



NEW INTEGRABILITY CONDITIONS OF DERIVATION EQUATIONS IN A SUBSPACE OF ASYMMETRIC AFFINE CONNECTION SPACE

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Abstract. In the work [10] we obtained derivational equations of a submanifold of a space L_N with asymmetric affine connection. Based on asymmetry of the connection we define four kinds of covariant derivative and obtain four kinds of derivational equations.

In [20] are examined integrability conditions of derivational equations, using the 1st and the 2nd kind of derivative, and in the present work we do it on the base of the 3rd and 4th kind.

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1. INTRODUCTION

In 1922 Cartan was put forward a modification of General Relativity Theory (GRT), by relaxing the assumption that the affine connection has vanishing the antisymmetric part (torsion tensor), and relating the torsion to the density of intrinsic angular momentum. Also, the torsion is implicit in the 1928 Einstein theory of gravitation with teleparallelism.

From 1923 to the end of his life Einstein worked on various variants of Unified Field Theory (Non-symmetric Gravitational Theory-NGT) [3]. This theory had the aim to unite gravitation theory and the theory of electromagnetism. Introducing different variants of his NGT, Einstein used a complex basic tensor, with a symmetric real part and a skew-symmetric imaginary part. Starting from 1950, Einstein used the real non-symmetric basic tensor g , sometimes called *generalized Riemannian metric/manifold*.

Notice that in NGT the symmetric part g_{ij} of the basic tensor g_{ij} ($g_{ij} = g_{ij} + g_{ij}$) is related to gravitation, and the skew-symmetric one g_{ij} to electromagnetism.

While on a Riemannian space the connection coefficients are expressed by virtue of the metric, g_{ij} , in Einstein's work on NGT the connection between these

magnitudes is determined by the so-called *Einstein metricity condition*, i.e. the non-symmetric metric tensor g and the connection components L_{ij}^k are connected with the equations

$$\frac{\partial g_{ij}}{\partial x^m} - L_{im}^p g_{pj} - L_{mj}^p g_{ip} = 0. \quad (1.1)$$

The choice of a connection in NGT is not uniquely determined. In particular, in NGT there exist two kinds of the covariant derivative. For example, for tensor a_j^i :

$$a_{j|m}^+ = \frac{\partial a_j^i}{\partial x^m} + L_{pm}^i a_j^p - L_{jm}^p a_p^i; \quad a_{j|m}^- = \frac{\partial a_j^i}{\partial x^m} + L_{mp}^i a_j^p - L_{mj}^p a_p^i,$$

where the lowering and the rising of indices one defines by equations

$$g_{pi} g^{pj} = g_{ip} g^{jp} = \delta_i^j, \quad (1.2)$$

Einstein considered only one curvature tensor:

$$R_{klm}^i = L_{kl,m}^i - L_{km,l}^i - L_{sl}^i L_{km}^s + L_{sm}^i L_{kl}^s \quad (1.3)$$

and proved Bianchi type identity for covariant curvature tensor (see [2])

$$R_{-+--+}^{iklm|n} + R_{-++-+}^{ikmn|l} + R_{-+--+}^{iknl|m} = 0,$$

where is $R_{iklm} = g_{si} R^s_{klm}$.

Afterwards, several mathematicians dealt with non-symmetric affine connection, for example, Eisenhart [4], Prvanović [13], Minčić [9–12], Stanković [16, 17] etc. Sinyukov [14] introduced the concept of almost geodesic mappings between affine connected spaces without torsion. Mikeš [1], [5–8, 13], [15], [18] gave some significant contributions to the study of geodesic and almost geodesic mappings of affine connected, Riemannian and Einstein spaces.

Let L_N be a space with asymmetric affine connection L_{jk}^i (in local coordinates), and torsion tensor T_{jk}^i .

In L_N one can define four kinds of covariant derivatives. For example, for a tensor a_j^i we have

$$\begin{aligned} a_{j|_1 m}^i &= a_{j,m}^i + L_{pm}^i a_j^p - L_{jm}^p a_p^i, & a_{j|_2 m}^i &= a_{j,m}^i + L_{mp}^i a_j^p - L_{mj}^p a_p^i, \\ a_{j|_3 m}^i &= a_{j,m}^i + L_{pm}^i a_j^p - L_{mj}^p a_p^i, & a_{j|_4 m}^i &= a_{j,m}^i + L_{mp}^i a_j^p - L_{jm}^p a_p^i. \end{aligned}$$

In [6] is proved very important theorem

Theorem 1. *Let $(L_N, g_{ij} = g_{ij} + g_{ij})$ be an asymmetric affine connection space and Γ_{jk}^i be the Levi-Civita connection of g_{ij} . Let L_{jk}^i be a linear connection with torsion T_{jk}^i . Then L_{jk}^i is uniquely determined by the following formula*

$$L_{i,jk} = \Gamma_{i,jk} + \frac{1}{2}(T_{i,jk} + T_{k,ij} - T_{j,ki}) - \frac{1}{2}(g_{jk|i} + g_{ki|j} - g_{ji|k}), \quad (1.4)$$

where $L_{i,jk} = g_{pi} L_{jk}^p$.

A submanifold $X_M \subset L_N$ is defined by equations

$$x^i = x^i(u^1, \dots, u^M) = x^i(u^\alpha), \quad i = \overline{1, N}.$$

Partial derivatives $B_\alpha^i = \frac{\partial x^i}{\partial u^\alpha}$ ($\text{rank}(B_\alpha^i) = M$) define tangent vectors on X_M .

Consider $N - M$ contravariant vectors C_A^i ($A, B, C, \dots, \in \{M+1, \dots, N\}$) defined on X_M and linearly independent, and let the matrix $\begin{pmatrix} B_\alpha^i \\ C_A^i \end{pmatrix}$ be inverse for the matrix (B_α^i, C_A^i) provided that the following conditions are satisfied [19, 20]:

$$\begin{aligned} a) B_\alpha^i B_i^\beta &= \delta_\alpha^\beta; & b) B_\alpha^i C_i^A &= 0; & c) B_i^\alpha C_A^i &= 0; \\ d) C_A^i C_i^B &= \delta_A^B; & e) B_\alpha^i B_j^\alpha + C_A^i C_j^A &= \delta_j^i; \end{aligned}$$

The magnitudes B_α^i, B_i^α are **projection factors (tangent vectors)**, and the magnitudes C_A^i, C_i^A are **affine pseudonormals** [10, 20].

The **induced connection** on X_M is [10, 19, 20]:

$$\tilde{L}_{\beta\gamma}^\alpha = B_i^\alpha (B_{\beta,\gamma}^i + L_{jk}^i B_\beta^j B_\gamma^k),$$

where $B_{\beta,\gamma}^i = \partial B_\beta^i / \partial u^\gamma = \partial^2 x^i / \partial u^\beta \partial u^\gamma$. Because L is asymmetric by virtue of j, k , \tilde{L} is asymmetric in β, γ too. The submanifold X_M endowed with \tilde{L} becomes L_M and we write $L_M \subset L_N$.

The set of pseudonormals of the submanifold $X_M \subset L_N$ makes a **pseudonormal bundle** of X_M , and we note it X_{N-M}^N . We have defined in [10] **induced connections of pseudonormal bundle** with coefficients

$$\bar{L}_{B\mu}^A = C_i^A (C_{B,\mu}^i + L_{jk}^i C_B^j B_\mu^k), \quad \bar{L}_{B\mu}^A = C_i^A (C_{B,\mu}^i + L_{kj}^i C_B^j B_\mu^k).$$

As the coefficients L, \tilde{L}, \bar{L} are generally asymmetric, we can define four kinds of covariant derivative for a tensor, defined in the points of X_M . For example:

$$\begin{aligned} t_{j\beta B}^{i\alpha A}|_\mu &= t_{j\beta B,\mu}^{i\alpha A} + L_{pm}^i t_{j\beta B}^{p\alpha A} - L_{jm}^p t_{p\beta B}^{i\alpha A} + \tilde{L}_{\pi\mu}^\alpha t_{j\beta B}^{i\pi A} - \tilde{L}_{\beta\mu}^\pi t_{j\pi B}^{i\alpha A} + \bar{L}_{P\mu}^A t_{j\beta B}^{i\alpha P} - \bar{L}_{\beta\mu}^P t_{j\beta P}^{i\alpha A}, \\ t_{j\beta B}^{i\alpha A}|_\mu &= t_{j\beta B,\mu}^{i\alpha A} + L_{mp}^i t_{j\beta B}^{p\alpha A} - L_{mj}^p t_{p\beta B}^{i\alpha A} + \tilde{L}_{\mu\pi}^\alpha t_{j\beta B}^{i\pi A} - \tilde{L}_{\mu\beta}^\pi t_{j\pi B}^{i\alpha A} + \bar{L}_{P\mu}^A t_{j\beta B}^{i\alpha P} - \bar{L}_{\beta\mu}^P t_{j\beta P}^{i\alpha A} \\ t_{j\beta B}^{i\alpha A}|_\mu &= t_{j\beta B,\mu}^{i\alpha A} + L_{pm}^i t_{j\beta B}^{p\alpha A} - L_{mj}^p t_{p\beta B}^{i\alpha A} + \tilde{L}_{\pi\mu}^\alpha t_{j\beta B}^{i\pi A} - \tilde{L}_{\mu\beta}^\pi t_{j\pi B}^{i\alpha A} + \bar{L}_{P\mu}^A t_{j\beta B}^{i\alpha P} - \bar{L}_{\beta\mu}^P t_{j\beta P}^{i\alpha A} \\ t_{j\beta B}^{i\alpha A}|_\mu &= t_{j\beta B,\mu}^{i\alpha A} + L_{mp}^i t_{j\beta B}^{p\alpha A} - L_{jm}^p t_{p\beta B}^{i\alpha A} + \tilde{L}_{\mu\pi}^\alpha t_{j\beta B}^{i\pi A} - \tilde{L}_{\mu\beta}^\pi t_{j\pi B}^{i\alpha A} + \bar{L}_{P\mu}^A t_{j\beta B}^{i\alpha P} - \bar{L}_{\beta\mu}^P t_{j\beta P}^{i\alpha A} \end{aligned}$$

In this manner four connections ∇_{θ} on $X_M \subset L_N$ are defined. We shall note the obtained structures $(X_M \subset L_N, \nabla_{\theta}, \theta \in \{1, \dots, 4\})$.

2. NEW INTEGRABILITY CONDITIONS OF DERIVATIONAL EQUATIONS FOR TANGENTS

2.0. We have obtained in [20] integrability conditions of derivational equations and corresponding Gauss-Codazzi equations in the structure $(X_M \subset L_N, \nabla_{\theta}, \theta \in \{1, 2\})$. In the present work we are solving this problem for $\theta \in \{3, 4\}$. Based on the Theorem 2.2. in [10], the following derivational equations for tangents are in force

$$B_{\alpha|\mu}^i = \Omega_{\alpha\mu}^P C_P^i, \quad \theta \in \{3, 4\} \quad (2.1)$$

$$B_{i|\mu}^{\alpha} = \widehat{\Omega}_{P\mu}^{\alpha} C_i^P, \quad \theta \in \{3, 4\}, \quad (2.2)$$

and, (see [20]), for pseudonormals

$$C_{A|\mu}^i = -\widehat{\Omega}_{A\mu}^{\pi} B_{\pi}^i, \quad \theta \in \{3, 4\} \quad (2.3)$$

$$C_{i|\mu}^A = -\Omega_{\pi\mu}^A B_i^{\pi}, \quad \theta \in \{3, 4\}. \quad (2.4)$$

In this manner, one obtains

$$B_{\alpha|\mu|\omega}^i - B_{\alpha|\omega|\mu}^i = (\widehat{\Omega}_{P\mu}^{\pi} \Omega_{\omega}^P - \widehat{\Omega}_{P\omega}^{\pi} \Omega_{\mu}^P) B_{\pi}^i + (\Omega_{\alpha\mu|\omega}^P - \Omega_{\alpha\omega|\mu}^P) C_P^i, \quad \theta, \omega \in \{3, 4\}. \quad (2.5)$$

and analogously

$$B_{i|\mu|\omega}^{\alpha} - B_{i|\omega|\mu}^{\alpha} = (\Omega_{\pi\mu}^P \widehat{\Omega}_{P\omega}^{\alpha} - \Omega_{\pi\omega}^P \widehat{\Omega}_{P\mu}^{\alpha}) B_i^{\pi} + (\widehat{\Omega}_{P\mu|\omega}^{\alpha} - \widehat{\Omega}_{P\omega|\mu}^{\alpha}) C_i^P, \quad \theta \in \{3, 4\}.$$

where [10]

$$\begin{aligned} \Omega_{\alpha\mu}^P &= C_i^P (B_{\alpha,\mu}^i + L_{pm}^i B_{\alpha}^p B_{\mu}^m) = \Omega_{\alpha\mu}^P, \\ \Omega_{\alpha\mu}^P &= C_i^P (B_{\alpha,\mu}^i + L_{mp}^i B_{\alpha}^p B_{\mu}^m) = \Omega_{\alpha\mu}^P, \end{aligned} \quad (2.6)$$

$$\widehat{\Omega}_{P\mu}^{\alpha} = C_P^i (B_{i,\mu}^{\alpha} + L_{mi}^p B_P^{\alpha} B_{\mu}^m) = \widehat{\Omega}_{P\mu}^{\alpha}, \quad \widehat{\Omega}_{P\mu}^{\alpha} = C_P^i (B_{i,\mu}^{\alpha} + L_{im}^p B_P^{\alpha} B_{\mu}^m) = \widehat{\Omega}_{P\mu}^{\alpha}.$$

From the Theorem 2.2. in [10] the induced connection \widetilde{L} in the structure $(X_M \subset L_N, \nabla_{\theta}, \theta \in \{3, 4\})$ is symmetric, i.e.

$$\widetilde{T}_{\beta\gamma}^{\alpha} = 0, \quad (2.7)$$

and based on the Theorem 3.2. in this case is

$$\overline{L}_1 = \overline{L}_2 = \overline{L}, \quad (2.8)$$

that is there exists an unique connection in the pseudonormal bundle.

2.1. Using the Ricci-type identities (12,13) in [9], by virtue of (2.7), we obtain

$$B_{\alpha|\mu|\nu}^i - B_{\alpha|\nu|\mu}^i = R_{\theta-2}^i{}_{pmn} B_{\alpha}^p B_{\mu}^m B_{\nu}^n - \widetilde{R}_{\alpha\mu\nu}^{\pi} B_{\pi}^i + (-1)^{\theta} \widetilde{T}_{\mu\nu}^{\pi} B_{\alpha|\pi}^i, \quad \theta \in \{3, 4\}. \quad (2.9)$$

where

$$\begin{aligned} a) \quad R_1^i{}_{jmn} &= L_{jm,n}^i - L_{jn,m}^i + L_{jm}^p L_{pn}^i - L_{jn}^p L_{pm}^i, \\ b) \quad R_2^i{}_{jmn} &= L_{mj,n}^i - L_{nj,m}^i + L_{mj}^p L_{np}^i - L_{nj}^p L_{mp}^i, \end{aligned} \quad (2.10)$$

are **curvature tensors of the 1st respectively the 2nd kind** of the L_N and $\widetilde{R}_{\beta\mu\nu}^{\alpha}$ is, with respect of (2.8), curvature tensor of $L_M^o \subset L_N$, where L_M^o is a subspace with symmetric affine connection.

Further, we examine integrability conditions for derivational equations of B_{α}^i and B_i^{α} , i.e. for $B_{\alpha|\mu}^i$, $B_{i|\mu}^{\alpha}$, $\theta \in \{3, 4\}$.

Substituting $\theta = \omega \in \{3, 4\}$ into (2.5) and comparing with (2.9), taking into consideration (2.7), we get **the 1st and the 2nd kind integrability condition of derivational equation (2.1)** in the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$:

$$R_{\theta-2}^i{}_{pmn} B_{\alpha}^p B_{\mu}^m B_{\nu}^n - \widetilde{R}_{\alpha\mu\nu}^{\pi} B_{\pi}^i = (\widehat{\Omega}_{P\mu}^{\pi} \Omega_{\theta}^P{}_{\alpha\nu} - \widehat{\Omega}_{P\nu}^{\pi} \Omega_{\theta}^P{}_{\alpha\mu}) B_{\pi}^i + (\Omega_{\theta}^P{}_{\alpha\mu|\nu} - \Omega_{\theta}^P{}_{\alpha\nu|\mu}) C_{\pi}^i, \quad \theta \in \{3, 4\}. \quad (2.11)$$

Multiplying this equation equation with B_i^{λ} and taking into consideration (1.2), we obtain

$$R_{\theta-2}^i{}_{pmn} B_i^{\lambda} B_{\alpha}^p B_{\mu}^m B_{\nu}^n - \widetilde{R}_{\alpha\mu\nu}^{\lambda} = \widehat{\Omega}_{P\mu}^{\lambda} \Omega_{\theta}^P{}_{\alpha\nu} - \widehat{\Omega}_{P\nu}^{\lambda} \Omega_{\theta}^P{}_{\alpha\mu}, \quad \theta \in \{3, 4\}. \quad (2.12)$$

i.e the **Gauss equation of the 1st and the 2nd kind** in the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$.

If one multiplies (2.11) with C_i^L , it follows that

$$R_{\theta-2}^i{}_{pmn} C_i^L B_{\alpha}^p B_{\mu}^m B_{\nu}^n = \Omega_{\theta}^L{}_{\alpha\mu|\nu} - \Omega_{\theta}^L{}_{\alpha\nu|\mu}, \quad \theta \in \{3, 4\}. \quad (2.13)$$

which are **the 1st Codazzi equation of the 1st and the 2nd kind** for the cited structure.

2.1' Further, if we use the Ricci type identities ([9], equation (12))

$$\begin{aligned} B_{i|\mu|_3}^\alpha - B_{i|_4}^\alpha|_{\mu|_4} &= -R_{2imn}^P B_p^\alpha B_\mu^m B_\nu^n + \widetilde{R}_{\pi\mu\nu}^\alpha B_i^\pi, \\ B_{i|\mu|_4}^\alpha - B_{i|_4}^\alpha|_{\mu|_4} &= -R_{1imn}^P B_p^\alpha B_\mu^m B_\nu^n + \widetilde{R}_{\pi\mu\nu}^\alpha B_i^\pi, \end{aligned} \quad (2.14)$$

and (2.3') for $\theta = \omega \in \{3, 4\}$, we have

$$\begin{aligned} -R_{2imn}^P B_p^\alpha B_\mu^m B_\nu^n + \widetilde{R}_{\pi\mu\nu}^\alpha B_i^\pi &= (\Omega_{3\pi\mu}^P \widehat{\Omega}_{3P\nu}^\alpha - \Omega_{\theta\pi\nu}^P \widehat{\Omega}_{\theta P\mu}^\alpha) B_i^\pi + (\widehat{\Omega}_{3P\mu|_3}^\alpha - \widehat{\Omega}_{3P\nu|_3}^\alpha) C_i^P, \\ -R_{1imn}^P B_p^\alpha B_\mu^m B_\nu^n + \widetilde{R}_{\pi\mu\nu}^\alpha B_i^\pi &= (\Omega_{4\pi\mu}^P \widehat{\Omega}_{4P\nu}^\alpha - \Omega_{\theta\pi\nu}^P \widehat{\Omega}_{\theta P\mu}^\alpha) B_i^\pi + (\widehat{\Omega}_{4P\mu|_4}^\alpha - \widehat{\Omega}_{4P\nu|_4}^\alpha) C_i^P \end{aligned} \quad (2.15)$$

which are **the 1st and the 2nd kind integrability conditions of derivational equation (2.2)** in the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$.

a') By multiplying the previous equation with B_λ^i , one obtains

$$\begin{aligned} \widetilde{R}_{\lambda\mu\nu}^\alpha - R_{2imn}^P B_p^\alpha B_\lambda^i B_\mu^m B_\nu^n &= \Omega_{3\lambda\mu}^P \widehat{\Omega}_{3P\nu}^\alpha - \Omega_{3\lambda\nu}^P \widehat{\Omega}_{3P\mu}^\alpha, \\ \widetilde{R}_{\lambda\mu\nu}^\alpha - R_{1imn}^P B_p^\alpha B_\lambda^i B_\mu^m B_\nu^n &= \Omega_{4\lambda\mu}^P \widehat{\Omega}_{4P\nu}^\alpha - \Omega_{4\lambda\nu}^P \widehat{\Omega}_{4P\mu}^\alpha, \end{aligned} \quad (2.16)$$

which is **another form of (2.12)**.

If we exchange at (2.12) $i \leftrightarrow p$, $\alpha \leftrightarrow \lambda$, for $\theta = 3$ from that equation one gets

$$R_{1imn}^P B_p^\alpha B_\mu^m B_\nu^n - \widetilde{R}_{\lambda\mu\nu}^\alpha = \widehat{\Omega}_{3P\mu}^\alpha \Omega_{3\lambda\nu}^P - \widehat{\Omega}_{3P\nu}^\alpha \Omega_{3\lambda\mu}^P.$$

Summing this equation with 1st case in (2.16), one concludes:

$$(R_{1imn}^P - R_{2imn}^P) B_p^\alpha B_\lambda^i B_\mu^m B_\nu^n = 0. \quad (2.17)$$

Putting $\theta = 4$ at (2.12), we get the 1st case from (2.16), and together with 2nd case it follows.

If we multiply (2.15) with C_L^i , it follows that

$$\begin{aligned} -R_{2imn}^P B_p^\alpha C_L^i B_\mu^m B_\nu^n &= \widehat{\Omega}_{3L\mu|_3}^\alpha - \widehat{\Omega}_{3L\nu|_3}^\alpha, \\ -R_{1imn}^P B_p^\alpha C_L^i B_\mu^m B_\nu^n &= \widehat{\Omega}_{4L\mu|_4}^\alpha - \widehat{\Omega}_{4L\nu|_4}^\alpha, \end{aligned} \quad (2.18)$$

and this is **another form for the 1st Codazzi equation of the 1st and the 2nd kind in the cited structure**.

2.2. Further we use the Ricci-type identity (equation (46) from [9])

$$B_{\alpha|\mu|_3}^i - B_{\alpha|_4}^i|_{\mu|_4} = R_{4p\mu\nu}^i B_\alpha^p - \widetilde{R}_{\alpha\mu\nu}^\pi B_\pi^i, \quad (2.19)$$

where

$$R_{4j\mu\nu}^i = (L_{jm,n}^i - L_{jn,m}^i + L_{jm}^p L_{np}^i - L_{nj}^p L_{pm}^i) B_\mu^m B_\nu^n + T_{jm}^i (B_{\mu,\nu}^m - \tilde{L}_{\mu\nu}^\pi B_\pi^m), \quad (2.20)$$

is the 4th kind curvature tensor of L_N with respect to $X_M \subset L_N$. On the other hand, putting into (2.5) $\theta = 3, \omega = 4$ and comparing with (2.19), we get the 3rd kind integrability condition of derivational equation (2.1) in the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$:

$$R_{4p\mu\nu}^i B_\alpha^p - \tilde{R}_{\alpha\mu\nu}^\pi B_\pi^i = (\hat{\Omega}_{3P\mu}^\pi \hat{\Omega}_{4\alpha\nu}^P - \hat{\Omega}_{4P\nu}^\pi \hat{\Omega}_{3\alpha\mu}^P) B_\pi^i + (\hat{\Omega}_{3\alpha\mu|v}^P - \hat{\Omega}_{4\alpha\nu|3}^P) C_P^i. \quad (2.21)$$

a) If one multiplies this equation with B_i^λ , we have

$$R_{4p\mu\nu}^i B_\alpha^p B_i^\lambda - \tilde{R}_{\alpha\mu\nu}^\lambda = \hat{\Omega}_{3P\mu}^\pi \hat{\Omega}_{4\alpha\nu}^P - \hat{\Omega}_{4P\nu}^\pi \hat{\Omega}_{3\alpha\mu}^P \quad (2.22)$$

and this is Gauss equation of the 3rd kind in the cited structure.

b) Multiplying (2.21) with C_i^L , we get

$$R_{4p\mu\nu}^i C_i^L B_\alpha^p = \hat{\Omega}_{3\alpha\mu|v}^P - \hat{\Omega}_{4\alpha\nu|3}^P, \quad (2.23)$$

i.e. the 1st Codazzi equation of the 3rd kind in the same structure.

2.2'. Based on equation (46) in [9], we have

$$B_{34}^{\alpha|_{\mu|v}} - B_{43}^{\alpha|_{v|\mu}} = R_{3i\mu\nu}^p B_p^\alpha + \tilde{R}_{\pi\mu\nu}^\alpha B_i^\pi,$$

where

$$R_{3j\mu\nu}^i = (L_{jm,n}^i - L_{nj,m}^i + L_{jm}^p L_{np}^i - L_{nj}^p L_{pm}^i) B_\mu^m B_\nu^n + T_{jm}^i (B_{\mu,\nu}^m - \tilde{L}_{\nu\mu}^\pi B_\pi^m), \quad (2.24)$$

is the 3rd kind curvature tensor of L_N with respect to $X_M \subset L_N$.

Simultaneously, putting into (2) $\theta = 3, \omega = 4$, and comparing with (2), we obtain the 3rd integrability condition of derivational equation (2.2) in the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$:

$$R_{3i\nu\mu}^p B_p^\alpha + \tilde{R}_{\pi\mu\nu}^\alpha B_i^\pi = (\hat{\Omega}_{3\pi\mu}^P \hat{\Omega}_{4P\nu}^\alpha - \hat{\Omega}_{4\pi\nu}^P \hat{\Omega}_{3P\mu}^\alpha) B_i^\pi + (\hat{\Omega}_{3P\mu|v}^\alpha - \hat{\Omega}_{4P\nu|3}^\alpha) C_i^P. \quad (2.25)$$

a') From here, multiplying with B_λ^i

$$R_{3i\nu\mu}^p B_p^\alpha B_\lambda^i + \tilde{R}_{\pi\mu\nu}^\alpha = \hat{\Omega}_{3\pi\mu}^P \hat{\Omega}_{4P\nu}^\alpha - \hat{\Omega}_{4\pi\nu}^P \hat{\Omega}_{3P\mu}^\alpha, \quad (2.26)$$

which is another form of the 3rd kind Gauss equation in the cited structure.

b') Similarly as the equation (2.23), we get

$$R_{4\ p\nu\mu}^i C_i^L B_\alpha^P = \Omega_{3\ \alpha\mu|_\nu}^P - \Omega_{4\ \alpha\nu|_\mu}^P, \quad (2.27)$$

and that is another form of (2.23).

Now, we can state the following theorems

Theorem 2. *The 1st and 2nd kind integrability conditions derivational equations (2.1) and (2.1') for submanifold $X_M \subset L_N$ with the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$, for the connections ∇_θ are given in (2.11) and (2.15) respectively. The 3rd kind of these conditions are given in (2.21).*

Theorem 3. *Gauss equations of the 1st and the 2nd kind are given in (2.12), and of the 3rd one in (2.22). The 1st Codazzi equations of the 1st and the 2nd kind are given in (2.18), and of the 3rd kind in (2.23). The equations (2.16), (2.18), (2.26), (2.27) are another forms of previous equations.*

3. INTEGRABILITY CONDITIONS OF DERIVATIONAL EQUATIONS FOR PSEUDONORMALS

3.0. Further, we use the similar procedure on derivational equation of pseudonormals.

Using (2.3), (2.1), we get

$$C_{\theta|\mu|_\nu}^i - C_{\omega|_\nu|_\mu}^i = -(\widehat{\Omega}_{\theta A\mu|_\nu}^\pi - \widehat{\Omega}_{\omega A\nu|_\mu}^\pi) B_\pi^i - (\widehat{\Omega}_{\theta A\mu}^\pi \widehat{\Omega}_{\omega}^P \pi_\nu - \widehat{\Omega}_{\omega A\nu}^\pi \widehat{\Omega}_{\theta}^P \pi_\mu) C_P^i \quad (3.1)$$

and from (2.4), (2.2):

$$C_{\theta|\mu|_\nu}^A - C_{\omega|_\nu|_\mu}^A = -(\Omega_{\theta \pi\mu|_\nu}^A - \Omega_{\omega \pi\nu|_\mu}^A) B_\pi^i - (\Omega_{\theta \pi\mu}^A \widehat{\Omega}_{\omega}^P \pi_\nu - \Omega_{\omega \pi\nu}^A \widehat{\Omega}_{\theta}^P \pi_\mu) C_i^P \\ \theta, \omega \in \{3, 4\}.$$

We need also the Ricci type identities ([11], Equation (2.19))

$$C_{\theta|\mu|_\nu}^i - C_{\omega|_\nu|_\mu}^i = R_{\theta-2\ pmn}^i C_A^p B_\mu^m B_\nu^n - \bar{R}_{A\mu\nu}^P C_P^i, \quad \theta \in \{3, 4\}, \quad (3.2)$$

and in the same way

$$C_{3\ 3}^A|_{\mu|_\nu} - C_{3\ 3}^A|_{\nu|_\mu} = \bar{R}_{P\mu\nu}^A C_i^P - R_{1\ imn}^P C_p^A B_\mu^m B_\nu^n, \\ C_{4\ 4}^A|_{\mu|_\nu} - C_{4\ 4}^A|_{\nu|_\mu} = \bar{R}_{P\mu\nu}^A C_i^P - R_{2\ imn}^P C_p^A B_\mu^m B_\nu^n, \quad (3.3)$$

where R, \bar{R} are given at (2.10) and \bar{R} at [11]

$$\widetilde{R}_{B\mu\nu}^A = \bar{L}_{B\mu,\nu}^A - \bar{L}_{B\nu,\mu}^A + \bar{L}_{B\mu}^P \bar{L}_{P\nu}^A - \bar{L}_{B\nu}^P \bar{L}_{P\mu}^A. \quad (3.4)$$

It follows that

$$C_{A|\mu|v}^i - C_{A|v|\mu}^i = R_{pmn}^i C_A^p - \bar{R}_{A\mu\nu}^P C_P^i, \quad (3.5)$$

and analogously

$$C_{i|\mu|v}^A - C_{i|v|\mu}^A = R_{imn}^P C_A^p + \bar{R}_{P\nu\mu}^A C_i^P,$$

The magnitude $\bar{R}_{B\mu\nu}^A$ is **curvature tensor of L_N with respect to the pseudonormal submanifold X_{N-M}^N** in the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$.

3.1. Taking $\theta = \omega \in \{3, 4\}$ in (3.1) and comparing with (3.2), we obtain **the 1st and the 2nd kind integrability condition of derivational equation (2.3)** (for pseudonormals) in the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$:

$$\begin{aligned} R_{\theta-2}^i{}_{pmn} C_A^p B_\mu^m B_\nu^n - \bar{R}_{A\mu\nu}^P C_P^i \\ = -(\hat{\Omega}_{A\mu|v}^\pi + \hat{\Omega}_{\theta A\nu|\mu}^\pi) B_\pi^i - (\hat{\Omega}_{\theta A\mu}^\pi \hat{\Omega}_{\theta \pi\nu}^P + \hat{\Omega}_{\theta A\nu}^\pi \hat{\Omega}_{\theta \pi\mu}^P) C_P^i, \quad \theta \in \{3, 4\}, \end{aligned} \quad (3.6)$$

a) Multiplying (3.6) with B_i^λ , one gets:

$$R_{\theta-2}^i{}_{pmn} B_i^\lambda C_A^p B_\mu^m B_\nu^n = -\hat{\Omega}_{A\mu|v}^\lambda + \hat{\Omega}_{\theta A\nu|\mu}^\lambda, \quad \theta \in \{3, 4\}. \quad (3.7)$$

which is **one more form of the 1st Codazzi equation (2.13)**.

b) If we multiply (3.6) with C_i^L , one obtains

$$R_{\theta-2}^i{}_{pmn} C_i^L C_A^p B_\mu^m B_\nu^n - \bar{R}_{A\mu\nu}^L = -\hat{\Omega}_{\theta A\mu}^\pi \hat{\Omega}_{\theta \pi\nu}^L + \hat{\Omega}_{\theta A\nu}^\pi \hat{\Omega}_{\theta \pi\mu}^L, \quad \theta \in \{3, 4\}. \quad (3.8)$$

and that is **the 2nd Codazzi equation of the 1st and the 2nd kind** in the cited structure.

3.1'. If one takes $\theta = \omega \in \{3, 4\}$ in (3) and compare with (3.3) we obtain **the 1st and the 2nd kind integrability condition of derivational equation (2.4)** in the structure $(X_M \subset L_N, \nabla, \theta \in \{3, 4\})$:

$$\begin{aligned} R_{2}^p{}_{imn} C_p^A B_\mu^m B_\nu^n - \bar{R}_{P\mu\nu}^A C_i^P &= (\hat{\Omega}_{3\pi\mu|v}^A - \hat{\Omega}_{3\pi\nu|\mu}^A) B_i^\pi + (\hat{\Omega}_{3\pi\mu}^A \hat{\Omega}_{3P\nu}^\pi - \hat{\Omega}_{3\pi\nu}^A \hat{\Omega}_{3P\mu}^\pi) C_i^P \\ R_{1}^p{}_{imn} C_p^A B_\mu^m B_\nu^n - \bar{R}_{P\mu\nu}^A C_i^P &= (\hat{\Omega}_{4\pi\mu|v}^A - \hat{\Omega}_{4\pi\nu|\mu}^A) B_i^\pi + (\hat{\Omega}_{4\pi\mu}^A \hat{\Omega}_{4P\nu}^\pi - \hat{\Omega}_{4\pi\nu}^A \hat{\Omega}_{4P\mu}^\pi) C_i^P. \end{aligned} \quad (3.9)$$

which is **the 1st and the 2nd kind integrability conditions ($\theta = 1, 2$) of derivational equation (2.4)**.

a') By multiplying of the previous equation with B_λ^i we get

$$\begin{aligned} R_{2imn}^p C_p^A B_i^\lambda B_\mu^m B_\nu^n &= \Omega_{\lambda\mu|_3}^A - \Omega_{\lambda\nu|_3}^A \\ R_{1imn}^p C_p^A B_i^\lambda B_\mu^m B_\nu^n &= \Omega_{\lambda\mu|_4}^A - \Omega_{\lambda\nu|_4}^A. \end{aligned} \quad (3.10)$$

that is **one more form of the 1st Codazzi equation (2.18)**.

b') Multiplying (3.9) with C_L^i , we get

$$\begin{aligned} R_{2imn}^p C_p^A C_L^i B_\mu^m B_\nu^n - \bar{R}_{L\mu\nu}^A &= \Omega_{\pi\mu}^A \hat{\Omega}_{L\nu}^\pi - \Omega_{\pi\nu}^A \hat{\Omega}_{L\mu}^\pi \\ R_{1imn}^p C_p^A C_L^i B_\mu^m B_\nu^n - \bar{R}_{L\mu\nu}^A &= \Omega_{\pi\mu}^A \hat{\Omega}_{L\nu}^\pi - \Omega_{\pi\nu}^A \hat{\Omega}_{L\mu}^\pi \end{aligned} \quad (3.11)$$

and this is **another form of the 2nd Codazzi equation of the 1st and the 2nd kind**.
3.2. If one takes $\theta = 3, \omega = 4$ in (3.1) and compares obtained equation with (3.5), we obtain **3rd integrability condition of derivational equation (2.3)** in the structure $(X_M \subset L_N, \nabla_\theta, \theta \in \{3, 4\})$:

$$R_{4p\mu\nu}^i C_A^p - \bar{R}_{A\mu\nu}^P C_P^i = -(\hat{\Omega}_{A\mu|_4}^\pi - \hat{\Omega}_{A\nu|_3}^\pi) B_\pi^i - (\hat{\Omega}_{A\mu}^P \hat{\Omega}_{4\pi\nu}^P - \hat{\Omega}_{A\nu}^P \hat{\Omega}_3^P \pi_\mu) C_P^i. \quad (3.12)$$

a) Multiplying (3.12) with B_i^λ , we get

$$R_{4p\mu\nu}^i B_i^\lambda C_A^p = -\hat{\Omega}_{A\mu|_4}^\lambda + \hat{\Omega}_{A\nu|_4}^\lambda, \quad (3.13)$$

which is **one more form of (2.17)**.

b) Multiplying (3.12) with C_i^L , we have

$$R_{4p\mu\nu}^i C_i^L C_A^p - \bar{R}_{A\mu\nu}^L = \Omega_{\pi\mu}^L \hat{\Omega}_{A\nu}^\pi - \Omega_{\pi\nu}^L \hat{\Omega}_{A\mu}^\pi. \quad (3.14)$$

which is **the 2nd Codazzi equation of the 3rd kind**.

3.2'. Endly, we put $\theta = 3, \omega = 4$ into (3) and compare obtained equation with (3.4'). In that manner, one obtains **the 3rd kind integrability condition of derivational equation (2.4)** in the structure $(X_M \subset L_N, \nabla_\theta, \theta \in \{3, 4\})$:

$$R_{3iv\mu}^p C_p^A + \bar{R}_{P\mu\nu}^A C_i^P = -(\Omega_{\pi\mu|_4}^A - \Omega_{\pi\nu|_3}^A) B_i^\pi - (\Omega_{\pi\mu}^A \hat{\Omega}_{P\nu}^\pi - \Omega_{\pi\nu}^A \hat{\Omega}_3^P \pi_\mu) C_i^P.$$

a') If one multiplies (3) with B_λ^i , it follows that

$$R_{3iv\mu}^p C_p^A B_\lambda^i = -\Omega_{\pi\mu|_4}^A + \Omega_{\pi\nu|_3}^A, \quad (3.15)$$

and this is **another form of (2.27) or (3.13)**.

$b')$ Multiplying (3) with C_L^i , we have

$$R_{i\mu\nu}^p C_A^p C_L^i - \bar{R}_{L\mu\nu}^A = \hat{\Omega}_{3L\mu}^\pi \Omega_{\pi\nu}^A - \hat{\Omega}_{4L\nu}^\pi \Omega_{3\pi\mu}^A. \quad (3.16)$$

which is **another form of the 2nd Codazzi equation of the 3rd kind** i.e. of (3.14).

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