

# SCIENTIFIC REPORTS

OPEN

## Electron irradiation induced amorphous SiO<sub>2</sub> formation at metal oxide/Si interface at room temperature; electron beam writing on interfaces

S. Gurbán<sup>1</sup>, P. Petrik<sup>1</sup>, M. Serényi<sup>1</sup>, A Sulyok<sup>1</sup>, M. Menyhárd<sup>1</sup>, E. Baradács<sup>2</sup>, B. Parditka<sup>2</sup>, C. Cserhádi<sup>2</sup>, G. A. Langer<sup>2</sup> & Z. Erdélyi<sup>2</sup>

Al<sub>2</sub>O<sub>3</sub> (5 nm)/Si (bulk) sample was subjected to irradiation of 5 keV electrons at room temperature, in a vacuum chamber (pressure  $1 \times 10^{-9}$  mbar) and formation of amorphous SiO<sub>2</sub> around the interface was observed. The oxygen for the silicon dioxide growth was provided by the electron bombardment induced bond breaking in Al<sub>2</sub>O<sub>3</sub> and the subsequent production of neutral and/or charged oxygen. The amorphous SiO<sub>2</sub> rich layer has grown into the Al<sub>2</sub>O<sub>3</sub> layer showing that oxygen as well as silicon transport occurred during irradiation at room temperature. We propose that both transports are mediated by local electric field and charged and/or uncharged defects created by the electron irradiation. The direct modification of metal oxide/silicon interface by electron-beam irradiation is a promising method of accomplishing direct write electron-beam lithography at buried interfaces.

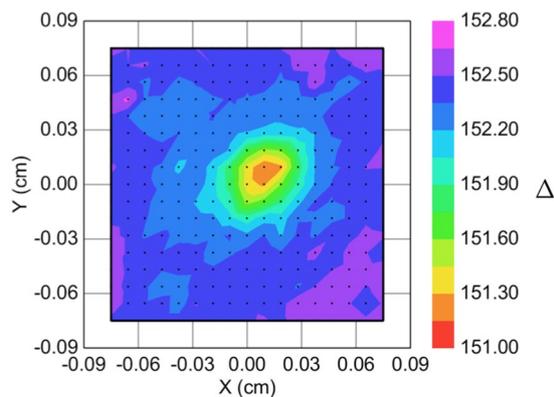
Due to the paramount role of SiO<sub>2</sub> in integrated circuits its formation is extremely well studied and understood. The really good quality silicon dioxide is produced at elevated temperatures to cope with the activation energy of its formation. It is well known, however, that even at room temperature native oxide forms at the surface of Si. The growth of native oxide terminates at a thickness of 2–3 nm. It has also been shown that SiO<sub>2</sub> growth at the surface of Si at low temperatures greatly facilitated if the oxygen molecule is excited by some means. This excitation can be carried out by various ways e.g. various plasma processes<sup>1</sup>, noble-gas ion bombardment<sup>2</sup> electron irradiation<sup>3,4</sup>, etc. Much less is known about the formation of SiO<sub>2</sub> at metal oxide/silicon interface.

Aluminum oxide/Si interface has been frequently studied for at least two reasons: a./in a quest of high dielectric constant material to replace the SiO<sub>2</sub><sup>5</sup>, b./for passivation of the Al<sub>2</sub>O<sub>3</sub>/Si interface in photovoltaic applications<sup>6–9</sup> and quantum dots<sup>10</sup>. In both cases, the presence of fixed charges inherent to Al<sub>2</sub>O<sub>3</sub> causes a problem<sup>11</sup>, which can be addressed by various means<sup>12</sup>. The interface of the as deposited (chemical vapor deposition at 400 °C) system of a 3.5 nm thick Al<sub>2</sub>O<sub>3</sub> film on Si(100) was studied by Klein *et al.*<sup>13</sup>, using nuclear reaction resonance profiling. They could detect the presence of a silicate layer. The theoretical calculation of Xiang *et al.*<sup>14</sup> showed the existence of a sharp Al<sub>2</sub>O<sub>3</sub>/Si interface exhibiting Si-O and Si-Al covalent bonds. Thus, in these works Si-O bond formation has been observed at the interface, but SiO<sub>2</sub> formation has not been reported.

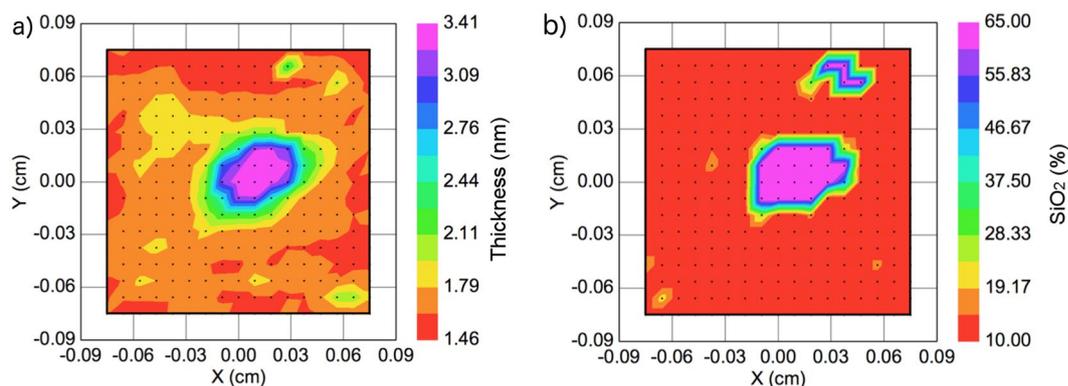
On the other hand medium energy electron and/or X ray irradiation might cause various chemical changes (including oxidation), defect formation, structural changes in the surface close regions of the irradiated material, mainly in the case of insulators, where recombination and healing are limited. Defect formation<sup>15,16</sup> is explained by bond breaking, while structural changes like amorphization<sup>15,17</sup> is explained by excitation of the chemical bonds and subsequent atomic movements.

In this paper, we will report on the electron irradiation induced amorphous SiO<sub>2</sub> growth at the Al<sub>2</sub>O<sub>3</sub>/Si interface. To demonstrate this process samples were made by growing a 5 nm thick Al<sub>2</sub>O<sub>3</sub> layer onto a carefully cleaned

<sup>1</sup>Institute for Technical Physics and Materials Science, Centre for Energy Research Hungarian Academy of Sciences, P.O.B. 49, H-1525, Budapest, Hungary. <sup>2</sup>Department of Solid State Physics, University of Debrecen, P.O. Box 400, H-4002, Debrecen, Hungary. P. Petrik, M. Menyhárd and Z. Erdélyi contributed equally to this work. Correspondence and requests for materials should be addressed to M.M. (email: [menyhard.miklos@energia.mta.hu](mailto:menyhard.miklos@energia.mta.hu))



**Figure 1.** Map of the measured ellipsometric angles,  $\Delta$ , around the illuminated spot at the wavelength of 300 nm.



**Figure 2.** Maps of thickness of the interface layer, (a) and the volume fraction of  $a\text{:SiO}_2$  in the interface layer, (b).

(100) Si substrate by two methods, by atomic layer deposition (ALD) and by radio frequency (RF) sputtering. The samples were then irradiated by electrons in a vacuum system (pressure of  $10^{-9}$  mbar) at room temperature. The parameters of electron irradiation were: energy 5 keV, current density up to  $6 \times 10^{16}$  electrons/cm<sup>2</sup>/s, total fluence up to  $6 \times 10^{21}$  electrons/cm<sup>2</sup>. The vacuum system was equipped with electron and ion guns and an electron analyzer. The effect of irradiation was studied *in situ* by Auger Electron Spectroscopy (AES), AES depth profiling, and *ex situ* by spectroscopic ellipsometry (SE), and transmission electron microscopy (TEM).

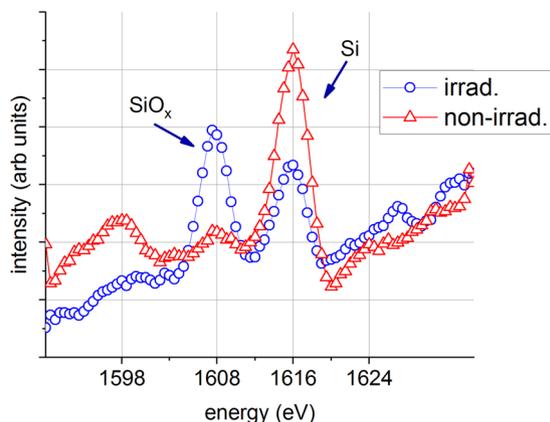
## Results

**Transmission Electron Microscopy (TEM).** TEM images were taken before and after irradiation; they did not show any crystalline phase except the Si substrate. Thus, both phases,  $\text{Al}_2\text{O}_3$  and  $\text{SiO}_2$ , observed were amorphous. To emphasize that the  $\text{SiO}_2$  produced is amorphous, we will sign it as  $a\text{:SiO}_2$ .

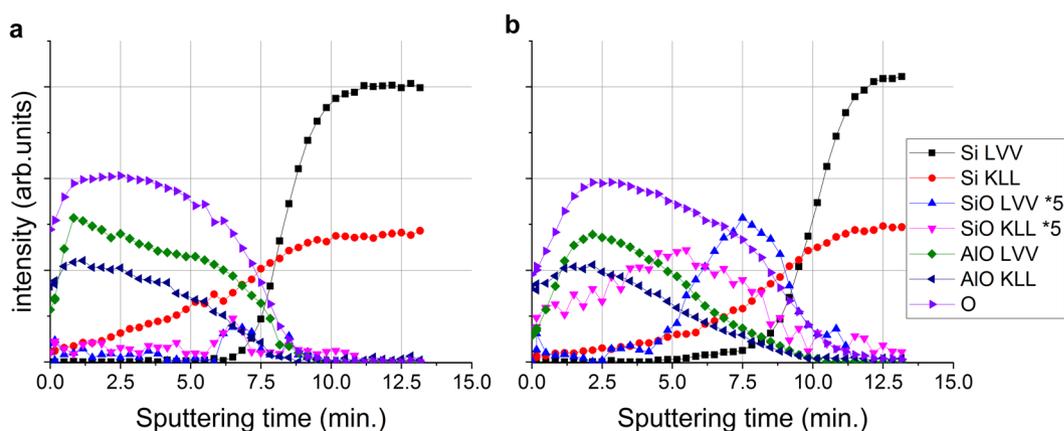
**Spectroscopic ellipsometry (SE).** The irradiated spot can clearly be identified on the map of measured ellipsometric angles,  $\Delta$ , in Fig. 1, revealing that the irradiation caused a significant change in the optical properties of the layers. The difference in  $\Delta$  between the irradiated and non-irradiated regions is larger than one degree, thus more than an order of magnitude higher than the sensitivity of the measurement.

For determination of the thickness of the amorphous  $\text{SiO}_2$  ( $a\text{:SiO}_2$ ) the interface was modeled as an effective medium composition of  $\text{Al}_2\text{O}_3$ ,  $\text{SiO}_2$  and Si<sup>18</sup>. Figure 2 shows that both the interface thickness,  $a$ , and the volume fraction of  $a\text{:SiO}_2$ ,  $b$ , increase in the irradiated spot. Due to small thickness and the correlation between the  $\text{Al}_2\text{O}_3$  and the  $a\text{:SiO}_2$  components, there might be a significant error in the determined volume fraction of  $a\text{:SiO}_2$ . However, the trend of the increasing interface thickness and volume fraction of  $a\text{:SiO}_2$  in the irradiated part is very clear from the ellipsometry results.

**Auger Electron Spectroscopy (AES).** The samples were produced either by ALD or RF sputtering. The results did not depend on the production technology and thus in the following the type of sample will not be specified. Figure 3 shows the Auger spectra,  $N(E)$ , in the  $1600 \pm 20$  eV range (where the  $\text{Si}_{\text{KLL}}$  Auger line appears independently from its chemical state) recorded on the surface of the irradiated ( $1.32$  electrons/s/ $\text{Å}^2 \times 20$  hours =  $9.4 \times 10^{20}$  electrons/cm<sup>2</sup>) and non-irradiated regions of the sample. It is clear that the intensity of the KLL Auger electrons emitted by Si in oxide environment increased considerably due to the irradiation.



**Figure 3.** The raw data, measured intensity ( $N(E)$ ) vs energy ( $E$ ) in the vicinity of the Si KLL transition obtained from the non-irradiated region of the surface (non-irrad.) and after an irradiation of  $9.4 \times 10^{20}$  electrons/cm<sup>2</sup> (irrad.).

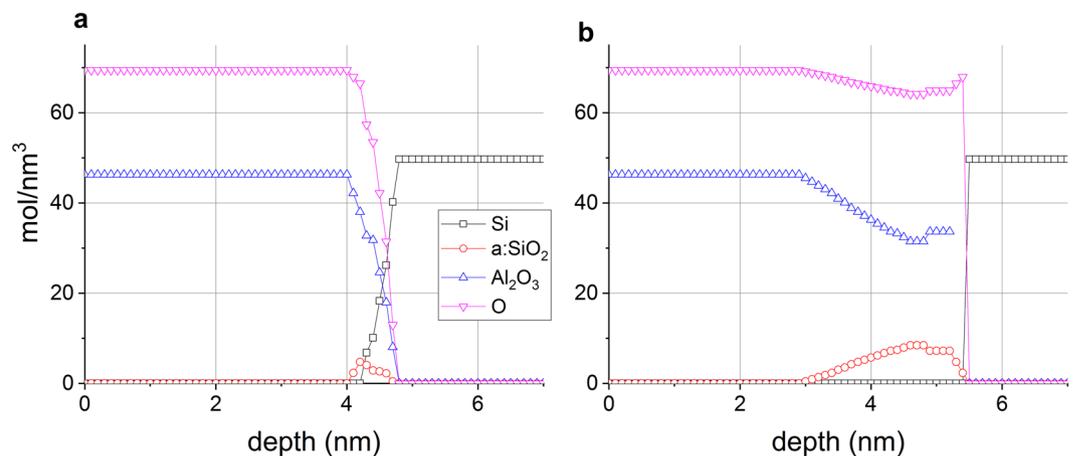


**Figure 4.** The raw data, peak-to-peak intensities measured in the  $N'(E)$  curve, of AES depth profiling obtained from non-irradiated (a) and irradiated (b) regions of the sample. SiO and AlO mean Auger intensities of Si and Al being in oxygen environment, respectively.

Simultaneously, the intensity of the KLL Auger electrons emitted by Si in metal environment decreased. The rough interpretation of the finding is simple; within the information depth of the  $Si_{KLL}$  Auger electrons a fraction of the Si has been converted to  $SiO_x$ . The x value cannot be accurately determined by AES but it is not far from 2. From the TEM results it is also clear that the compound formed is amorphous. Thus we will sign the compound as  $a:SiO_2$ , which means an amorphous phase with composition close to the stoichiometric. However, since the measured Auger intensity is a weighted sum of the intensities emitted from various depths, the actual depth distribution of the  $a:SiO_2$  produced cannot be derived from this measurement.

For the determination of the depth distribution of the  $a:SiO_2$  produced by the electron irradiation AES depth profiling has been utilized. In AES depth profiling the surface of the sample is etched away in predefined steps by using ion sputtering and the newly exposed surface is analyzed by AES which provides elemental and some chemical information.

Figure 4 shows the raw data (measured Auger peak-to-peak intensities in the differentiated,  $N'(E)$  curve) obtained from the non-irradiated (a) and irradiated (b) (total electron dose of  $9.4 \times 10^{20}$  electrons/cm<sup>2</sup>) regions of the sample; the two depth profiles show considerable differences. The most important one is that while in the case of the non-irradiated region the intensity of the Si Auger signal in oxide environment (signed by SiO) is just higher than the noise level around the  $Al_2O_3/Si$  interface, in the case of the irradiated region the same signal is easily measurable, proving that due to the irradiation  $a:SiO_2$  has formed. It should also be noted that longer sputtering time is necessary to remove the layer at the irradiated region than that necessary in the case of the non-irradiated one, demonstrating that the thickness of the top most layer in the irradiated region (pure  $Al_2O_3$  + altered layer +  $a:SiO_2$ ) of the sample is greater than that of the virgin part. Figure 5 shows the concentration distributions, which best correspond to the depth profiles shown in Fig. 4 measured on the non-irradiated (virgin) and irradiated region of the sample, respectively, calculated by our trial and error evaluation method<sup>19,20</sup>. Note that the irradiation took place at room temperature.



**Figure 5.** The concentration distributions of the sample at the non-irradiated (a) and irradiated (b) regions, respectively derived from the depth profiles shown in Fig. 4a and b.

Figure 5a shows that the structure of the non-irradiated sample is the expected one; there is an about 5 nm thick  $\text{Al}_2\text{O}_3$  layer on the top of the Si substrate. The small amount of a:SiO<sub>2</sub> which is at the interface region is due to either the non-proper cleaning of the Si substrate or oxidation at the beginning of the deposition of the  $\text{Al}_2\text{O}_3$ . On the other hand, the irradiated region provides a strongly altered structure. Though the surface is covered by a pure  $\text{Al}_2\text{O}_3$  layer its thickness is only about 3 nm instead of the 5 nm initial thickness. This cover layer is followed by an  $\text{Al}_2\text{O}_3$  and a:SiO<sub>2</sub> mixture in which the concentration of the a:SiO<sub>2</sub> increases toward the interface. The overall thickness of the altered layer and remainder  $\text{Al}_2\text{O}_3$  layer is larger than that of the initial  $\text{Al}_2\text{O}_3$  layer.

## Discussion

Both methods (SE and AES) clearly show that the irradiated part of the sample has considerably changed: a:SiO<sub>2</sub> has been built due to the irradiation. Since the non-destructive ellipsometric measurement cannot produce a:SiO<sub>2</sub> it is clear that the a:SiO<sub>2</sub> is the result of the electron irradiation.

An obvious question is whether the subsequent AES studies combined with ion bombardment might impair the actual distribution of the a:SiO<sub>2</sub> formed during the electron beam irradiation experiment. Considering Fig. 5b we can estimate the overall rate of the process; it is rather low being in the range of  $10^{-4}$ – $10^{-5}$  a:SiO<sub>2</sub>/5 keV electrons. Thus, the production of SiO<sub>2</sub> during AES measurement can be ignored since the irradiation fluence during the recording of the AES spectra is 300 times lower than that of the irradiation experiment.

The ion bombardment can also produce SiO<sub>2</sub>. The efficiency of the process can also be easily checked experimentally. The measured Auger intensities recorded on the surface before Auger depth profiling cannot be used to determine the distribution of the emitting elements. On the other hand, the depth distributions determined by the AES depth profiling can easily be integrated and they should give the intensities measured at the surface before the AES depth profiling. This integral has been calculated for all measurements and compared with the intensities measured on the surface. The amount of a:SiO<sub>2</sub> calculated from the AES measurement performed on the surface before the AES depth profiling started, and that obtained for the AES depth profile agreed within an error of 20%, meaning that the possible effect of the ion bombardment on the a:SiO<sub>2</sub> content is typically still less than 20%, which will be ignored. Thus, the concentration distributions provided by the AES depth profiling can be accepted as those produced by the 5-kV electron irradiations.

The accurate determination of heating effect of the electron beam irradiation is rather difficult<sup>15</sup> in case of high current densities and small particles. In our case the current density was chosen to be rather low, which allows to make a rough overestimation of the temperature rise,  $\Delta T$ . It can be supposed that the whole irradiation energy is absorbed by the sample and used for heating. In this case in stationary case the input energy is equal to that carried away by heat conduction which results in the simple equation of  $\Delta T = Pr/K^{21}$ , where  $P$  is the irradiation power density,  $r$  is the radius of the electron beam, and  $K$  is the heat conduction. In the present case  $\Delta T$  is around 3 K, which is negligible, and consequently all processes are taking places at room temperature.

The basic features of the a:SiO<sub>2</sub> growth (due to the electron irradiation) were similar on samples prepared by different technologies. Thus, we ignore possible effects of the imperfections connected to the sample production technology, and the process will be explained assuming perfect interfaces. The growth of a:SiO<sub>2</sub> can be divided into two steps: creating charged and/or neutral oxygen and the growth process itself.

**Electron irradiation induced bond breaking & excitation.** Vast amount of experimental data has been amassed in transmission electron microscopy (TEM) studies concerning various electron beam-induced damages in insulators<sup>15</sup>. Several electron irradiation-induced effects, including phase transformation, decomposition, amorphization, oxidation, reduction, etc., have been identified. For the description of the phenomena two basic models have been developed: a./the mechanical interaction of the bombarding electron with the nuclei, knock-on mechanism, b./the electronic excitation by the electric field, radiolytic processes.

The knock-on process in our case will not be considered since the energy, 5 keV, of our bombarding electrons is far below the threshold of this mechanism<sup>15</sup>. On the other hand, the radiolytic processes have practically no

threshold since because of the multiple interactions nearly any type of excitations can happen. E.g., recently the dissolution of boehmite under the high energy electron irradiation (TEM studies) was explained by the electron-hole pairs created during the electron bombardment<sup>22</sup>. This explanation is rather similar to those developed long time ago for the description of low-energy electron-stimulated desorption (ESD)<sup>23–25</sup> which can readily be applied for the explanation of our observations. According to these models following the ionization of a core electron, a valence electron from the O decays to fill the resulting hole in frame of the Auger process. Therefore,  $O^{2-}$  transforms to  $O^0$ . It might also happen that by an additional Auger effect a further electron is emitted and  $O^+$  is formed for a short while<sup>23–25</sup>. The lifetime of both  $O^0$  and  $O^+$  is long, since fast recombination is not possible due to the insulator matrix. During the excitation processes, because of Columbic repulsion, structural changes might also occur. The electron irradiation this way produces neutral and/or charged O in the  $Al_2O_3$  matrix together with various charged and neutral crystal defects.

**SiO<sub>2</sub> growth.** If Si surface containing O<sub>2</sub> and/or H<sub>2</sub>O is irradiated by electrons, SiO<sub>2</sub> can be produced<sup>3,4</sup>. The description of the process is similar to that of ESD. In this case it is assumed that electrons attach to an adsorbed O<sub>2</sub> molecular precursor to form O<sub>2</sub><sup>-</sup>. The O<sub>2</sub><sup>-</sup> then decomposes to form O and O<sup>-</sup>, and one or both of these species cause rapid oxidation of the surface. In our case neutral and/or charged O are produced in the  $Al_2O_3$  matrix and for the compound formation to take place they should be transported to the interface.

The diffusion of O in  $\alpha$  alumina has been studied by Sokol *et al.*<sup>26</sup> showing, that due to the structural and spin configuration, the defect reaction energy can change by over 2 eV. This behavior affects the equilibrium defect concentrations by many orders of magnitude. Consequently, the diffusion processes in such materials may be more complicated, which has previously been assumed. Århammar *et al.* reached similar conclusion for amorphous oxides<sup>27</sup>. Still it seems that at room temperature neither the O nor the Si can be transported fast enough by thermal diffusion to explain the growth of SiO<sub>2</sub> layer at the  $Al_2O_3$ /Si interface.

It is well known that despite the high heat of the SiO<sub>2</sub> formation the native oxide forms on the surface of clean Si at room temperature in air. The thickness of this oxide is about 2 nm. This process is explained by the theory of Cabrera and Mott, which assumes that the electrons can pass freely from the Si to the oxide surface to ionize oxygen atoms. This establishes a uniform field within the oxide, which leads to a shift in the Fermi level of the oxide<sup>28</sup>. The same reasoning can be used if electron is placed to the surface from any other source. Nowak *et al.* have shown that electron bombardment, providing charges to build up electric field, induces oxide growth on tungsten nanowires at room temperature<sup>29,30</sup>. They use the explanation of Mott and Cabrera; the electric field created reduces the energy barriers for the migration of metal cations or oxygen anions into and through the oxide, allowing significant material transport and thus growth of the oxide layer at low temperature. It is also evident that this is a self-controlled process; after reaching a given thickness the strength of electric field is no more sufficient to drive the diffusion.

Our  $Al_2O_3$  layer is 5 nm thick that is thicker than the native SiO<sub>2</sub>. In our case, however, the charged defects, charged and neutral oxygen produced by the electron bombardment are distributed evenly in the layer. That is, in our case various distributions of local fields and charged as well as neutral oxygen is produced. It is evident that those oxygen atoms, which are close to the interface can be transported to the substrate Si atoms and form oxide at a high probability; thus the oxide growth starts from the interface. Similarly, if the charged defect is close to the interface, then an electric field of sufficient strength is built up initiating the diffusion of Si into the  $Al_2O_3$  layer. Since in the defected  $Al_2O_3$  layer there are charged and/or neutral oxygen atoms, compound formation might take place. The a:SiO<sub>2</sub> grain grows by adding additional Si and O atoms. The primary Si source is the substrate. It should be remembered, however, that the self-limitation process is activated after the a:SiO<sub>2</sub> grain reaches a certain thickness and the local field strength is not sufficient anymore to drive the diffusion. It seems that our particles are somewhat larger than that of the typical thickness of the native oxide. On the other hand, one should also consider that the already produced a:SiO<sub>2</sub> grains are also subjected to electron irradiation producing charged defects, excited oxygen atoms and quasi free Si in the a:SiO<sub>2</sub> grain. The quasi free Si can utilize the local fields to diffuse to the surface of the a:SiO<sub>2</sub> grain resulting in further growth. The probability of this process is lower than that of the primary one, however, and the rate of growth in this phase is much lower resulting in only some additional growth resulting in a 3–4 nm thick a:SiO<sub>2</sub> grains.

## Conclusions

We have shown that bombarding 5 nm thick  $Al_2O_3$ /Si structure by 5 keV electrons at room temperature, amorphous silicon dioxide is produced. The amorphous SiO<sub>2</sub> grains grow from the interface toward the  $Al_2O_3$  matrix; their amount depends on the irradiated charge. The phenomenon was explained considering electron bombardment-induced bond breaking in  $Al_2O_3$ , electric field driven diffusion of Si and O in defected regions of the  $Al_2O_3$  based on the Cabrera-Mott theory.

The direct modification of metal oxide/silicon interface by electron irradiation is a promising method of accomplishing direct write lithography at buried interfaces.

**Samples and Methods.** *Samples.* Samples were made by growing an  $Al_2O_3$  layer on a Si (100) substrate ALD and RF sputtering.

The ALD layers were made by a Beneq TFS 200 ALD reactor in the plasma-enhanced deposition mode. Trimethylaluminium (TMA – from Sigma Aldrich) together with high purity oxygen gas was used for deposition. Prior to the sample preparation, the deposition chamber had been heated up to 150 °C. During the deposition, the pressure inside the vacuum chamber was 9.5 mbar, while in the reactor chamber 1.1 mbar. The RF power of the plasma was set to be 50 W and the flow rate of the oxygen was 100 sccm. The ALD cycle was the following: 150 ms TMA, which was followed by a 2 s purge, then a 2 s oxygen plasma at 50 W and at the end of the cycle another 2 s purge.

	LVV		KLL	
	Oxide	Metal	Oxide	Metal
Al	54	68	1389	1396
Si	78	92	1610	1619

**Table 1.** The LVV and KLL Auger lines energies (eV) of Al and Si in pure and oxidized forms.

The RF sputter deposition was carried out in a Leybold Z400 apparatus evacuated to  $5 \times 10^{-5}$  mbar. Sputtering was performed under a mixture of high purity argon and oxygen gases with an applied RF power of 255 W yielding a plasma pressure of  $2.5 \times 10^{-2}$  mbar. Oxygen was continuously let into the sputtering chamber at flow rates resulting in a partial oxygen pressure of 6%. The deposited amorphous  $\text{Al}_2\text{O}_3$  film has a high refractive index and low absorption coefficient<sup>31</sup>.

**Electron irradiation and Auger Electron Spectroscopy.** All electron irradiation experiments have been carried out in our standard vacuum system equipped with electron energy analyzer (STAIB OPC 105) and various electron sources. The electron irradiation was facilitated by an electron gun (STAIB EK-10-M) with the following parameters: energy 5 keV, current density up to  $6 \times 10^{16}$  electrons/cm<sup>2</sup>/s fluence about up to  $6 \times 10^{21}$  electrons/cm<sup>2</sup>, angle of incidence 54° (with respect to surface normal). The irradiated area was  $100 \times 100 \mu\text{m}^2$ . For AES analysis the same electron gun was used for the excitation with a current density and energy of  $1 \times 10^{16}$  electrons/cm<sup>2</sup>/s and 5 keV, respectively. The fluence during the recording of the Auger spectra is about 300 times less than that used for irradiation.

The Auger spectra,  $N(E)$ , were recorded in counting mode. The recorded spectrum was numerically differentiated,  $N'(E)$ , for performing the concentration calculation.

The Auger signals of  $\text{Al}_{\text{KLL}}$ ,  $\text{Al}_{\text{LVV}}$ ,  $\text{Si}_{\text{KLL}}$ ,  $\text{Si}_{\text{LVV}}$ , C and O, were measured; by measuring the high (KLL) and low (LVV) escape depths Auger electrons the quality of evaluation of the depth distributions in AES depth profiling considerably improves. The energies of the Al and Si Auger electrons strongly depend on their chemical environment; the corresponding Auger electron energies (in eV) are shown in Table 1.

Therefore it is easy to determine the metal and oxide fractions of the elements by measuring either the LVV or KLL Auger electrons allowing the determination of the depth distributions for the metal and oxide components separately.

**AES depth profiling.** A low energy ion gun of Technoorg Linda was used for AES depth profiling. The parameters of the ion bombardment used were: energy 1 keV, projectile Ar, angle of incidence (with respect to the surface normal) 80° and specimen rotation during ion bombardment. The ion beam was scanned in an area of  $1.5 \times 1.5 \text{ mm}^2$ . Using these parameters, the ion bombardment induced roughening and mixing is minimal<sup>32</sup>.

**Determination of the concentration distribution from AES spectra.** In AES, generally simple expressions are applied for the evaluation of the composition using the measured peak-to-peak amplitudes of the differentiated,  $N'(E)$ , curve<sup>33</sup>. This expression assumes that the composition within the escape depth of signal electrons is homogeneous; in any other case it cannot be applied. This is however the situation presently since the thickness of the Al oxide layer is only 5 nm, while the inelastic mean free path (IMFP) of the  $\text{Al}_{\text{KLL}}$  and  $\text{Si}_{\text{KLL}}$  Auger electrons in  $\text{Al}_2\text{O}_3$  are 3.2 and 3.6 nm, respectively<sup>34</sup>. We used a trial and error approach to determine the composition distribution of our sample<sup>19,20</sup>. The essence of the method is that we assume a composition distribution along the depth and calculate the Auger intensities assuming that the transport of electrons can be described by the exponential attenuation law, not considering the elastic scattering. (Neglecting the elastic scattering creates an error in the range of 10–15%, which will not affect the description of the phenomena.) The composition distributions are varied until the simulated depth profile is close enough to the measured one. If an element emits high (high IMFP) and low energy (low IMFP) Auger electrons, as in the present case, the accuracy of the method is rather good.

In case of depth profiling the above procedure is repeated for all spectra obtained after each sputtering steps assuming that the ion bombardment used for removing the material does not cause serious changes to the material. This is a reasonable assumption since the removed layer thickness is less than 8 nm and all alterations scale with the removed layer thickness<sup>32</sup>.

**Spectroscopic ellipsometry (SE).** Auger depth profiling uses ions and electrons to reveal the depth distribution of the  $\text{SiO}_2$ . Both projectiles may initiate the formation of  $\text{SiO}_2$ . Though it will be shown that these are low probability processes, still we have applied SE, a non-destructive method, to verify the presence of the  $\text{SiO}_2$  produced by electron irradiation. The SE measurements have been carried out by a Woollam M-2000DI rotating compensator ellipsometer at an angle of incidence of 55°. The microspot option was used to focus the light into a spot with a diameter of approximately 0.2 mm. The surface around the irradiated region was mapped with steps smaller than the spot size, in order to precisely locate the modified spot on the sample surface, and to measure only the irradiated region of the sample.

## References

- Pan, P., Nesbit, L. A., Douse, R. W. & Gleason, R. T. The Composition and Properties of ECVD Silicon Oxide Films. *J. Electrochem. Soc.* **132**, 2012–2019 (1985).
- Engstrom, J. R., Bonser, D. J. & Engel, T. The reaction of atomic oxygen with Si (100) and Si (111): II. Adsorption, passive oxidation and the effect of coincident ion bombardment. *Surf. Sci.* **268**, 238–264 (1992).
- Xu, J., Choyke, W. J. & Yates, J. T. Jr. Enhanced silicon oxide film growth on Si (100) using electron impact. *J. Appl. Phys.* **82**, 6289–6292 (1997).

4. Popova, I., Zhukov, V. & Yates, J. T. Jr. Electron-stimulated conversion of chemisorbed O to Al<sub>2</sub>O<sub>3</sub> on Al (111). *Appl. Phys. Lett.* **75**, 3108–3110 (1999).
5. Wilk, G. D., Wallace, R. M. & Anthony, J. M. High-gate dielectrics: Current status and materials properties considerations. *J. Appl. Phys.* **89**, 5243–5275 (2001).
6. Choi, S. *et al.* Structural evolution of tunneling oxide passivating contact upon thermal annealing. *Sci. Rep.* **7**, 12853 (2017).
7. Agostinelli, G. *et al.* Very low surface recombination velocities on p-type silicon wafers passivated with a dielectric with fixed negative charge. *Sol. Energy Mater. Sol. Cells* **90**, 3438–3443 (2006).
8. Chen, M. J., Shih, Y. T., Wu, M. K. & Tsai, F. Y. Enhancement in the efficiency of light emission from silicon by a thin Al<sub>2</sub>O<sub>3</sub> surface-passivating layer grown by atomic layer deposition at low temperature. *J. Appl. Phys.* **101**, 033130 (2007).
9. Hoex, B. *et al.* Excellent passivation of highly doped p-type Si surfaces by the negative-charge-dielectric Al<sub>2</sub>O<sub>3</sub>. *Appl. Phys. Lett.* **91**, 112107 (2007).
10. Spruijtenburg, P. C., Sergey, V. A., Mueller, F., van der Wiel, W. G. & Zwanenburg, F. A. Passivation and characterization of charge defects in ambipolar silicon quantum dots. *Sci. Rep.* **6**, 38127 (2016).
11. Johnson, R. S., Lucovsky, G. & Baumvol, I. Physical and electrical properties of noncrystalline Al<sub>2</sub>O<sub>3</sub> prepared by remote plasma enhanced chemical vapor deposition. *J. Vac. Sci. Technol. A* **19**, 1353–1360 (2001).
12. Benick, J. *et al.* Effect Of A Post-Deposition Anneal On Al<sub>2</sub>O<sub>3</sub>/Si Interface Properties. <https://doi.org/10.1109/PVSC.2010.5614148> (2010).
13. Klein, T. M. *et al.* Evidence of aluminum silicate formation during chemical vapor deposition of amorphous Al<sub>2</sub>O<sub>3</sub> thin films on Si (100). *Appl. Phys. Lett.* **75**, 4001–4003 (1999).
14. Xiang, H. J., Da Silva, J. L. F., Branz, H. M. & Wei, S. H. Understanding the Clean Interface Between Covalent Si and Ionic Al<sub>2</sub>O<sub>3</sub>. *Phys. Rev. Lett.* **103**, 116101 (2009).
15. Jiang, N. Electron beam damage in oxides: a review. *Rep. Prog. Phys.* **79**, 16501 (2016).
16. Gergely, G. & Menyhard, M. Charging of NaCl under low electron bombardment. *Acta Tech. Hung.* **80**, 309–315 (1975).
17. Janossy, I. & Menyhard, M. LEED study of quartz crystals. *Surf. Sci.* **25**, 647–649 (1971).
18. Petrik, P. *et al.* Comparative measurements on atomic layer deposited Al<sub>2</sub>O<sub>3</sub> thin films using *ex situ* table top and mapping ellipsometry, as well as X-ray and VUV reflectometry. *Thin Solid Films* **541**, 131–135 (2013).
19. Kotis, L., Gurban, S., Pecz, B., Menyhard, M. & Yakimova, R. Determination of the thickness distribution of a graphene layer grown on a 2" SiC wafer by means of Auger electron spectroscopy depth profiling. *App. Surf. Sci.* **316**, 301–307 (2014).
20. Zommer, L., Jablonski, A., Kotis, L., Safran, G. & Menyhard, M. Simulation and measurement of AES depth profiles; a case study of the C/Ta/C/Si system. *Surf. Sci.* **604**, 633–640 (2010).
21. Hofmann, S. & Zalar, A. Electron beam effects during the sputter profiling of thin Au/Ag films analysed by Auger Electron spectroscopy. *Thin Solid Films* **56**, 331–342 (1979).
22. Conroy, M. *et al.* Importance of interlayer H bonding structure to the stability of layered minerals. *Sci. Rep.* **7**, 13274 (2017).
23. Menzel, D. & Gomer, R. Desorption from metal surfaces by low-energy electrons. *J. Chem. Phys.* **41**, 3311–3328 (1964).
24. Hobbs, L. W. & Pascucci, M. R. Radiolysis and defect structure in electron-irradiated α-quartz. *J. Physique* **C6**, 237–242 (1980).
25. Knotek, M. L. & Feibelman, P. J. Ion desorption by Core-Hole Auger Decay. *Phys. Rev. Lett.* **40**, 964–967 (1978).
26. Sokol, A. A., Walsh, A., Richard, C. & Catlow, A. Oxygen interstitial structures in close-packed metal oxides. *Chem. Phys. Lett.* **492**, 44–48 (2010).
27. Århammar, C. *et al.* Unveiling the complex electronic structure of amorphous metal oxides. *PNAS* **108**, 6355–6360 (2011).
28. Atkinson, A. Transport processes during the growth of oxide films at elevated temperature. *Rev. Mod. Phys.* **57**, 437–470 (1985).
29. Nowak, C., Kirchheim, R. & Schmitz, G. Electric-field-induced low temperature oxidation of tungsten nanowires. *Appl. Phys. Lett.* **89**, 143104 (2006).
30. Nowak, C., Schmitz, G. & Kirchheim, R. Electric field effect on low temperature nanoscale oxidation. *Surf. Sci.* **604**, 641–648 (2010).
31. Serényi, M., Lohner, T., Sáfrán, G. & Szívós, J. Comparison in formation, optical properties and applicability of DC magnetron and RF sputtered aluminum oxide films. *Vacuum* **128**, 213–218 (2016).
32. Menyhard, M. High-depth-resolution Auger depth profiling/atomic mixing. *Micron* **30**, 255–265 (1999).
33. Childs, K.D. *et al.* Handbook of Auger Electron Spectroscopy, 3rd ed D.G. Watson Eden Prairie, Minnesota, 1995.
34. Tanuma, S., Powell, C. J. & Penn, D. R. Calculations of electron inelastic mean free paths for 31 materials. *Surf. Inter. Anal.* **11**, 577–589 (1988).

## Acknowledgements

The work is supported by the GINOP-2.3.2-15-2016-00041, NN 114422 M-ERANET “GRACE” and TÉT\_16-1-2016-0100 projects. The project is co-financed by the European Union and the European Regional Development Fund.

## Author Contributions

Z.E. and M.M., co-wrote the paper, headed and coordinated the work. E.B. and B.P. produced the ALD layers. C.Cs. and G.A.L. performed preliminary e-beam irradiation and quality test measurements. M. S. made the layers by RF sputtering. P.P. made the ellipsometry measurements and evaluation and co-wrote the paper. E.B., C. Cs., G.A.L., A. S. contributed to the discussion of the results. S.G., A.S., M.M. performed the AES studies.

## Additional Information

**Competing Interests:** The authors declare that they have no competing interests.

**Publisher's note:** Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.



**Open Access** This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons license and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this license, visit <http://creativecommons.org/licenses/by/4.0/>.

© The Author(s) 2018