

# Nanosegregation in $\text{Na}_2\text{C}_{60}$

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## Abstract.

There is continuous interest in the nature of alkali metal fullerides containing  $\text{C}_{60}^{4-}$  and  $\text{C}_{60}^{2-}$ , because these compounds are believed to be nonmagnetic Mott–Jahn–Teller insulators. This idea could be verified in the case of  $\text{A}_4\text{C}_{60}$ , but  $\text{Na}_2\text{C}_{60}$  is more controversial. By comparing the results of infrared spectroscopy and X-ray diffraction, we found that  $\text{Na}_2\text{C}_{60}$  is segregated into 3–10 nm large regions. The two main phases of the material are insulating  $\text{C}_{60}$  and metallic  $\text{Na}_3\text{C}_{60}$ . We found by neutron scattering that the diffusion of sodium ions becomes faster on heating. Above 470 K  $\text{Na}_2\text{C}_{60}$  is homogeneous and we show IR spectroscopic evidence of a Jahn–Teller distorted  $\text{C}_{60}^{2-}$  anion.

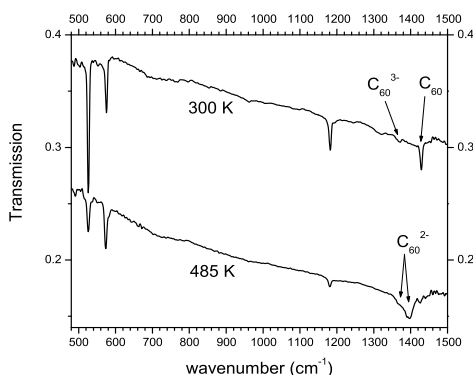
## INTRODUCTION

The proposed [1] Mott–Jahn–Teller ground state of  $\text{A}_4\text{C}_{60}$  ( $\text{A} = \text{K}, \text{Rb}, \text{Cs}$ ) fullerides was experimentally verified by EELS [2], NMR [3] and infrared (IR) [4] spectroscopies. Theory predicts a similar effect in fullerides with  $\text{C}_{60}^{2-}$  ions, the electron-hole inverted analogue of  $\text{C}_{60}^{4-}$ .  $\text{Na}_2\text{C}_{60}$  is the only known alkali-metal fulleride that is believed to contain the  $\text{C}_{60}^{2-}$  dianion and there were several attempts to verify the Mott–Jahn–Teller insulating ground state in this compound as well. Experimental results were not unambiguous [3], e.g. a very weak metallic behavior was found in  $\text{Na}_2\text{C}_{60}$ , but it was concluded that  $\text{K}_4\text{C}_{60}$  and  $\text{Na}_2\text{C}_{60}$  have essentially the same Mott–Jahn–Teller ground state.

In contrast, our infrared, X-ray and neutron scattering experiments indicate that  $\text{Na}_2\text{C}_{60}$  is not a good model system to investigate the ground state of  $\text{A}_2\text{C}_{60}$  fullerides. Above 470 K  $\text{Na}_2\text{C}_{60}$  is homogeneous and at high temperatures we find indication for Jahn–Teller distorted  $\text{C}_{60}^{2-}$  ions. The IR spectroscopy and X-ray diffraction results at ambient temperatures show, however, a nanosegregation of  $\text{Na}_2\text{C}_{60}$  into two phases.

## EXPERIMENTAL

$\text{Na}_2\text{C}_{60}$  was obtained by the reaction of stoichiometric amounts of Na and  $\text{C}_{60}$  at high temperature in a stainless steel capsule. The typical annealing sequence was first 23 days



**FIGURE 1.** The IR spectra of  $\text{Na}_2\text{C}_{60}$  at 300 K and at 485 K. At high temperature the absorption of the  $T_{1u}(4)$  mode could be fitted with two Lorentzians centered at the indicated positions, corresponding to  $\text{C}_{60}^{2-}$ , while at 300 K lines corresponding to  $\text{C}_{60}$  and  $\text{C}_{60}^{3-}$  appear.

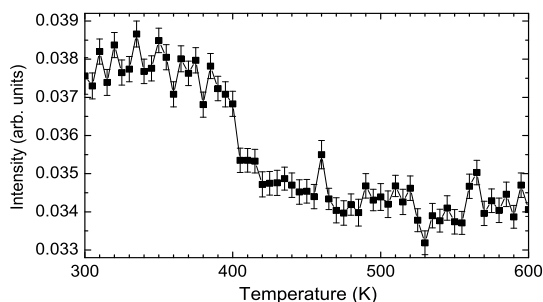
at 350 °C and then 7 days at 450 °C. To homogenize the samples, we reground them about once every five days. Due to the air sensitivity of  $\text{Na}_2\text{C}_{60}$ , the reaction and the handling of the samples were carried out under inert atmosphere.

The sample was characterized by X-ray diffraction, which showed the material to be single phase  $Pa\bar{3}$  (simple cubic)  $\text{Na}_2\text{C}_{60}$ . This result is identical to those in the literature [5].

Infrared measurements were performed on pressed KBr pellets in a Bruker IFS 28 FTIR instrument in a cryostat under dynamic vacuum. For neutron scattering measurements 1.2 g of  $\text{Na}_2\text{C}_{60}$  powder was placed in an annular aluminum sample holder. Temperature dependent elastic fixed-window scans were taken on the High Flux Backscattering Spectrometer of the NIST Center for Neutron Scattering [6]. In this measurement, the incident neutron energy was fixed at 2.08 meV and scattering processes were detected near the elastic line within the 1  $\mu\text{eV}$  energy window of the resolution of the instrument[7]. Thus when a dynamic process became faster than the corresponding timescale of 0.8 ns, the measured intensity decreased.

## RESULTS AND DISCUSSION

The charge state of  $\text{C}_{60}$  anions can be evaluated from the quasilinear relationship between the charge state and the line position of the  $T_{1u}(4)$  IR mode [8]. We expect from this relation a line of  $\text{C}_{60}^{2-}$  at about 1380  $\text{cm}^{-1}$ . Contrary to this expectation, there is no such line in the measured 300 K spectrum (Fig. 1) instead lines characteristic of neutral  $\text{C}_{60}$  and of  $\text{C}_{60}^{3-}$  appear. In contrast to the narrow line of  $\text{C}_{60}$ , that of  $\text{C}_{60}^{3-}$  is smeared into a broad line. This kind of line broadening is common for  $\text{A}_3\text{C}_{60}$  phases and is caused by metallic electrons. The metallic character of an aggregate as small as a few  $\text{C}_{60}^{3-}$  mole-



**FIGURE 2.** Elastic fixed window scan intensity of  $\text{Na}_2\text{C}_{60}$  at  $Q = 1\text{\AA}^{-1}$ .

cules is sufficient to broaden the IR line. Thus  $\text{C}_{60}^{3-}$  and  $\text{C}_{60}$  are not homogeneously distributed in the crystal lattice.

The observation of segregated  $\text{Na}_3\text{C}_{60}$  and  $\text{C}_{60}$  regions in the material by IR spectroscopy apparently contradicts the observation of a single phase by X-ray diffraction. However, if the size of segregated regions is smaller than about 10 nm, then X-ray diffraction cannot resolve the two phases but detects their average. Thus from the comparison of IR and X-ray diffraction measurements we conclude that  $\text{Na}_2\text{C}_{60}$  is nanosegregated at ambient temperatures.

The amount of neutral  $\text{C}_{60}$  in the material could be determined by its selective extraction with toluene. The concentration of the obtained  $\text{C}_{60}$  solution was measured with high-pressure liquid chromatography (HPLC). From this we could estimate the  $\text{C}_{60}$  content to be 26-33 %  $\text{C}_{60}$  in the nominally  $\text{Na}_2\text{C}_{60}$  material.

Heating the material to high temperatures proves that the stoichiometry of the sample is indeed  $\text{Na}_2\text{C}_{60}$  and that at high temperatures there is no neutral  $\text{C}_{60}$  left. The lines of neutral  $\text{C}_{60}$  disappear above 470 K, and a pair of lines appears at  $1369\text{ cm}^{-1}$  and  $1394\text{ cm}^{-1}$  (Fig. 1). This line pair is at about the expected frequency for the  $T_{1u}(4)$  line of  $\text{C}_{60}^{2-}$ . Thus at room temperature the  $\text{C}_{60}$  content is not material left unreacted during the high temperature synthesis, but rather the product of the reaction  $3\text{C}_{60}^{2-} \longrightarrow \text{C}_{60} + 2\text{C}_{60}^{3-}$  taking place on cooling after preparation. This reaction is reversible on heating and cooling, though a total retransformation at ambient temperature following treatment at high temperature is reached only after about two weeks.

The line pair of  $\text{C}_{60}^{2-}$  is the lower frequency analogue of the line pair of  $\text{C}_{60}^{4-}$  [4], indicating that  $\text{C}_{60}^{2-}$  is distorted to a  $D_{3d}$  or a  $D_{5d}$  geometry due to the molecular Jahn-Teller effect [4].

The synproportion reaction on heating ought to be accompanied by diffusion of Na ions in the lattice. This could be proven by neutron scattering. Figure 2 shows the temperature dependence of the elastic line at  $Q = 1\text{\AA}^{-1}$ . The intensity shows a Debye-Waller-type overall decrease, but has an unusual drop near 400 K. We interpret this with the increase of the jump diffusion of sodium ions between tetrahedral and off-centered octahedral sites. Above 400 K it becomes fast enough to be resolved by the

instrument and the incoherent scattering contribution of the sodium ions is removed from the fixed-window intensity. The different transition temperatures found in the IR and neutron measurements may be due to the different timescales.

## CONCLUSION

We propose that  $\text{Na}_2\text{C}_{60}$  is nanosegregated at room temperature. The two main phases of  $\text{Na}_2\text{C}_{60}$  are insulating  $\text{C}_{60}$  and metallic  $\text{Na}_3\text{C}_{60}$ . The size of the homogeneous regions is about 3-10 nm. The segregation disappears on heating when the jump diffusion of sodium ions becomes faster. A similar segregated phase, the "intermediate phase" is known for  $\text{KC}_{60}$  [9]. In  $\text{Na}_2\text{C}_{60}$  at high temperatures, the sodium distribution is homogeneous and a Jahn–Teller distortion of  $\text{C}_{60}^{2-}$  ions is observed.

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## REFERENCES

1. Fabrizio, M., and Tosatti, E., *Phys. Rev. B*, **55**, 13465 (1997).
2. Knupfer, M., and Fink, J., *Phys. Rev. Lett.*, **79**, 2714 (1997).
3. Brouet, V., Alloul, H., Garaj, S., and Forró, L., *Phys. Rev. B*, **66**, 155122 (2002).
4. Kamarás, K., Klupp, G., Tanner, D. B., Hebard, A. F., Nemes, N. M., and Fischer, J. E., *Phys. Rev. B*, **65**, 052103 (2002).
5. Yildirim, T., Hong, S., Harris, A. B., and Mele, E. J., *Phys. Rev. B*, **48**, 12262 (1993).
6. Meyer, A., Dimeo, R. M., Gehring, P. M., and Neumann, D. A., *Rev. Sci. Instrum.*, **74**, 2759 (2003).
7. Becker, T., and Smith, J. C., *Phys. Rev. E*, **67**, 021904 (2003).
8. Pichler, T., Winkler, R., and Kuzmany, H., *Phys. Rev. B*, **49**, 15879 (1994).
9. Faigel, G., Bortel, G., Tegze, M., Gránásy, L., Pekker, S., Oszlányi, G., Chauvet, O., Baumgartner, G., Forro, L., Stephens, P. W., Mihály, G., and Jánosy, A., *Phys. Rev. B*, **52**, 3199 (1995).