) of the two S-band observations with an asterisk were corrupted and unusable. The horizontal lines separate the epochs during which the three observed types of X-band profile are seen (see §3.2 and Figure 2).

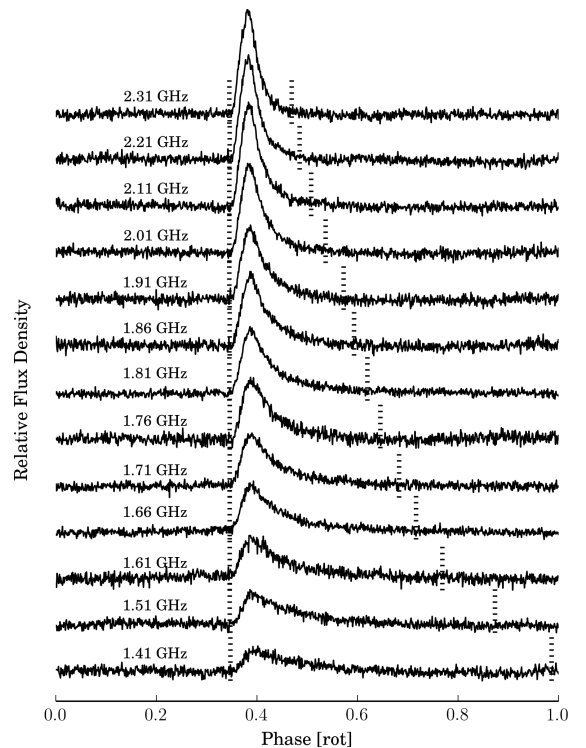


Figure 1. Examples of L- and S- band profiles averaged over all epochs. The profiles are shown with 1024 phase bins for clarity. These data are aligned via a wideband portrait model, as described in §3.5. In general, the un-averaged profiles were also of good quality, with only minor systematics in the baseline. The total bandwidth covered across these two bands is about 1 GHz, from ~ 1.4 to 2.4 GHz; 25 MHz of data were averaged for each of these profiles, with their center frequencies shown. The profiles were very well described by a single scattered Gaussian component, and so we do not over-plot the wideband model. The vertical dotted lines show examples of on-pulse regions used for the flux density measurements. See §3.4 for details.

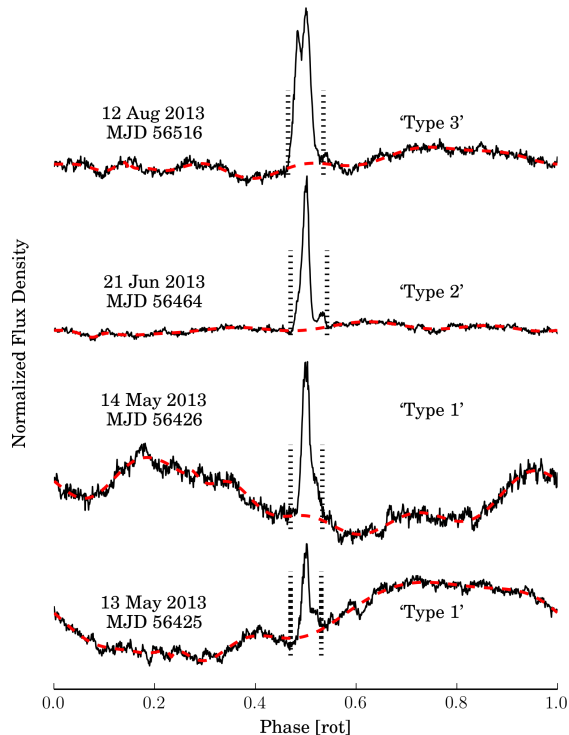


Figure 2. Examples of time- and frequency-averaged X-band profiles. The profiles are shown with 1024 phase bins for clarity. The baseline variations were removed on a profile-to-profile basis by fitting a high degree polynomial (red dashed lines) to the off-pulse region (outside the dotted lines) in order to make measurements of the flux density (see §3.4). The on-pulse phase window varied in size between about 6 and 8%. The profile evolved monotonically from one “type” to the next (see §3.2 and Table 1).

dra Interactive Analysis of Observations software package (CIAO, version 4.6, Fruscione et al. 2006) and the calibration files in the CALDB release 4.5.9.

In both observations, J1745–2900 was bright enough to cause pile-up in the ACIS detector. A “pile-up map” created with the CIAO tool `pileup_map` confirmed that mild pile-up was present. Exclusion of data near the center of the point-spread function (PSF) from the analysis would have resulted in the loss of too many photons (63% of the source counts were in the two central pixels). Moreover, the external part of the PSF contained a substantial number of counts from Sgr A*. We thus decided to proceed as follows.

We extracted the source counts from a circular region centred on J1745–2900 with a $1.5''$ radius (see Figure 3); this region includes the piled-up events. This area covers $\sim 85\%$ of the *Chandra* PSF (encircled energy fraction) at 4.5 keV. A larger radius of 2– $2.5''$ would let in more counts from Sgr A* and would only marginally increase the encircled energy fraction. Because of the complex environment, the background spectrum needed to be extracted close to the source. We used a thin annulus (with radii of $2''$ and $4''$), excluding a bright area associated with Sgr A*. The spectra, the ancillary response files and the spectral redistribution matrices were created using `specextract`. Following Rea et al. (2013), we adopt a pure blackbody for the spectral shape. We corrected the spectra using the pile-up model by Davis (2001), as implemented in the modeling and fitting package SHERPA (Freeman et al. 2001). The pile-up fraction,

Table 2
Summary of Simultaneous *Chandra* Observations

Obs. ID	Radio epoch [MJD]	Exposure time [ks]	Net source counts [10^3]	RMS pulsed fraction [%]
15041	56500	45.4	15.7	28.8 ± 1.5
15042	56516	45.7	14.4	28.9 ± 1.8

Note. — The 1σ uncertainties for the RMS pulsed fractions were determined from Monte Carlo simulations (cf. Gotthelf et al. (1999)). By another measure, the pulsed fractions — defined as the difference between the profile maximum and minimum divided by their sum — are $\sim 48\%$. The folded profiles are shown in Figure 5.

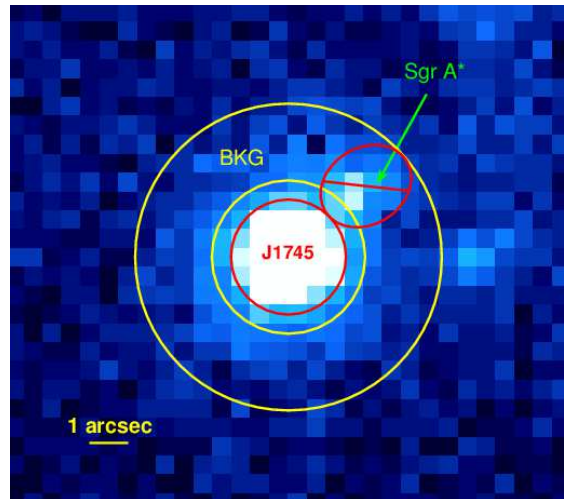


Figure 3. *Chandra* field of J1745–2900 for observation 15041. 1 ACIS pixel = $0.492''$. The source counts were taken from the central-most encircled region (red circle). Background counts were extracted from the annulus between the outer two (yellow) circles, excluding the area marked as “Sgr A*”. We account for pile-up as described in §2.2.

estimated by fitting the `jdpileup` model, is 3.7% for the first observation, and 4.1% for the second. We did not attempt any correction of the light curves; the pile-up fraction is modest and, in general, pile-up affects spectra more than it does light curves and pulse profiles²⁰. The spectral model fits were acceptable only when the pile-up model component was included. A summary of the spectral fits is given in Table 3.

3. RESULTS

3.1. Transient Events

J1745–2900 is known to show narrow individual pulses (Spitler et al. 2014; Bower et al. 2014; Lynch et al. 2014), similar to the radio magnetar J1622–4950 (Levin et al. 2012). We performed a cursory analysis of J1745–2900’s individual pulses in our X-band data, seeking only to find anomalous burst-like events in the radio data that might be coincident or correlated with X-ray features or flares. For this, we took two approaches. In the first case, we folded the raw data into single-rotation integrations, approximately maintaining the original temporal resolution, averaging over

²⁰ This is true unless the pulse profiles are strongly dependent on energy, which is not the case for J1745–2900, though we refer the reader to Coti Zelati et al. (2015) for further details.

Table 3
Chandra Spectral Results

Obs. ID	μ^a	f^a	N_{H} [10^{23} cm $^{-2}$]	kT^b [keV]	R^b [km]	Observed flux ^c [10^{-12} erg cm $^{-2}$ s $^{-1}$]	Luminosity ^c [10^{35} erg s $^{-1}$]	χ^2_{red} (dof)
15041	$0.50^{+0.31}_{-0.05}$	99.8%	1.26 ± 0.03	0.82 ± 0.01	$2.34^{+0.13}_{-0.17}$	8.9 ± 1.2	$3.2^{+0.4}_{-0.3}$	1.00 (288)
15042	$0.48^{+0.25}_{-0.08}$	97.1%	1.23 ± 0.03	0.83 ± 0.01	$2.16^{+0.13}_{-0.16}$	$8.1^{+1.4}_{-1.1}$	2.8 ± 0.4	1.00 (287)

Note. — The abundances used in the absorbed blackbody model are those of Anders & Grevesse (1989) and photoelectric absorption cross-sections are from Balucinska-Church & McCammon (1992). See Coti Zelati et al. (2015) for a complete treatment of these observations in the context of a long-term X-ray monitoring campaign. Parameter uncertainties in the table are 1σ .

^a Parameters of the `jdpileup` SHERPA pile-up model; μ is the grade-migration parameter and f is the fraction of the PSF treated for pile-up, required to be in the range 85–100%. For details, see Davis (2001) and “The *Chandra* ABC Guide to Pileup”.

^b The blackbody temperature and radius are calculated at infinity and assuming $D = 8.3$ kpc (Genzel et al. 2010), which is assumed throughout this work.

^c In the 0.3–8 keV energy range; for the luminosity we again assumed $D = 8.3$ kpc.

frequency, and summing the polarizations. These data were inspected visually. In the second case, we analyzed the raw data with the PRESTO²¹ pulsar software package. Here, we applied an RFI mask to the raw data with `rfifind`. We then made a dedispersed²², frequency-averaged time-series with `prepdata` for each X-band epoch, and searched for single pulses with the boxcar-convolution algorithm implemented in `single_pulse_search.py`. We repeated this process on the unmasked raw data.

Single pulses were detected; indeed, one to several pulses are visible by eye during almost every rotation. However, we saw no anomalously large single pulses or other bursts. The distributions of estimated single pulse energies all peak at $\lesssim 1$ times the average profile energy and were inconsistent with power-law distributions. The phases of the single pulse arrival times were consistent with occurring within the on-pulse window, and the distributions of the resolved single pulse widths peaked near 3–4 samples ≈ 2 ms, in agreement with the X-band scattering timescales found by Bower et al. (2014).

Similarly, the (unfolded) X-ray light curves during the two simultaneous observations, binned from 0.5 to 5000 s, were featureless and constant. The χ^2 probability of constancy was high for both observations, regardless of the choice of binning ($>30\%$ and frequently approaching 100%). Due to the uniformly poor quality of the X-band data as previously described, we refrain from further analysis or discussion of this aspect of J1745–2900 and direct the reader to the observations of its X-band single pulses as observed with the Very Large Array and the GBT in Bower et al. (2014) and Lynch et al. (2014), respectively.

3.2. Profile Variability

Figure 2 shows examples of the three general types of observed X-band profiles, as well as corresponding examples of our baseline removal and on-pulse determination. The transition between “Type 1”, with a single main component having a trailing-side shoulder and a more quickly rising leading edge, and “Type 2”, with the main component having a leading-side shoulder and a nub fea-

ture on the trailing side, happens more than three weeks after the X-ray burst on MJD 56407 and more than two weeks before the burst on MJD 56450. The “Type 1” shape was seen as early as a week after the discovery (Eatough et al. 2013a) and published in Eatough et al. (2013b). Similarly, the transition between “Type 2” and “Type 3”, which has a larger two-peaked component, happened more than two weeks after the burst on MJD 56450 and more than three weeks before the burst on MJD 56509. For these reasons, we do not associate the profile types (which are most likely not absolutely discretized) with the observed X-ray bursts. Within a single observation, the profile shape did not change between 5 min subintegrations.

Lynch et al. (2014) also documented the time-variability of J1745–2900’s X-band profile as seen with the GBT. As their first observation is coincident with our last observation, they have also seen the “Type 3” shape, which persists and evolves during most of what they have labeled a “stable state”. This “stable state” is characterized by relatively smooth profile transitions, a gradual flux evolution, and a phase-connected timing solution — all in contrast to what they call an “erratic state”, which is onset sometime after MJD 56682. Later in their observations, during epochs with MJDs 56794 and 56865 (both in the “erratic-state”), they see a profile resembling what we have labeled “Type 1”. We note that we did not witness any of the very sporadic profile variability seen in Lynch et al. (2014) associated with the “erratic state” (e.g., the drastic profile changes seen in their last two observations, separated by only eleven days), but rather we observed each of these three types only for a single interval of time.

3.3. Timing

Between having bursts, glitches, unstable profiles, and timing noise, magnetars are notoriously some of the hardest pulsars to time (cf. the original radio magnetar J1809–1943 (Camilo et al. 2006), or see a recent review of magnetars in Mereghetti (2013)). As is evident from the X-ray and radio timing in Coti Zelati et al. (2015), Lynch et al. (2014), and Kaspi et al. (2014), obtaining a single phase-connected timing solution for J1745–2900 is difficult, due to a significant level of timing noise. Here, we measure an overall average spin-down for the purpose of summing the data in each epoch.

Pulse times-of-arrival (TOAs) were measured by cross-correlating the time- and frequency-averaged data pro-

²¹ <http://www.cv.nrao.edu/~sransom/presto/>

²² We used the Eatough et al. (2013b) DM of 1778 cm $^{-3}$ pc, and compared the results to those from times-series dedispersed at 0 cm $^{-3}$ pc and twice the nominal DM in order to discriminate between transient RFI and candidate pulses.

files with smoothed, “noise-free” template profiles using standard PSRCHIVE routines. The templates are generated by arbitrarily aligning and averaging all of the data for which the template is used. Single templates were used for the L- and S-band data, but three separate templates were used for X-band, depending on the profile observed, as discussed in §3.2. Arbitrary phase offsets were fit between TOAs measured from all of the different templates as part of the timing models. These phase offsets serve to align the template profiles, but do so indiscriminately with respect to pulse broadening from interstellar scattering; this has the effect of biasing DM estimates if one tries also to measure the dispersive delay between TOAs of different frequencies. See §3.5.2 for our DM measurements based on wideband modeling of the L- and S-band data.

Figure 4 shows the measured values of the spin frequency f as a function of time. The average measured spin-down of $\dot{f}_{avg} = -8.3(2) \times 10^{-13} \text{ Hz s}^{-1}$ was sufficient to average the data in each epoch with negligible smearing for the flux measurements (§3.4), and is a reasonable approximation for the overall trend in the spin evolution²³. This average value also lies between the two \dot{f} values presented in Table 2 of Kaspi et al. (2014) for the same range of dates.

Although we are not interested in a full timing solution for these data in this work, we found corroborative results when following the suggestion in Kaspi et al. (2014) that there is an abrupt change in \dot{f} around the time of the *Swift*/BAT-observed X-ray burst on MJD 56450. Namely, while a single, predictive timing solution was not found, our pre- and post-burst TOAs are described by two simple phase-coherent solutions with parameters $\dot{f}_{pre} = -5.005(1) \times 10^{-13} \text{ Hz s}^{-1}$, $\dot{f}_{post} = -9.4799(5) \times 10^{-13} \text{ Hz s}^{-1}$, and $\ddot{f}_{post} = -2.696(6) \times 10^{-20} \text{ Hz s}^{-2}$. These values are in good agreement with those in Coti Zelati et al. (2015), Kaspi et al. (2014), and Rea et al. (2013), although we were not sensitive to \ddot{f}_{pre} . We could only obtain a single phase-connected timing solution for all of the TOAs by using five spin frequency derivatives, which is not a predictive ephemeris.

3.3.1. Profile Alignment

In Figure 5, we present the absolute alignment between the *Chandra* 0.3–8 keV X-ray profiles and the GBT radio profiles in L-, S-, and X-bands. We determined an independent ephemeris for each of the two epochs from the radio data by fitting TOAs from each day for only the spin frequency, fixing the spin-down parameter at the average value reported above. These TOAs were measured from the frequency-averaged data with 5 min subintegration resolution. The phase-zero time was referenced to the arrival of infinite-frequency radiation at the Solar System barycenter, which assumes a constant dispersion measure of $1778 \text{ cm}^{-3} \text{ pc}$ between the two observations. The X-ray photon arrival times were barycentered also using the sky position given in §2.1 and the JPL Planetary Ephemeris DE-405. These events were folded into pulse profiles with 64 phase bins using the corresponding epoch-specific ephemeris by the `prepfold` program of

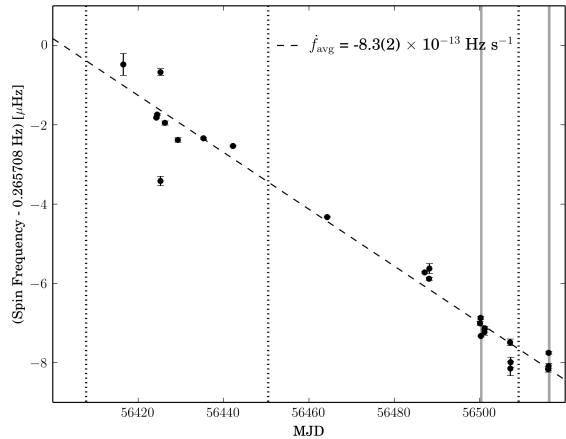


Figure 4. Average spin evolution of J1745–2900. The three vertical dotted lines correspond to the three X-ray bursts detected by *Swift*/BAT. The two vertical grey bars cover our *Chandra* observations. Measurements from the two early S-band observations are not included, nor from the X-band epoch on MJD 56425, as they were very significant outliers. The quoted uncertainty does not include the residual scatter.

PRESTO. The alignment based on folding using a single ephemeris for both epochs — either the post-burst or the multiple frequency derivative ephemeris — yielded indistinguishable results. This is reasonable, since the RMS timing residual from either of those ephemerides is on the level of individual bins. On the other hand, it may be surprising that there seemed to be no interruption in the “post-burst” ephemeris; the third detected X-ray burst occurred at the midpoint between the two simultaneous radio/X-ray epochs.

We modeled each of the two X-ray profiles with four Gaussian components to measure the relative offsets with respect to the radio profiles. The offsets and their uncertainties were determined from Monte Carlo trials, where “offset” here refers to the phase that maximizes a cross-correlation such as the one prescribed in Taylor (1992). There was a small offset between the X-ray models, $\lesssim 0.02$ rot. A difference in DM would shift the relative phase between the X-ray profile and the S-band profile (our fiducial profile) only by $\sim 3 \times 10^{-4}$ rot per unit DM [$\text{cm}^{-3} \text{ pc}$]. Even for the DM difference of $\sim 17 \text{ cm}^{-3} \text{ pc}$ measured between these epochs (see §3.5.2 and Figure 8), the phase difference is ~ 0.005 rot. The remaining offset can be explained by a combination of the variability of the X-ray profile and timing noise, with the former being dominant. After removing this difference, the offsets with respect to the radio profiles do not change between the two days within the variance of the measurements. The phase offset relative to the S-band profile is approximately $0.15(1)$ rot. The radio magnetars J1809–1943 and J1550–5418 both also show rough alignment of pulsed radio emission with their X-ray profiles (Camilo et al. 2007b; Halpern et al. 2008), whereas no pulsed X-ray emission has been detected from J1622–4950 (Anderson et al. 2012).

The two double-peaked X-ray profiles appear essentially featureless. The RMS pulsed fractions are given in Table 2. There are not sufficient data to decompose the profiles into energy bands to look for meaningful spectral dependencies, although we wish to point out a possible transient feature that appears in the *XMM-Newton*

²³ Quantities in parentheses represent the 1σ uncertainty on the last digit in the respective measurement throughout the paper.

data recently published by Coti Zelati et al. (2015). In the energy-dependent *XMM-Newton* profiles of Figure 4 from Coti Zelati et al. (2015), there is a conspicuous narrow feature on the leading edge of the double-humped X-ray profile that is close to the phase of radio emission (within ~ 0.05 rot). It appears most prominently around phase 0.55 in the 0.3–3.5 keV profile of the third *XMM-Newton* observation (with Obs. ID 0724210501). It is also seen in two of the other three energy-dependent profiles (except for the highest energy 6.5–10.0 keV profile), contributing to the integrated flux in the energy-averaged profile. A similar feature is seen at the same phase in the first *XMM-Newton* observation (with Obs. ID 0724210201) to a lesser extent. According to their table, these observations were separated by 23 days, with the first occurring 19 days after the *Chandra* observations presented here (which are also included in Coti Zelati et al. (2015)). The three *Chandra* observations and the one *XMM-Newton* observation taken during these 23 days show no obvious feature, despite covering the same range of energies, although *Chandra* recorded only between 10 and 50% of the counts as did by *XMM-Newton*. Therefore, without additional observations, it remains only a peculiarity.

3.4. Radio Flux Density

From the radio data, we made measurements of J1745–2900’s flux density as a function of time and frequency. We measured the mean flux densities in 50 MHz wide channels and used a weighted average of these measurements to obtain representative flux densities for each band, per epoch.

For all of the L- and S-band profiles, we defined “on-pulse” regions as follows. A model pulse profile for each frequency was determined from the wideband modeling described in §3.5. We then found the smallest range of pulse phases that contained 99% of the integrated flux density of the model profile. Examples of the on-pulse windows for the scattered L- and S-band profiles can be seen in Figure 1. The mean flux density was calculated by averaging the observed flux density in the window and scaling it by the duty cycle. The uncertainties were estimated by measuring the mean noise level in the last quarter of each profile’s power spectrum²⁴. We accounted for systematics in the residual profile by adding the scaled, residual mean flux density to the uncertainty in quadrature. These corrections were small, as the reduced χ^2 values of the residuals were usually < 1.5 and always < 2.0 .

The measurement of the X-band flux densities was complicated by the dynamic baseline variations mentioned in §2.1, as well as the intrinsic variability of the profile shape. We used polynomial functions to remove the baseline variations on a profile-to-profile basis (e.g., see Figure 2). For these profiles, we first centered each profile to be near phase 0.5 to avoid edge-effects of the polynomial fit from affecting the on-pulse region. A high degree polynomial function was fit to the baseline of each profile, where in the first iteration an on-pulse window with a duty cycle of 6% was blanked out from the fit to avoid initially over-estimating the noise²⁵. The level

of the residual off-pulse noise was calculated, and then the on-pulse window was widened until the flux density at the edges of the on-pulse region dropped below the noise level. The baseline polynomial was then refit to the original profile, but with the new on-pulse window blanked out. The mean flux density and its uncertainty were calculated in these baseline-removed, on-pulse windows as described for the lower frequency data above, but a systematic error was added in quadrature to the uncertainty that represented the mean flux density across the on-pulse phase window removed by the polynomial fit. This tested method gives dependable, conservatively estimated X-band flux densities.

3.4.1. Flux Evolution

The radio flux evolution of J1745–2900 is shown in Figure 6. The mean X-band flux density increases rapidly in the first half of our observations, increasing by at least a factor of ~ 6 over fifty days, and then tapers off at the 1 mJy level. The earliest reported measurement of J1745–2900’s X-band flux density was ~ 0.2 mJy, taken with the Effelsberg 100-m Radio Telescope, consistent with our GBT measurement two days later (Eatough et al. 2013a). Our data show a similar increase in the low frequency flux densities. The S-band flux increases by about an order of magnitude over ninety days, and in our last five observations covering about thirty days, the average L- and S-band fluxes increase by a factor of two. Given the measured scattering timescales for J1745–2900 (see §3.5.1) and the recently measured proper motion of the pulsar, the timescale for refractive scintillation to be important is much larger than the span of our observations (see Bower et al. (2015) for further discussion).

Having picked up where we left off, Lynch et al. (2014) increased the cadence of GBT X-band observations after MJD 56516 and found a similar, slow increase of the flux, up to ~ 3 mJy, over the next 170 days. As already mentioned, after this “stable state” of slow, steady flux increase, the authors found that J1745–2900 entered an “erratic state”, characterized in part by a larger and highly variable X-band flux, similar to what was seen in two other radio magnetars (Camilo et al. 2007b; Levin et al. 2012). Superimposed on top of this radio flux evolution is a relatively slow decay of the X-ray flux, compared to other magnetars (Rea et al. 2013; Kaspi et al. 2014; Lynch et al. 2014; Coti Zelati et al. 2015). Between our two simultaneous GBT/*Chandra* observations separated by ~ 15 days, the radio flux increased by $\sim 60\%$ while the X-ray flux decreased by $\sim 10\%$. This trend (seen here and in Lynch et al. (2014)) is opposite to those of the other radio magnetars, which show decreasing radio and X-ray flux with time over the course of an outburst Rea et al. (2012).

3.4.2. Radio Spectral Index

Because we have essentially simultaneous observations²⁶ of J1745–2900 in frequency bands spaced by two

a polynomial of degree 15 was used; this was the smallest degree polynomial that reasonably and automatically removed systematic baseline trends from all of the profiles without having to also vary the degree of the polynomial on a profile-to-profile basis.

²⁶ In one case, the X-band observation was taken a day earlier; see Table 1.

²⁴ This is a robust method to estimate the off-pulse variance, assuming the profile is resolved (e.g., see Demorest 2007).

²⁵ None of the profiles had a smaller duty cycle than 6% and

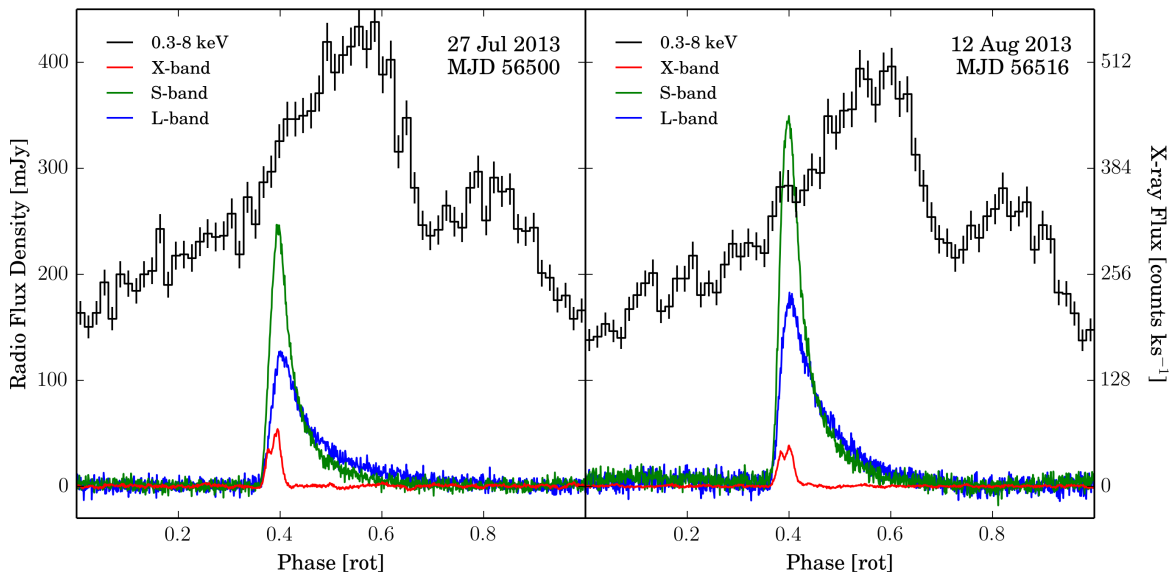


Figure 5. Absolute phase alignment of J1745–2900’s radio and X-ray profiles determined separately on two days. Note that the brightest radio profile is seen in S-band (see Figure 7). The profiles have 1024 and 64 phase bins, respectively, and are shown as they would be observed at the Solar System barycenter for phase-zero MJDs 56499.98000761 and 56515.96999979, referenced to infinite frequency. The assumed dispersion measure is $1778 \text{ cm}^{-3} \text{ pc}$. During two later *XMM-Newton* observations (presented in Coti Zelati et al. (2015)), there is a peculiar, narrow feature seen in the otherwise broad X-ray profile near the phase of radio emission as shown here (see text).

octaves, we can measure the spectral index γ , where $S_\nu \propto \nu^\gamma$ for flux density S_ν at frequency ν . The upper panel of Figure 6 shows γ as measured between the average X-band flux density and the combined average flux densities of the lower frequency band(s). The error bars were approximated by varying the average fluxes within their measurement uncertainties. The decorrelation bandwidth for diffractive scintillation is much smaller than even our native frequency resolution and will not be a source of variability here.

There is no large, obvious stochasticity, as opposed to, for example, J1809–1943 (Lazaridis et al. 2008), but there may be a trend. Shannon & Johnston (2013) report two early measurements of γ across the bands spanning 4.5–8.5 GHz and 16–20 GHz. The first measurement on MJD 56413 is close to -1.0 in the high frequency band, though it is closer to 0.0 in the lower frequencies, and the second on MJD 56443 is ~ -1.0 across both bands, consistent with our measurements more than two weeks prior. Our three later measurements indicate a significantly steeper spectrum. The average value for γ of -1.4 is tantamount to the average spectral index for normal pulsars across gigahertz frequencies as reported in Bates et al. (2013). Camilo et al. (2007c) and Anderson et al. (2012) both make mention of a general steepening of the spectral indices of J1809–1943 and J1622–4950, respectively, despite remaining much flatter than what is seen in J1745–2900. However, (Lazaridis et al. 2008) finds the opposite for J1809–1943 in later observations.

This finding apparently breaks the mold set by the other three radio magnetars, which have essentially flat (or inverted) spectra (Camilo et al. 2006, 2008; Levin et al. 2010; Keith et al. 2011). However, no firm conclusions can be drawn from this handful of measurements from early times in J1745–2900’s outburst, especially knowing that the other radio magnetars also show a variable radio spectrum (Camilo et al. 2007c;

Lazaridis et al. 2008; Anderson et al. 2012). In fact, at the time of writing, the findings of Torne et al. (2015) suggest that at much later times (a year after the present observations), the radio spectrum of J1745–2900 between 2 and 200 GHz was much flatter, with $\gamma = -0.4(1)$.

3.4.3. Spectral Shape

One example of J1745–2900’s radio spectrum is shown in Figure 7; the spectra from the other days are qualitatively similar. The spectrum shows a non-power-law increase in flux between 1.4 and 2.4 GHz, with a possible peak near 2 GHz. The inverted log-parabolic shape is reminiscent of what have been called “gigahertz-peaked spectra” (GPS) pulsars (Kijak et al. 2011, 2013; Dembska et al. 2014, 2015), although the GPS pulsars supposedly have a much broader spectral shape, over a dex in frequency. For reference, we fit a log-parabola to the low frequency points, the parameters of which are given in the figure.

It is difficult to explain the spectral shape we see in the lower frequencies. It is conceivable that the dense, unique environment near J1745–2900 in the Galactic Center significantly alters the spectral shape of radio emission between 1 and 10 GHz (e.g., via free-free absorption, although the detection of Sgr A* at 330 MHz implies a low free-free optical depth of $\lesssim 1$ (Nord et al. 2004)), but it is difficult to draw any conclusions without a dedicated set of observations.

Another possibility is that we have systematically under-estimated the flux: one well known source of bias comes from under-estimating the flux at low frequencies due to significant area in the scattering tails being lost in the calculation of the baseline flux. However, even at 1.4 GHz the scattering timescale is $\sim 500 \text{ ms} \approx 0.13 \text{ rot}$ (see §3.5.1). In the worst case of a Kolmogorov scattering index (-4.4), the scattering timescale at our lowest frequency is no more than $\sim 20\%$ of a rotation. As mentioned in Kijak et al. (2011) and treated graphically in

Macquart et al. (2010), the pulsed fraction drops by only $\sim 10\%$ when the scattering timescale is *half* the pulse period. Therefore, we can suggest that at worst we are underestimating the L-band flux densities at the $\sim 10\%$ level, but this still would imply a positive or approximately flat spectral index between 1.4 and 2.4 GHz; the observed flux density drops precipitously somewhere thereafter.

A more promising, albeit provisional possibility has been offered up by recent modeling of the Shannon & Johnston (2013) observations. Lewandowski et al. (2015b) make a case study of J1745–2900 to demonstrate the possibility of thermal free-free absorption as the explanation for the GPS. For J1745–2900, the authors suggest a combination of an expanding ejecta and/or an external absorber to explain the changing spectrum seen early after the initial outburst in Shannon & Johnston (2013). The free-free absorbed model spectra offer a reasonable explanation for the lack of low-frequency detections of J1745–2900 immediately after the initial outburst and detections above 4 GHz; our two early S-band observations may support this idea. Our spectra from three months later may also inform the story of an evolving or endemic free-free absorbing medium in the environment of J1745–2900.

3.5. Wideband Portrait Model

As is evident from Figure 1, J1745–2900 has a highly scattered, simple profile across a gigahertz bandwidth, from 1.4 to 2.4 GHz. For a nominal DM value of $1778 \text{ cm}^{-3} \text{ pc}$, there is a delay of ~ 0.66 rotations across this band, which is easily measurable. All of the average L- and S-band profiles showed prominent scattering tails from multipath propagation through the interstellar medium (ISM). The quality of the data permitted us to make “wideband” measurements of both the DM and the scattering timescale τ , as well as its power-law index α , on an epoch-to-epoch basis²⁷.

For this, we used the methods and augmented software described in Pennucci et al. (2014) to make a wideband “portrait”²⁸ model for each of the five epochs where we have both L- and S-band observations. For each of these epochs we combined the data from the two low frequency bands in a fit for a global portrait model that included a single scattered Gaussian component with profile evolution parameters, a constant baseline term, a phase offset between the bands, DMs for each band, and the scattering index α . The scattering timescale is defined in the usual way by assuming a one-sided exponential pulse broadening function for the ISM, so that an observed profile $p(\varphi)$ is the convolution given by

$$p(\varphi) = g(\varphi) * e^{-\frac{\varphi P_s}{\tau}} H(\varphi), \quad (1)$$

where φ is the rotational phase, P_s is the spin period, H is the Heaviside step function and $g(\varphi)$ is the intrinsic

total intensity profile shape. For a power-law spectrum of density inhomogeneities in the ionized ISM τ is expected to have a power-law dependence on frequency ν as

$$\tau(\nu) = \tau_{\nu_0} \left(\frac{\nu}{\nu_0} \right)^\alpha, \quad (2)$$

with reference frequency ν_0 . In all cases, the scattering timescale (133 ms at 2 GHz; see below) dominates the smearing from the process of incoherent dedispersion ($\sim 0.7 \text{ ms } (\nu/2 \text{ GHz})^{-3}$), the smearing from an incorrect DM when averaging channels ($\sim 25 \mu\text{s } (\delta\text{DM}/\text{cm}^{-3} \text{ pc}) (\nu/2 \text{ GHz})^{-3}$), and the temporal resolution (1.8 ms for 2048 profile bins), so we have not included those modifications of the pulse profile shape in the model. However, deviation from the simple timing models discussed in §3.3 (e.g., see Figure 4) during any of these epochs could add profile smearing in the integrated profiles (at the level of \sim tens of ms — a significant fraction of the scattering timescale). We avoided this source of bias by iterating over the timing model to remove the timing residual on a per-epoch basis.

We model g with a frequency-dependent Gaussian function,

$$g(\nu, \varphi) = A(\nu) \exp\left(-4\ln(2) \frac{(\varphi - \varphi_g(\nu))^2}{\sigma(\nu)^2}\right), \quad (3)$$

which is parameterized by its location $\varphi_g(\nu)$, full-width-at-half-maximum (FWHM) $\sigma(\nu)$, and amplitude $A(\nu)$.

As described in Pennucci et al. (2014), each of these parameters nominally has an additional parameter describing its frequency dependence. However, because this combined band has a fractional bandwidth of “only” ~ 0.5 , we assume $\varphi_g(\nu) = \varphi_0$ is a frequency-independent value. That is, we assume there is no drift intrinsic to the one component across the band.

Furthermore, when allowing for a frequency-dependent σ , we found no significant evolution, and so we chose also to fix the evolution $\sigma(\nu) = \sigma_0$ to be a frequency-independent fit parameter in our final portrait models. This choice was further justified by performing independent per-channel profile fits of a single, scattered Gaussian component and examining the frequency evolution of σ . Also, there are X-band observations for three of these epochs, and in these cases the FWHM of the X-band profiles (all of “Type 3”), when fitted with a single, unscattered Gaussian component, was always within the scatter of those measured from the lower frequency observations. These results are consistent with the weak (or lack of) frequency dependence of σ found in Spitler et al. (2014).

We normalized the intensities of the data to be fit by the maximum profile value in each frequency channel to remove the unusual spectral shape (see §3.4.3). This allowed the Gaussian amplitude to be easily modeled by a power-law function for $A(\nu)$. In all cases, the reduced χ^2 of the fit was < 1.1 , and a second Gaussian component was never justified by the residuals.

The combination of the quality of the X-band data, the variability of the profile, and the expected value of τ at 8.9 GHz ($\lesssim 1$ ms, comparable with our native time resolution) was such that we did not attempt to incorporate this high frequency data into our wideband profile model, nor did we measure the scattering timescale in either the

²⁷ The two earliest S-band observations were exceptions; corrupted data, low signal-to-noise ratios, and the lack of L-band data resulted in uninformative measurements of the DM, τ , and α . This is also the reason these observations were excluded from the average f measurement in §3.3.

²⁸ We use the word “portrait” to mean the total intensity profile as a function of frequency.

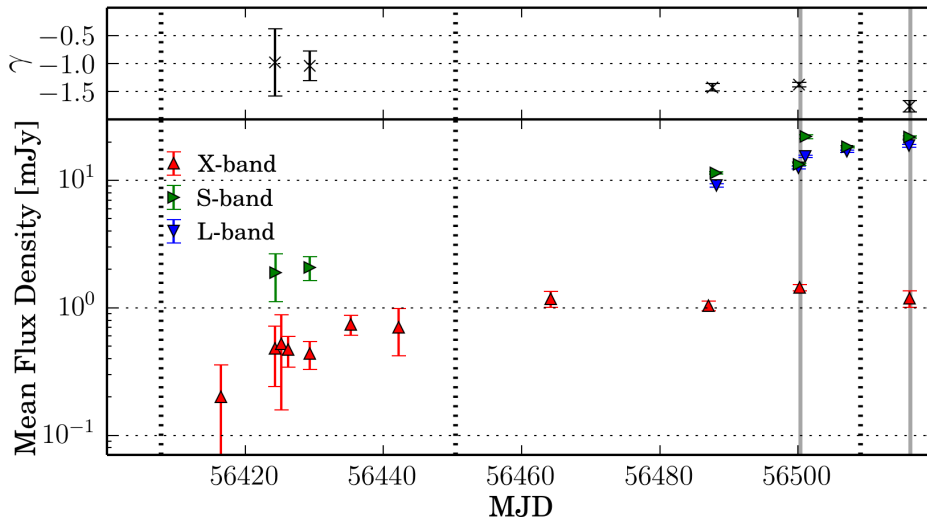


Figure 6. The early radio flux (bottom panel) and spectral (top panel) evolution of J1745–2900 over 100 days from the observations in Table 1. The vertical demarcations are the same as in Figure 4. Lynch et al. (2014) find a continuation of the slow, steady increase in X-band flux for another six months, which is followed by what they call an “erratic state”. The apparent excess average S-band flux density during MJD 56501 is explained by the fact that the lower half of the band was corrupted (see Table 1), and the pulsar’s flux density apparently increases with frequency in this range (see Figure 7). The average value of the spectral index γ is about -1.4 ; see text for details.

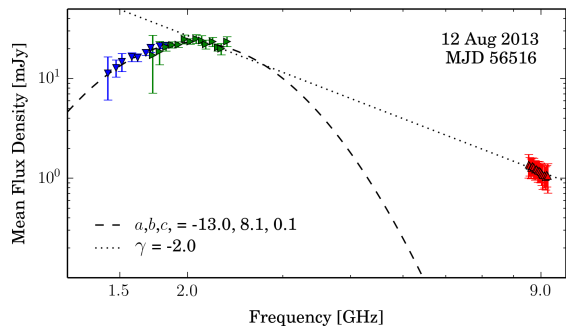


Figure 7. An example of J1745–2900’s radio spectrum from the brightest observed epoch, MJD 56516. The markers are as in Figure 6; note that the flux densities agree in the ~ 100 MHz overlap between L- and S- band. A similar inverted parabolic shape over log-frequency is seen during the other sets of (nearly) simultaneous observations, which is reminiscent of the so-called GPS pulsars. The coefficients a , b , and c of the fitted dashed parabola ($\log_{10}(S_\nu) = ax^2 + bx + c$, for $x = \log_{10}(\nu)$) are given in the plot, along with the spectral index γ , which for this plot was fitted between the peak of the parabola and the X-band data.

average profile or the single pulses. We refer the reader to Bower et al. (2014) and Spitler et al. (2014) for high frequency scattering measurements of J1745–2900.

3.5.1. Pulse Width & Scattering Parameters

The results from our wideband models are shown in Figure 8. There was no significant change in the measured FWHM of the unscattered profile, and our average (frequency-independent) value of $91.9(4)$ ms = $0.0244(1)$ rot is also consistent with what is reported in Spitler et al. (2014). The scattering timescale at 2 GHz, $\tau_{2\text{GHz}}$, appears to increase by $\sim 10\%$ over the four weeks, and the scattering index α deviates from its average value, first to a Kolmogorov value near -4.4 (the dash-dotted line), and then to a much shallower value near -3.0 . Both of these results are somewhat peculiar, but similar variations are also reported in Spitler et al. (2014), though they do not discuss the temporal evolution of either quantity. That is, their pub-

lished values of τ from a variety of epochs and frequencies cannot be unified by a single scattering timescale and index. In fact, their measurements of τ show more scatter over the course of their observations than those presented here, which have some overlap. When the authors combine all of their measurements, they find an average value for α of $-3.8(2)$.

We checked our measurements in two ways. First, we performed conventional profile fits of a single, scattered Gaussian component to each individual frequency channel, independent of any evolutionary constraint. The values of σ , τ , and α for each epoch were consistent with what we found by applying the wideband modeling method. In Figure 9, we show the measurements of τ measured in this way for the brightest observed epoch (MJD 56516) and over-plot the fitted power-law, which has the most extreme α value of the five epochs. There was nothing unusual about the data from this epoch in terms of RFI, data removal, calibration, or baseline variations. Second, as a check for our average values, we summed all of the data portraits together by coherently stacking the observations (having fit for a phase and DM in each epoch), and fit a single wideband model to the averaged data (with the same constraints as earlier). Using this method, we obtained similar average values: $\sigma = 0.0246(1)$ rot, $\tau_{2\text{GHz}} = 133.0(5)$ ms, and $\alpha = -3.71(2)$, the latter of which is in concert with the average α value from Spitler et al. (2014). Our extrapolated value of $\tau_{1\text{GHz}} = 1.74(3)$ s is only slightly at odds with their average value of $\tau_{1\text{GHz}} = 1.3(2)$ s, which is probably due to the temporal variability of τ . As others have noted (e.g., Bower et al. (2014)), the anticipated value for $\tau_{1\text{GHz}}$ along this line of sight based on empirical relationships, for a DM of 1778 cm $^{-3}$ pc, is about $600\times$ larger than what is observed (Krishnakumar et al. 2015; Lewandowski et al. 2015a).

Bower et al. (2014) determined $\tau_{8.7\text{GHz}} \lesssim 2$ ms from interferometric measurements of J1745–2900’s single pulses, implying that a value for α as shallow as -3

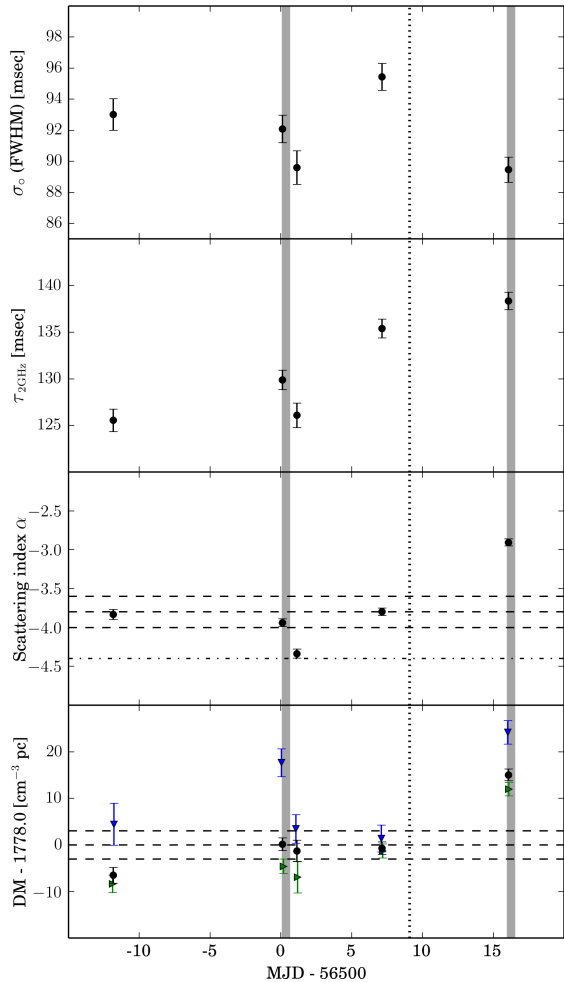


Figure 8. Our wideband measurements from five epochs. The vertical demarcations are the same as in Figure 4. The FWHM showed neither frequency nor temporal dependence. The trend in the scattering timescale τ is less scattered and more precise than the measurements presented in Spitler et al. (2014). The dash-dotted line in the panel for the scattering index α marks the fiducial Kolmogorov value of -4.4 , and the dashed lines mark the Spitler et al. (2014) measurement of $-3.8(2)$. The additional markers in the bottom panel are the same as in Figure 6: the blue/down-pointing triangles are L-band measurements, and the green/right-pointing triangles are S-band measurements — the dots are their weighted average. The dashed lines here are the Eatough et al. (2013b) DM of $1778(3) \text{ cm}^{-3} \text{ pc}$. See the text for a discussion of the DM measurements.

is not unbelievable. Furthermore, scattering measurements from two high-DM pulsars discovered near the Galactic Center (both within 0.3° and having DMs $1100\text{--}1200 \text{ cm}^{-3} \text{ pc}$) implied $\alpha = -3.0(3)$ (Johnston et al. 2006). It is not uncommon for pulsars to have $\alpha > -4$, particularly along special lines-of-sight, and it is empirically suggested that the highest DM pulsars may have an average scattering index significantly shallower than -4 (Löhmer et al. 2001, 2004; Lewandowski et al. 2015a). Note that observing $\alpha \neq -4.4$ does not necessarily imply a non-Kolmogorov spectrum of density inhomogeneities; rather, it could be that a non-thin-screen geometry may be responsible (Cordes & Lazio 2001; Lewandowski et al. 2013).

3.5.2. Dispersion Measures

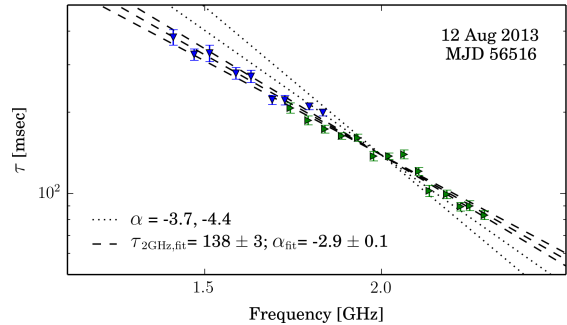


Figure 9. Independent per-channel measurements of the scattering timescale τ and the fitted scattering index α for the brightest set of L- and S-band observations, on MJD 56516; similar plots from the other four days have a significantly more negative slope. The dashed lines represent our measurement, whereas the dotted lines show our average value of α from wideband modeling and the fiducial Kolmogorov value of -4.4 , all with the same value of $\tau_{2\text{GHz}}$. Here, the measurement uncertainties have been inflated by the reduced $\chi^2 \sim 2$.

The bottom panel in Figure 8 shows the best-fit DMs as determined by the wideband models in the essentially simultaneous L-band and S-band observations (blue/down-pointing and green/right-pointing triangles, respectively). The black points are the weighted average of the two measurements; there is obviously some variance about the nominal value of $1778(3) \text{ cm}^{-3} \text{ pc}$, and our overall average value is $\sim 1781(1) \text{ cm}^{-3} \text{ pc}$. Without exception, the measured L-band DMs are greater than those measured in S-band. There is also only one epoch where the 1σ uncertainties have any overlap; the RMS variance of the differences is $\sim 6 \text{ cm}^{-3} \text{ pc}$. The absolute DM differences cause residual dispersion on the order of $\lesssim 20 \text{ ms} \sim 10$ bins (for 2048-bin profiles) across the corresponding band, and so they present significant profile deviations. To make sense of the discrepant DMs between the two frequency bands, we consider a number of possibilities.

First, the time-averaged data for each epoch showed few systematics with negligible baseline variations, so we do not believe that data quality was an issue here.

Next, as is well known, the measured absolute DM will be affected by the choice of profile alignment²⁹. We can rule out any simple, constant profile evolution as the source of the differing DMs because such a modification introduces a *constant* difference in the DMs; the changes in the measured DM should be the same independent of the choice of alignment. Even if our assumption that there is no intrinsic drift in the location of the (unscattered) profile component across the band is wrong, allowing for a drifting component will still reproduce discrepant DMs; we have verified this by allowing for frequency evolution in the location parameter of the Gaussian component.

A second confounding element from our modeling could be the use of different models for each epoch; if they are all systematically wrong in their alignments or

²⁹ For example, DMs are significantly biased when either assuming a constant profile shape in the presence of scattering, or aligning scattered profiles by conventional methods because the convolution of the ISM pulse broadening function with a profile of finite width introduces a delay that is a function of the scattering timescale (i.e. frequency). This is partly why proper wideband modeling is necessary.

representation, they could be wrong differently. One way to check this is to simply use one fixed model to remake the DM measurements. We used the average portrait model discussed earlier and confirmed that the DMs remain similarly extreme, within $\sim 2 \text{ cm}^{-3} \text{ pc}$, comparable with the measurement uncertainties. In fact, we tried a large number of fixed and variable portrait models, but never obtained either consistent DMs or DMs with a near constant offset. So, to the extent that τ and/or α are measurably changing, we are justified in keeping them as free parameters for each epoch’s model.

Similarly, the known profile variability that is seen in all of the radio magnetars could also play a role when using either a fixed or variable portrait model. However, besides the flux density, any underlying profile shape changes either with time or frequency are masked by the large level of scattering. As mentioned, the FWHM does not seem to change significantly in either time or frequency. Furthermore, the three X-band observations taken during these epochs show no large profile changes, and are all of the “Type 3” shape.

Next we can ask whether or not the slight asynchronicity could have any effect; that is, could the DM change so significantly on \sim hour timescales? We will return to this question below, but it is not an uncommon *a priori* assumption to expect that the observed DM does not change between observations separated by $\lesssim 4 \text{ hr}$.

One could also ask if the method by which we measure the DMs introduces a systematic error, where the error may depend on the exact values of τ and α , or even the spectral shape. To answer this, we performed a number of Monte Carlo simulations. In the simulations, we used the models from MJDs 56500 and 56516, which have the most extreme values for the difference between the DMs, and the most extreme values for τ and α , respectively. For each trial, we made fake L- and S-band observations by appropriately constructing the model for that band and scaling each frequency channel’s amplitude to match the spectral shape. We then added random frequency-dependent white noise to the model at the same level as measured from the data portraits and finally dispersed the fake data with a DM of $1778 \text{ cm}^{-3} \text{ pc}$. Visual inspection verified that the fake data were faithfully rendered. We used the same method to measure the DM (and phase), which is described in detail in Pennucci et al. (2014). In summary, the measured DMs were always in accord and unbiased, and the uncertainties were accurately estimated. We conclude that the measurement method produces accurate DMs, independent of the model parameters, provided that the model for the data is accurate.

We assume in our measurements that the phase offsets ($\Delta\phi$) incurred by finite-frequency signals due to propagation through the ionized interstellar medium scale as predicted by the usual cold-plasma dispersion law such that $\Delta\phi \propto \frac{\text{DM}}{F_s} \nu^{-2}$. This is certainly the case to first-order even over large, low frequency bandwidths (Hassall et al. 2012). However, to the extent that we understand the ISM to be *inhomogeneous* — after all, we do observe pulse broadening — then it is anticipated that the simple ν^{-2} dependence will be an insufficient description at some level for broadband DM measurements. When an inhomogeneous medium causes multi-path propagation

of radio waves where the path depends on frequency, the sampled column density of free electrons (the DM) will also be a function of frequency. Thus, we are left with the intriguing possible explanation that the DM inconsistencies we are seeing are the consequence of imposing a ν^{-2} dispersion law onto a frequency-dependent DM ($\text{DM}(\nu)$) due to an inhomogeneous ISM³⁰.

To our knowledge, the most recent claim for having observed frequency-dependent DMs was reported in Ahuja et al. (2007) for the slow, low DM pulsars B0329+54 and B1642–03, although they observed lower DMs at lower frequencies. However, the authors only made one set of simultaneous pairs of dual-frequency measurements per pulsar. We argue that to confidently segregate the effects of profile evolution, DM variations with time ($\text{DM}(t)$), $\text{DM}(\nu)$, and other potential confounding factors, many epochs of simultaneous, wide-band (large fractional bandwidth) observations of a stable, preferably high DM pulsar need to be made. A similar recommendation was recently made by Cordes et al. (2015) in their detailed study of frequency-dependent DMs, which makes theoretical predictions for the characteristic timescales and sizes of $\text{DM}(\nu)$ effects.

In their treatment of the problem, Cordes et al. (2015) predict the minimum scale of DM variations about a mean value,

$$\overline{\text{DM}}_{\text{rms}} \sim \phi_F^2 / \lambda r_e \sim 3.84 \times 10^{-8} \text{ cm}^{-3} \text{ pc } \nu_{\text{GHz}} \phi_F^2, \quad (4)$$

where ν_{GHz} is the frequency in GHz and ϕ_F is the size of the phase perturbations over the Fresnel scale, $l_F = \sqrt{(cD)/(2\pi\nu)}$, for the speed of light c and source distance D . For J1745–2900, which is in the strong scattering regime, ϕ_F will be very large. We estimate it from their prescription,

$$\phi_F(\nu) \approx 9.6 \text{ rad} \left(\frac{\nu/\Delta\nu_d}{100} \right)^{5/12}, \quad (5)$$

where $\Delta\nu_d$ is the scintillation bandwidth, which is readily estimated from our scattering measurements as $\sim 1.16/(2\pi\tau(\nu))$. For 1.4 and 2.4 GHz, we find $\overline{\text{DM}}_{\text{rms}} \sim 10$ and $5 \text{ cm}^{-3} \text{ pc}$, respectively. These can be compared to the RMS DM values as measured in L- and S-band of ~ 9 and $7 \text{ cm}^{-3} \text{ pc}$, respectively. The characteristic spatial size for the DM differences near 2 GHz will be several Fresnel scales, which can be converted to a characteristic time by using the recently measured proper motion of 236 km s^{-1} (Bower et al. 2015). For our range of frequencies, the characteristic timescale associated with the Fresnel scale size is $\sim 3 \text{ hr}$, comparable to the separation between the observations on a given epoch. Therefore, it may be that the small temporal gap between the observations contributes somewhat to the difference in the DMs, but we certainly do expect that the DMs vary significantly on different days, separated by many Fresnel timescales.

Finally, Cordes et al. (2015) make a prediction for the observed RMS difference between DMs at frequencies ν

³⁰ This is opposed to other supposed origins of $\text{DM}(\nu)$ relating to magnetospheric propagation effects or magnetic sweepback, which would likely have different statistics from an ISM induced $\text{DM}(\nu)$; see Hassall et al. (2012) or Ahuja et al. (2007) for an overview.

and ν' ,

$$\sigma_{\text{DM}}(\nu, \nu') \approx 4.42 \times 10^{-5} \text{ cm}^{-3} \text{ pc } F_{\beta}(r) \left(\frac{\nu \phi_F^2}{1000} \right), \quad (6)$$

where we have ignored a geometric factor of order unity and the function F_{β} contains the frequency dependence for $r \equiv \nu/\nu'$, given the power-law index β for the wavenumber spectrum of density inhomogeneities. For $\nu = 2.4$ GHz and $\nu' = 1.4$ GHz, $\sigma_{\text{DM}} \sim 4 \text{ cm}^{-3} \text{ pc}$, compared to our observed RMS difference of $\sim 6 \text{ cm}^{-3} \text{ pc}$.

That the predicted and observed values are similar may be coincidence, but we note the corroborating facts that J1745–2900 is the highest DM pulsar, is relatively bright, highly scattered, has a simple, easily modeled profile, and does not show significant profile evolution or stochastic profile variability (at least in these observations). Furthermore, we verified that our measurement method produces inconsistent (and biased) DMs between the bands by introducing non- ν^{-2} phase delays into our fake data simulations described earlier. After ruling out the other potential sources for the inconsistent DMs, we suggest that J1745–2900 may have an observable frequency-dependent dispersion measure.

A potential counter argument is that over many Fresnel timescales, one expects the sign of the DM differences to change, such that the observed low frequency DM becomes smaller than the high frequency DM. Between the small number and low density of epochs, the potentially incorrect portrait model, and ISM uncertainties (the predictions here are based on a thin-screen model with a Kolmogorov spectrum of density perturbations, which is partly supported by the findings in Bower et al. (2014)), it is conceivable that this observation is not inconsistent with a frequency-dependent DM as described.

Determining whether or not a difference in DM as seen in two frequency bands is intrinsically a $\text{DM}(\nu)$ effect is complicated by the issues described above, and with only five measurements we obviously cannot draw any definite or statistical conclusions, but future studies could potentially disentangle the evolution of $\text{DM}(t, \nu)$, the profile, and other ISM parameters. One strategy, as Cordes et al. (2015) note, is to model the frequency dependence of the dispersive delays as something other than ν^{-2} . This should be done for many epochs, at least as long as the timescale for refractive scintillation, over which time the specific frequency dependence of the average DM remains stable. For J1745–2900, this timescale is potentially many years.

4. SUMMARY & DISCUSSION

In this paper we have presented multi-epoch, multi-frequency wideband GBT observations of the Galactic Center radio magnetar J1745–2900 at 1.5, 2.0, and 8.9 GHz from the first ~ 100 days after it was discovered. After its initial X-ray burst on 25 April 2013, J1745–2900 underwent two additional bursts in the course of our observations. For two epochs, during which time we collected data from three radio bands, we also have simultaneous X-ray observations taken with *Chandra*. An analysis of the radio data, as well as a joint analysis with the X-ray data, yielded a few noteworthy results.

1. We found no anomalous radio bursts or giant-pulse-

like individual pulses in any of the X-band observations. Similarly, the smooth transitioning of the X-band profile between three broad categories seems to have also been unperturbed by the *Swift*-detected X-ray bursts.

2. Our simple radio timing analysis corroborates the findings of Kaspi et al. (2014), which are also supported by Coti Zelati et al. (2015). We presented the absolute alignment of the three radio and 0.3–8 keV profiles. The near-alignment of the radio components with the X-ray profile is similar to the two other radio magnetars that have published alignments. We also make note of a possible transient X-ray feature from Coti Zelati et al. (2015) because of its proximity to the phase of radio emission located ~ 0.2 in phase preceding the peak in the X-ray profile.
3. The evolution of our early radio flux measurements, showing a relatively stable growth from around the time of the initial outburst, is consistent with the continued GBT X-band observations presented in Lynch et al. (2014) and with what they have called a “stable state”³¹. The combination of the gradual flux evolution with the simple timing and profile variability results leads us to extrapolate J1745–2900’s “stable state” back to the time of its initial burst.
4. The shape of J1745–2900’s low frequency radio spectrum is potentially positive or flat, whereas it shows a “typical” spectral index of ~ -1.4 between ~ 2 and 9 GHz, at least during a brief period ~ 100 days after its initial outburst, around the times of two later X-ray bursts. This steep spectral index might indicate a different magnetospheric configuration during these times, although the evolving spectra may be a result of environmental factors and free-free absorption (Lewandowski et al. 2015b). The possible variability of γ means that dedicated observations covering several higher frequency bands need to be carried out over many epochs to confirm this (cf. Torne et al. 2015).
5. We made wideband models of J1745–2900’s low frequency radio “portrait” to measure the scattering timescale, scattering index, and the DM as a function of time. Our average measurements are consistent with what has been published in Spitler et al. (2014), though the ISM parameters may be variable. Time-variable scattering parameters would complicate the predicted sensitivities of future pulsar surveys of the Galactic Center. Lastly, we make a suggestion that our discrepant, nearly simultaneously determined DMs are a manifestation of an ISM-induced frequency-dependent dispersion measure, and that future observations

³¹ While our observations were taken over a shorter range of time (about a third), our cadence of observations is comparable to theirs taken during the “erratic state”, the onset of which was apparently unrelated to X-ray bursts.

could make a case study out of J1745–2900 to investigate $DM(\nu)$ — provided the pulsar remains visible and stable.

J1745–2900 shares several (but not all) properties with the other three radio magnetars, J1809–1943, J1550–5418, and J1622–4950 (Camilo et al. 2006, 2007a; Levin et al. 2010). Common properties of the pulsed radio emission from magnetars are: a) a delay in the appearance of the radio emission after the X-ray outburst onset, b) variable pulse profiles and radio flux on timescales from hours to days, c) a large rotational (spin-down) luminosity with respect to the quiescent X-ray luminosity, d) a decrease of the radio flux as the X-ray flux decays, and e) a flat radio spectrum over a wide range of frequencies. J1745–2900 grossly shares the first three properties with the rest of its class. However, while in all other cases the radio flux was observed to decay as the X-ray outburst was fading, the long-term radio and X-ray flux evolution of J1745–2900 is at variance with this trend. The radio flux shows a re-brightening hundreds of days after the outburst onset and the X-ray emission is decaying very slowly, challenging current crustal cooling models Coti Zelati et al. (2015). Furthermore, the recently published flux measurements by Torne et al. (2015) taken one year after those presented here suggest that the 8.35 GHz flux remained stable at the ~ 3 mJy level over thirty days. Another interesting peculiarity of this radio magnetar was the steep (and possibly free-free absorbed) radio spectrum seen in our observations, though the more recent observations in Torne et al. (2015) suggest that the spectrum has since flattened.

Of particular interest is J1745–2900’s low quiescent luminosity compared to its high rotational power ($L_{X,qui}/L_{rot} < 1$; Rea et al. 2013). This peculiarity of the four radio magnetars, which is at variance with canonical magnetars (for which the fact that $L_{X,qui}/L_{rot} > 1$ has always been used as proof of their magnetically dominated emission (Mereghetti & Stella 1995; Thompson & Duncan 1995; Mereghetti 2008)), has been viewed as evidence for a similar mechanism powering the radio emission from magnetars and normal pulsars alike. In fact, while normal radio pulsars have primarily dipolar-dominated magnetic fields (B_p), magnetars have a substantial toroidal component (B_ϕ) that is present in both the internal and external fields. This toroidal component is the main reason for their quiescent X-ray luminosities, hot surface temperatures, flaring emission, and outburst activity (Thompson et al. 2002; Beloborodov 2009). For a fixed dipolar field, the internal toroidal field has no significant effect on the luminosity unless $B_\phi > B_p$, as is the case for most magnetars (Viganò et al. 2013). Both radio magnetars and high- B radio pulsars have systematically lower toroidal fields and higher rotational energies than typical magnetars; this is in agreement with the former being fainter in quiescence and having a softer X-ray spectrum (a lower crustal toroidal field results in less heating produced by Joule dissipation in the crust, Pons et al. 2009). As for the energy powering the radio emission, simulations of high dipolar field pulsars that have a small toroidal component showed that the particle acceleration and subsequent ignition of the cascade process could proceed

as it does in normal pulsars, successfully reaching the open-field line region and generating pulsed radio emission (Medin & Lai 2010). On the other hand, for an extremely strong toroidal component, it is expected that the particle cascades cannot reach the open-field lines due to the powerful currents formed as a consequence of the twisted magnetosphere. Radio magnetars might lie in between, having a high enough rotational energy to power pair cascades as in normal pulsars, but also having toroidal components lower than typical magnetars, resulting in lower quiescent X-ray luminosities.

In the above picture, the possible radio flux increase, the steep spectrum, and the slow cooling of the X-ray outburst might be explained by the presence of a strongly twisted bundle, which can account for the radio emission and the additional heating by particles slamming onto the surface. If the radio emission is generated by acceleration of particles only in this part of the magnetosphere, then the radio flux and the X-ray flux might be unrelated. In particular, untwisting of the bundle during the outburst decay might induce fewer currents blocking the pair cascade generation, hence more radio emission from this region. However, these are only speculative, plausible hypotheses. Proof of this scenario would need a longer monitoring of the radio and X-ray emission, as well as detailed magnetohydrodynamical simulations of particle acceleration and pair cascades in a strongly magnetized and twisted bundle.

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